

Vacuum Arcs and Gradient Limits

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Outline

We need to understand gradient limits.

Critical for many fields

Seems vital to optimize accelerator designs

What frequency ?

What material properties ?, Materials ?

Description of simple (elegant ?) model

A few simple mechanisms

No variables

Everything can be calculated

We try to understand all arcing data

Many Predictions

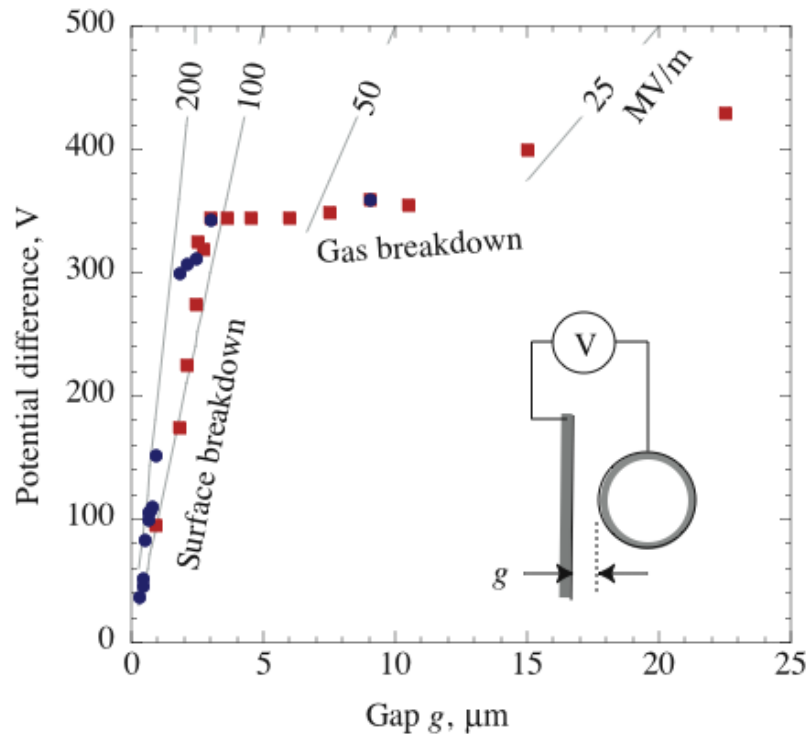
Details for much of this work have been numerically simulated.

Everything looks simple, but calculations can be difficult.

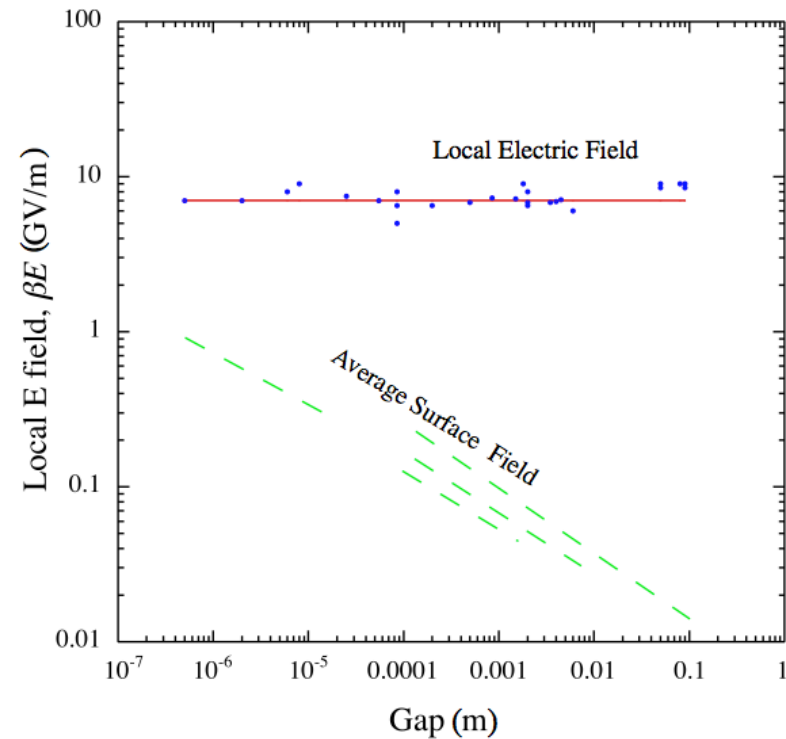
What's next?

Vacuum Arcs

Old data define the problem:
1900-1904



1964 summary



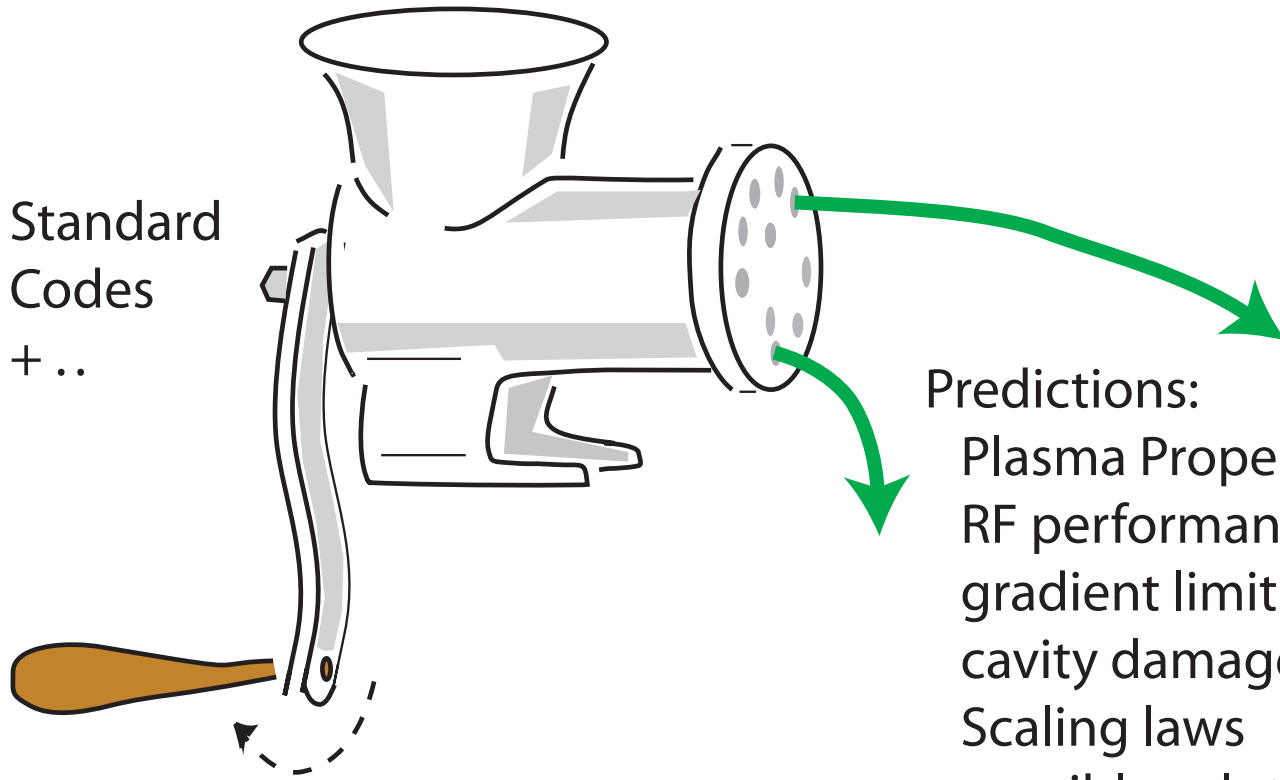
According to this data, vacuum arcs are: 1) different from gas breakdown (Michelson and Millikan) and, 2) are a single-surface phenomenon (Alpert et. al.).

They are not generally understood.

Our Model

Simple Assumptions:
Coulomb explosions
Unipolar arcs
Physics

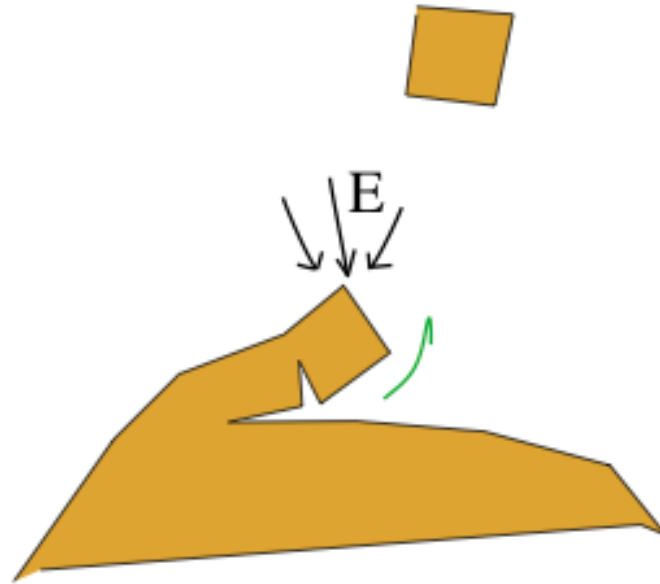
Standard
Codes
+ ..



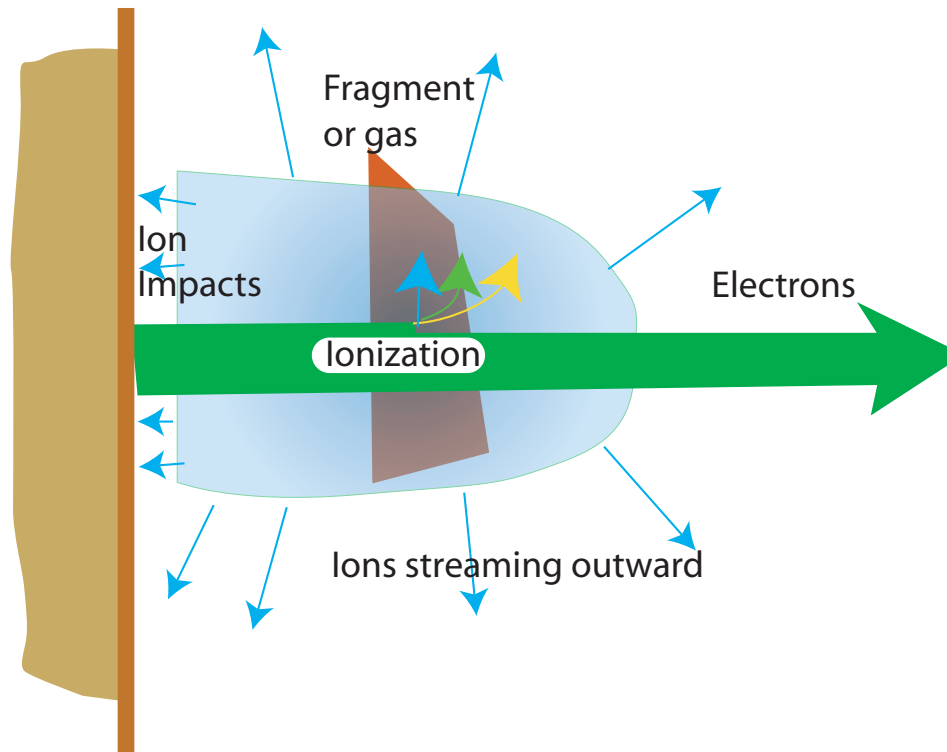
Predictions:
Plasma Properties
RF performance
gradient limits
cavity damage
Scaling laws
possible solutions

Assumptions

Coulomb Explosions



and Unipolar Arcs



Triggers for vacuum arcs

All vacuum arcs seem to look the same.

Single surface phenomena (Alpert et al)

Breakdown at ~ 10 GV/m surface field (Alpert, MTA, many others)

Tensile stress \sim tensile strength

Fatigue (creep) can contribute

Our highest surface electric fields were with 3 - 4.5 T magnetic fields

Compatible with

Fatigue, Electromigration (Antoine, Peauger), minor heating

Other arcs, Laser ablation, Lab plasmas, DC arcs, tokamaks.

No whiskers seen either in rf or DC expts.

Small emitters we see don't heat, EEE

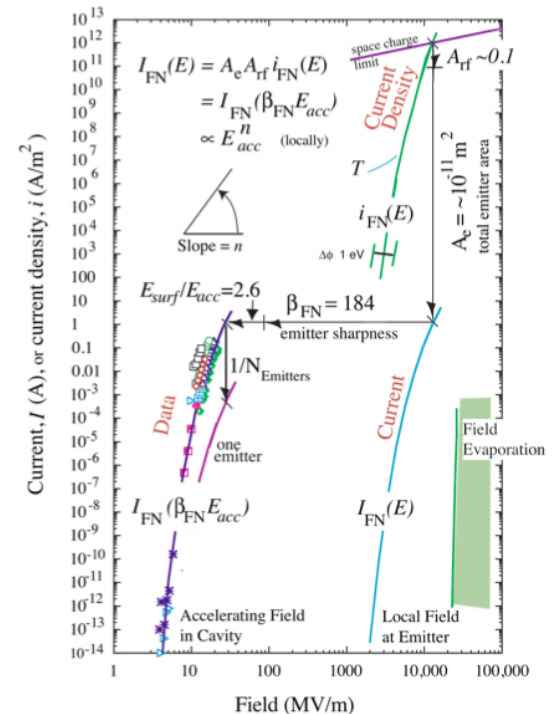
Pulse heating ~ 0 , if $d_{\text{breakdown site}} \ll$ skin depth

Breakdown is fully 3D, not 1D.

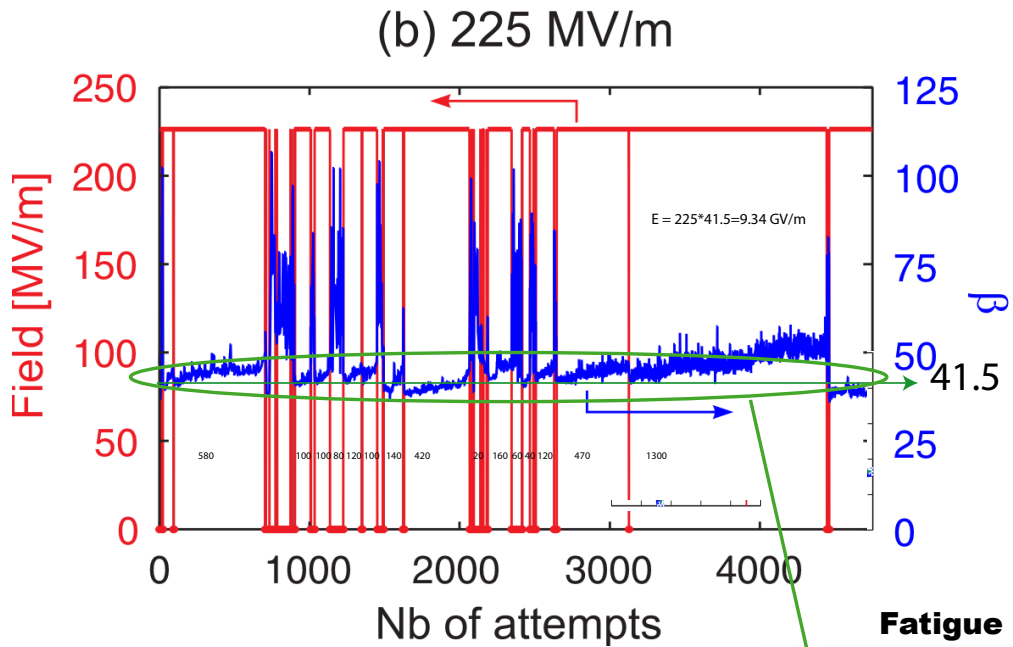
Vacuum breakdown is not gas breakdown

We do not see any keV ions anywhere.

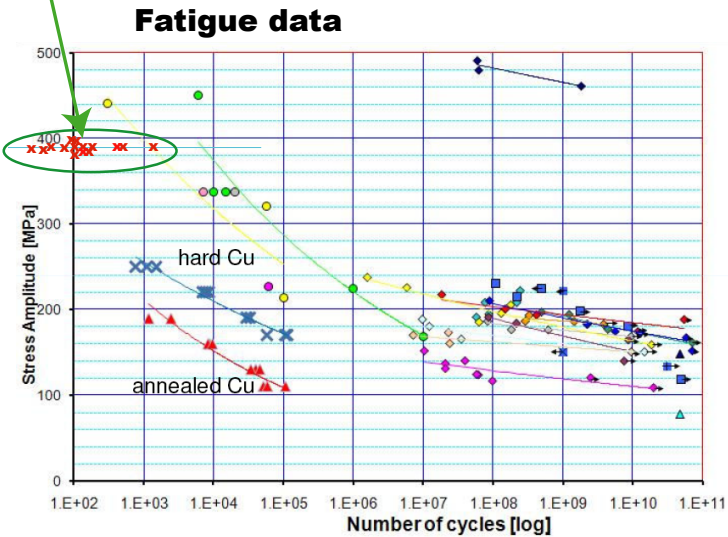
$$E_{\text{local}} \sim 10 \text{ GV/m @ } 805 \text{ MHz}$$



Coulomb explosions are compatible with fatigue
 Breakdown numbers from CERN small gap experiment.



$$P = E^2 \epsilon_0 / 2 = 386 \text{ MPa}$$



Unipolar arc properties

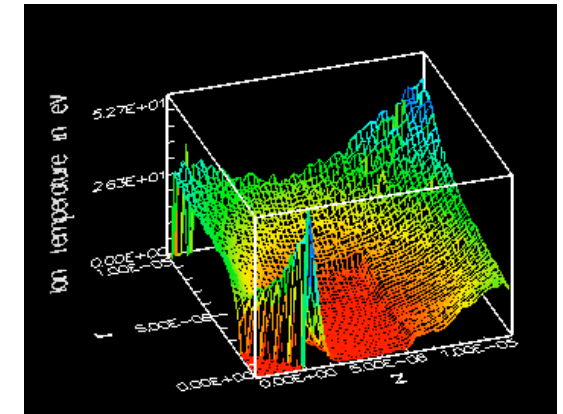
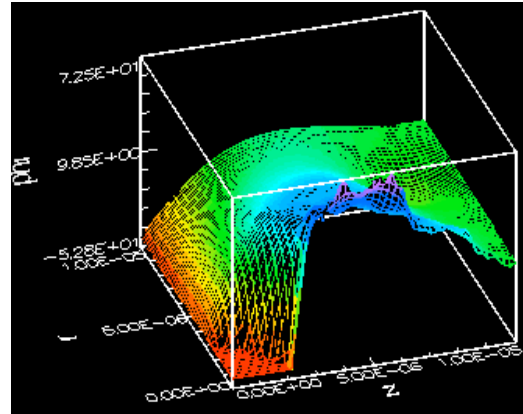
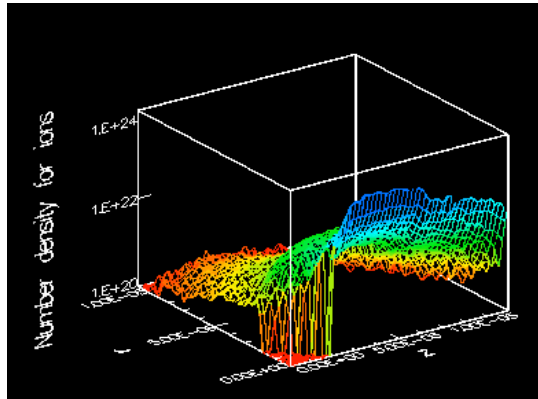
Dimensions matter

OOPIC Simulations (at ~ 6 ns), $r_{\max} = Z_{\max} = 10$ microns

Ion Density,

Phi,

T_I



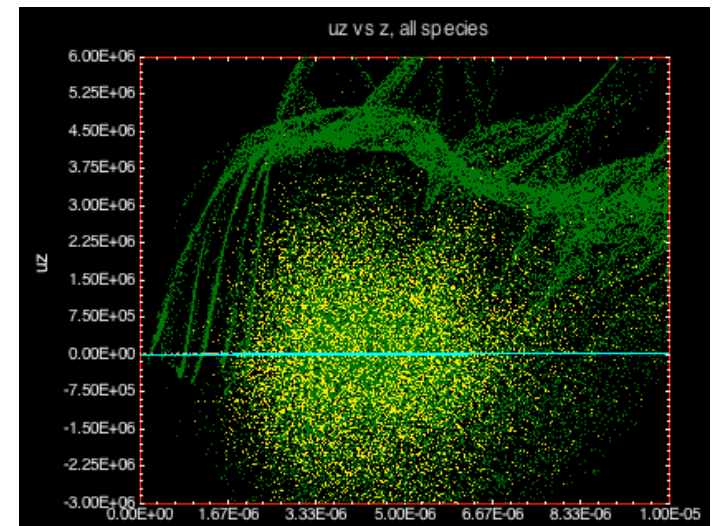
Arc dimensions a few microns.

Primary electron current

Space charge limit can be seen in v_z vs z

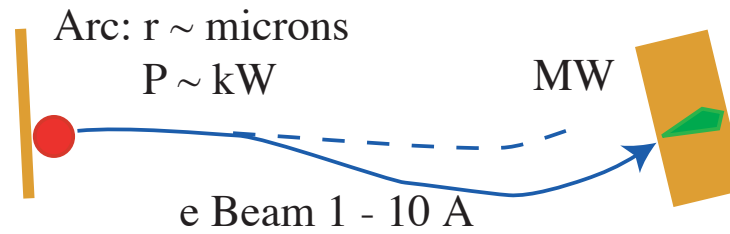
Plasma functions as a virtual cathode

Collision length remains constant $\sim 10 \mu$

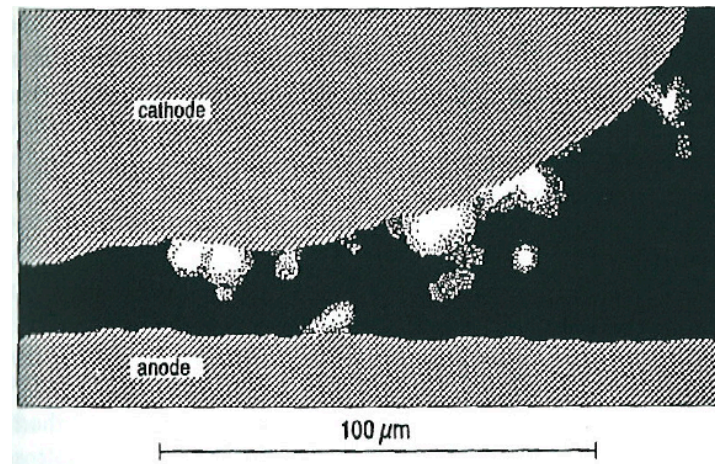


These (few micron) dimensions are consistent with a lot of data.

Our picture, from OOPIC

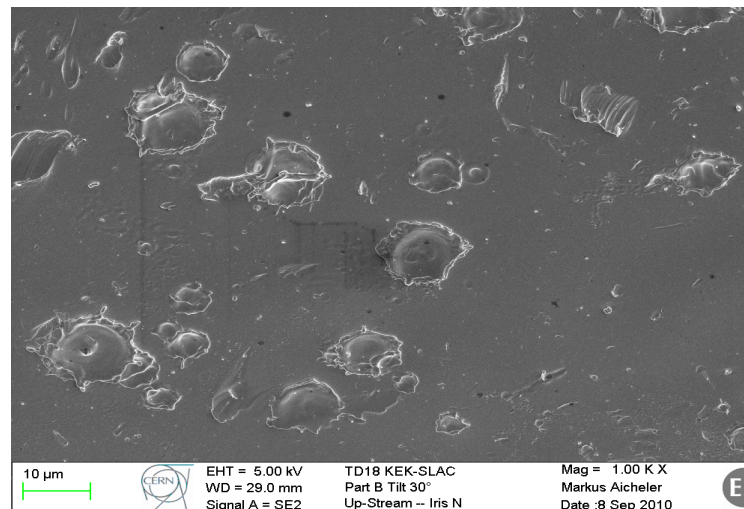


Boxman, Martin, Sanders,
Ch 3, Fig 22



CERN and SLAC
(Aicheler CERN '10)

.. etc.



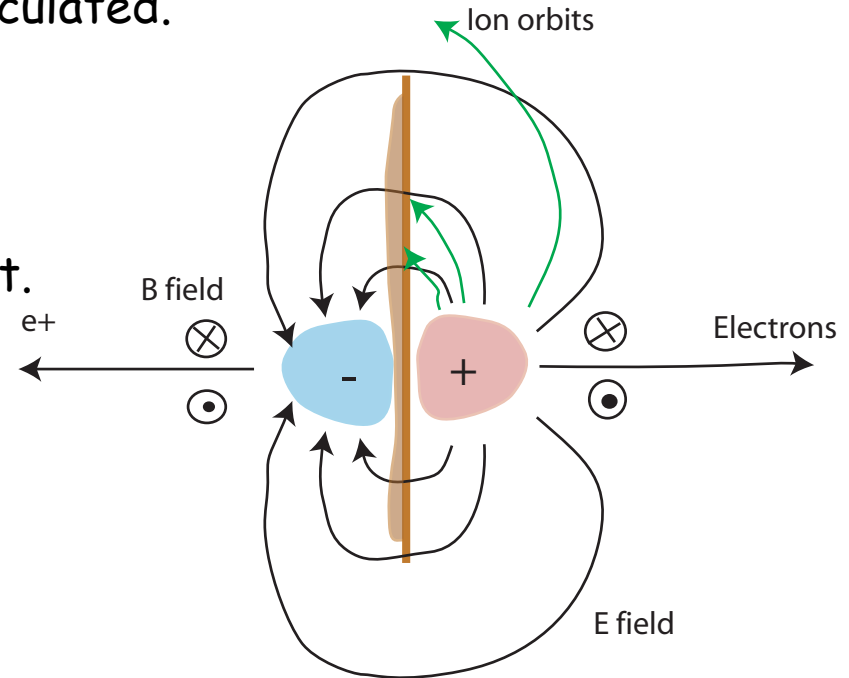
Electric field Distribution can be easily calculated.

Image charge gives boundary conditions

Ion motion determined by ϕ and $E(r,z)$

Radial dist of E, inc ion angle

Exact dimensions of plasma are important.



Time development:

n vs. time for arc

Trapped plasma develops

$E_{\text{average}} = \phi / \lambda_D$ (increases)

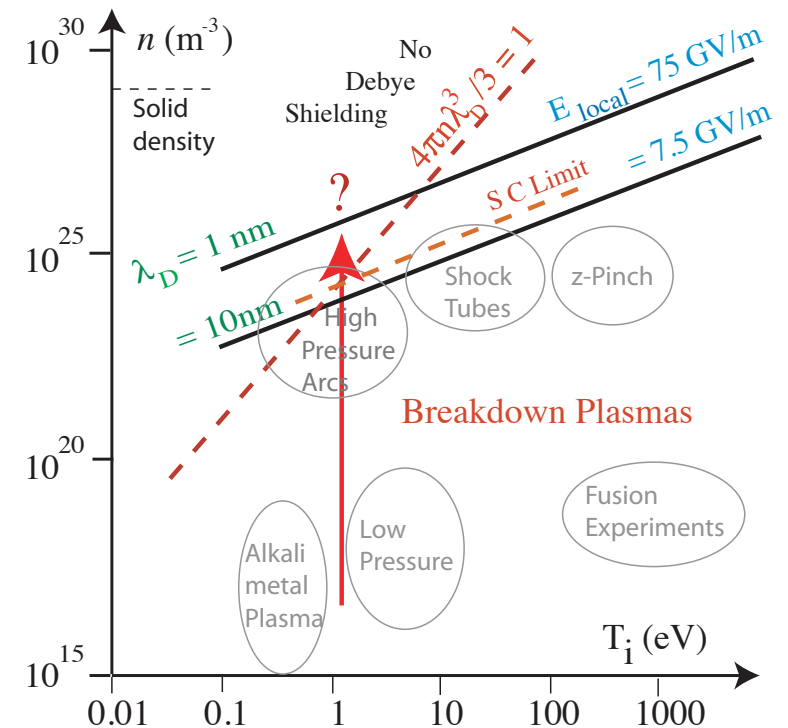
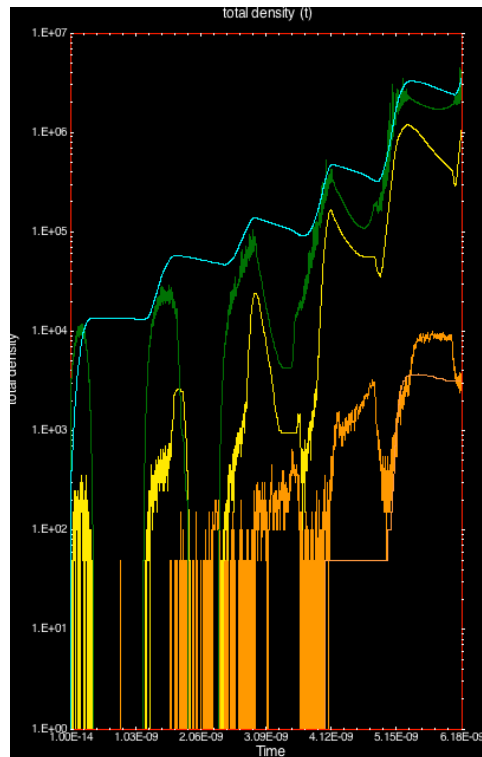
Non-Debye plasma ? ?

Increase in E drives

arc evolution

currents

density increase



This plasma is a NON "Local Thermal Equilibrium" (LTE) plasma

Ion & Electron Temperatures are very different

Ions are essentially thermal ($T \sim 0$), but stream in the electric field

Electrons stream through the plasma, but some are trapped.

Liquid surface deformation

Surface tension flattens surface

Electric tensile force pulls on surface, may be inhomogeneous

Plasma pressure pushes on surface

Spinodal decomposition causes ripples, can measure plasma surface properties

Non-Debye plasma properties are not understood

Plasma pressure is significant, $p = nkT$,

Generates particulates: Chapt 6 of Anders "Cathodic arcs"

Surface heating comes primarily from ion current

Plasma pressure forms small craters.

Power balance

Ion, electron fluxes change rapidly

Radiation flux goes like n^2

Surface heating is large and localized

Distribution of energy can be calculated.

B Field effects electrons and should change gradient limits.

Arc behavior is a function of B .

Focusing of shorting currents is evidently not harmful

Best gradient obtained with High collinear B field

Larmor focusing of electrons, $r_L = 0.3 [\mu] W_{eV}^{1/2}$

If $E \parallel B$, arc is more compact, more damaging,

If $E \perp B$, or $B = 0$ arc is more spread out

This mechanism should explain the B field dependence in muon accelerator cavities.

High local fields (megagauss) are seen

Ion radius much larger

less pressure

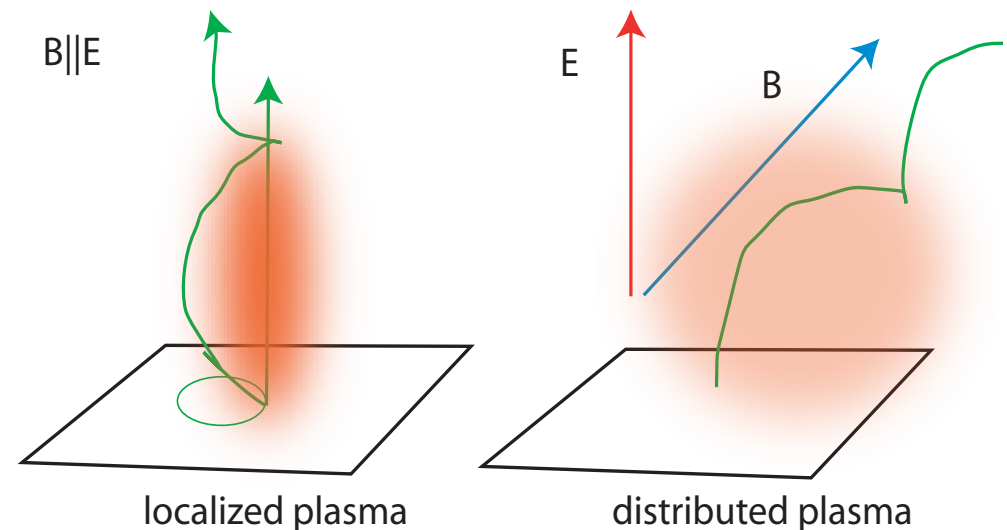
fewer particles, etc.

B fields are tough simulation problem

Requires VORPAL (or LSP)

Evidently parallel processing

We are trying to organize something.



Self-sputtering fuels the arc and can make it self-sustaining.

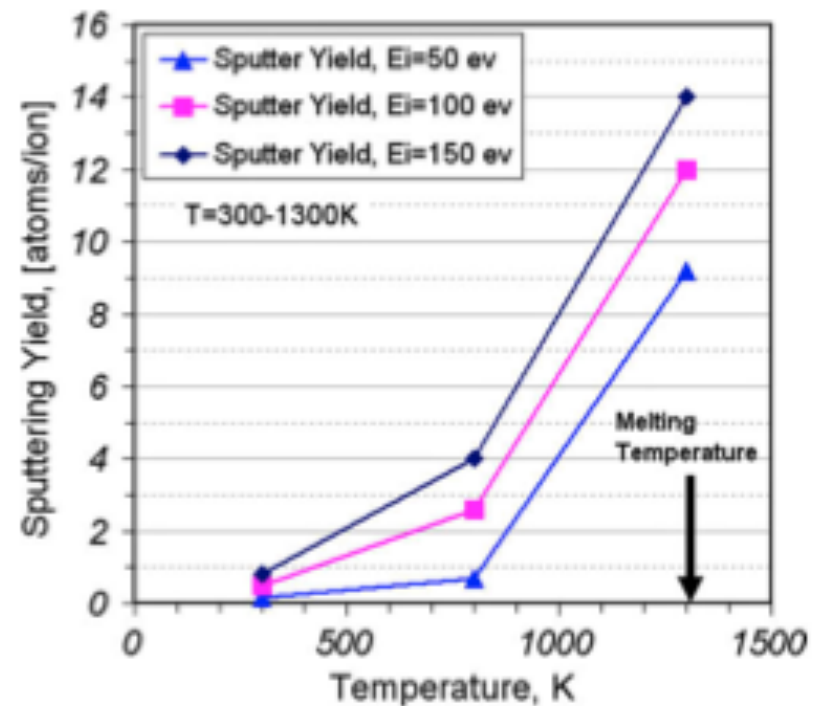
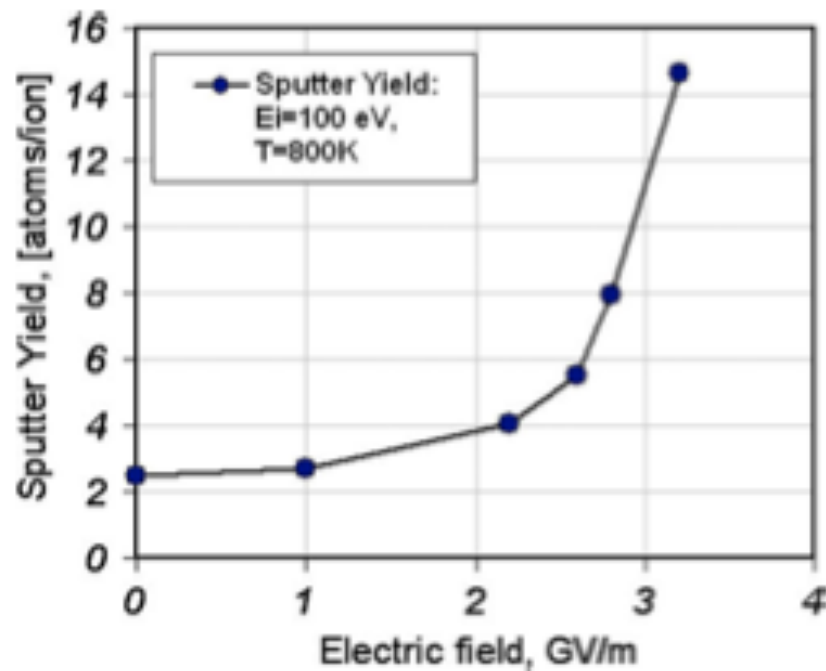
High E and T increase self-sputtering above 10 at low energies.

Grain orientation also seems to matter

We have calculations, are starting experiments

Electron sputtering also produces shorting currents.

Arcs must be self-sustaining in many non-accelerator environments.



We believe a unipolar arc is essentially a transient -
This is a little different from Schwirzke's picture.

Arcs seem to be inherently unstable.

Many candidates for a termination mechanism

Emitter melting

non-Debye hiccup

n_i/n_e imbalance

radiation cooling, . . .

Secondary sources appear nearby

Close: Ripples due to ion motion

Far: splashes from liquid particulates

Either way, fractal motion results

Space potential / sheath potential

Plasma is polarized because ions move more slowly than electrons

This would happen in free space or near a metallic boundary

Near a boundary the potential is essentially a normal sheath, even if nonthermal.

Ghost arcs (CERN) should be relics or precursors of Unipolar arcs.

Enhancement factors seem to be a source of confusion.

Fitting historical field emission data seems to give a wide range for β and A .

$$2 < \beta < 1000$$

$$1 \text{ nm}^2 < A < \text{many } \mu^2$$

(This wide range is not seen in cavities however.)

These values are not compatible with a whisker model of enhancement factors.

The validity of Fowler-Nordheim (and quantum mechanics) has been questioned.

We have a simple solution:

Emitters are small, ($A \sim 1 \text{ nm}^2$) with natural β s around 100.

They are formed at crack junctions and spattered particulates.

If they sit on other structures, their β s can be much larger.

If there are lots of them, the combined A will be much larger.

Surfaces are rough, so lots of structures to sit on and lots of spatters/cracks.

Fowler-Nordheim should be OK, as is Quantum Mechanics.

Damage

Damage mechanisms

Cracks

Particulates, splashes

Oblique ion fluxes and ripples

Damage from shorting currents

Erosion

The spectrum of $n(\beta)$ is known, in many environments.

Nobody sees whiskers, so breakdown sites must be blunt

Thermal properties of blunt corners

Time const for cooling ~ 10 fs

Heated volume is small $\sim 1 \text{ nm}^3$.

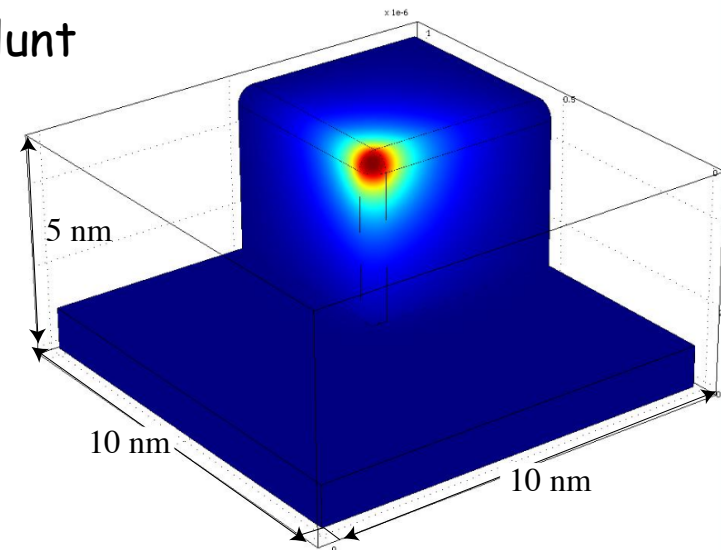
Large heat sink $\sim 1 \mu^3$ (@ ns)

VERY hard to heat

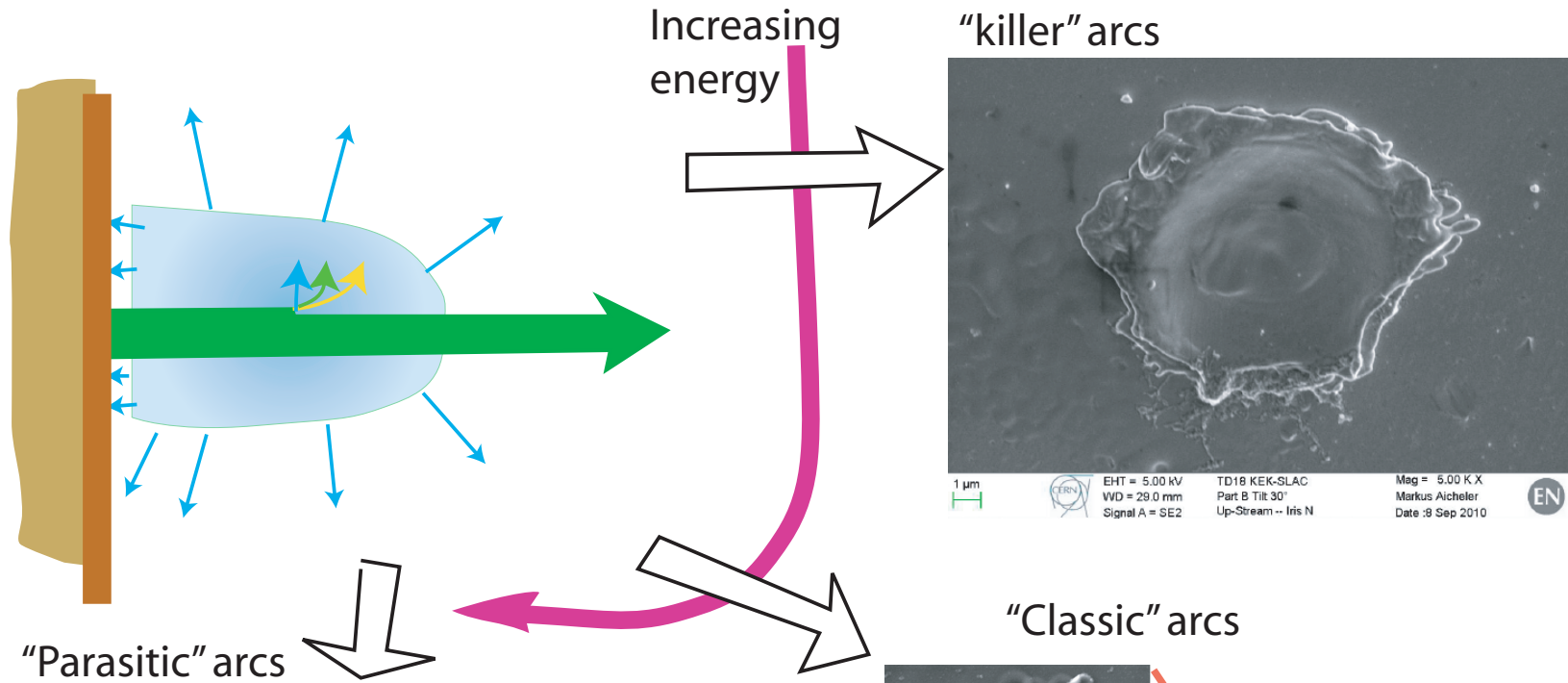
Two types of arcs

Killer - arc current shorts potential, removes power source.

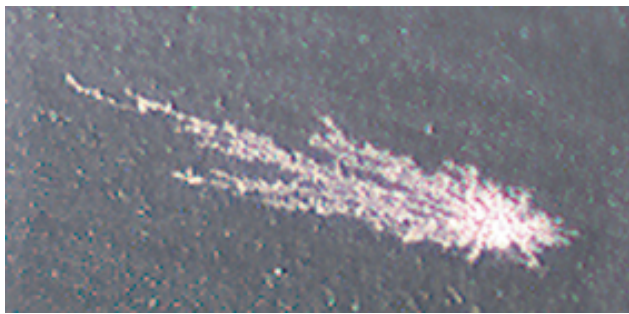
Parasitic - arc removes energy from plasma slowly



Effects of increasing Arc energy

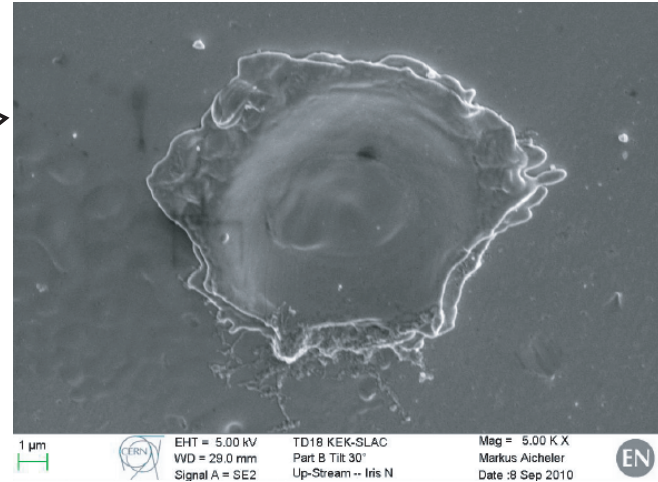


"Parasitic" arcs

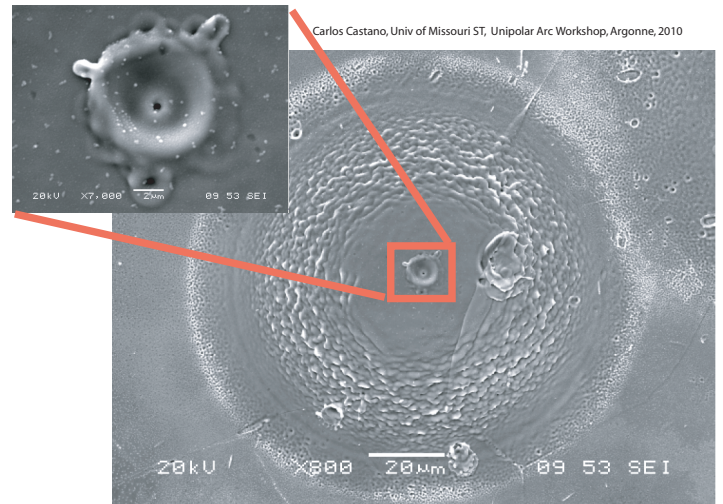


These arcs are ~1 cm long
201 rf coupler

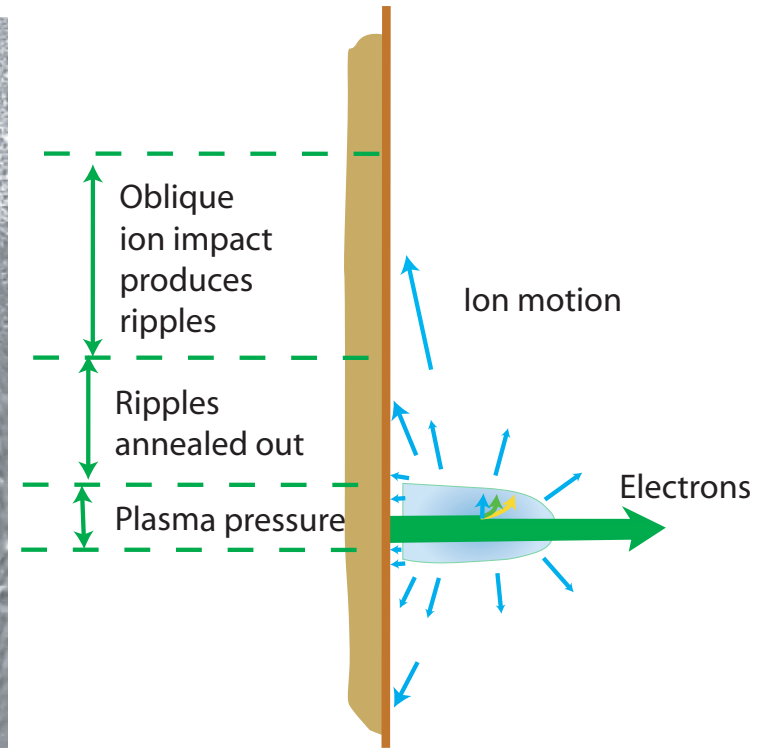
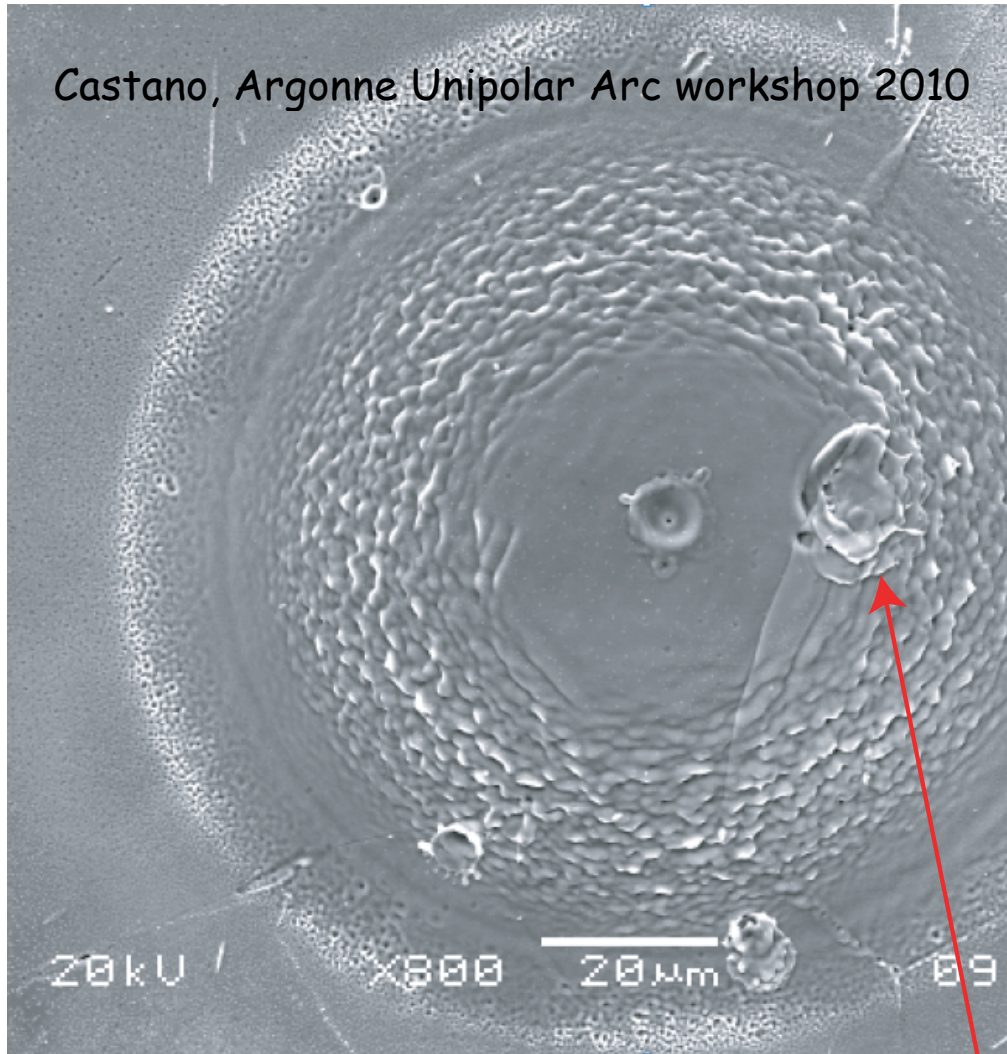
"killer" arcs



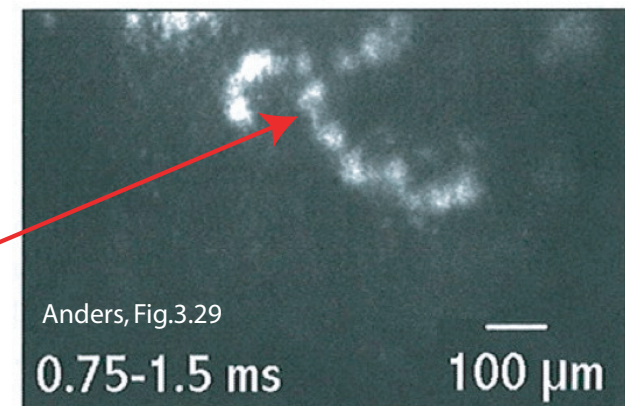
"Classic" arcs



Arc damage mechanisms



Secondary breakdown sites ~25 mm from primary site consistent with fractal arc motion

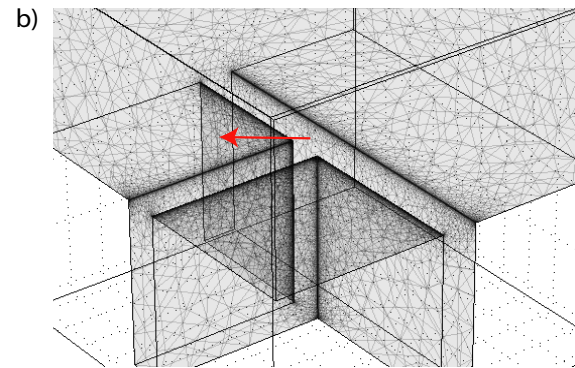
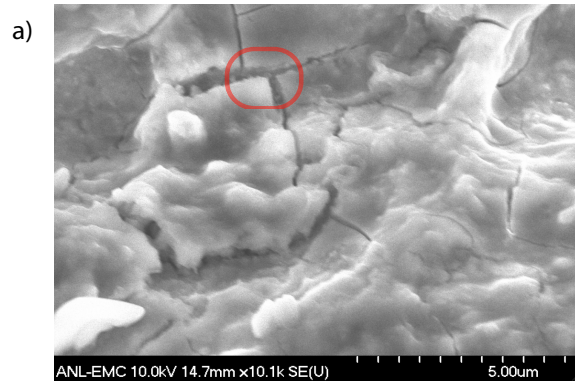


Cracks

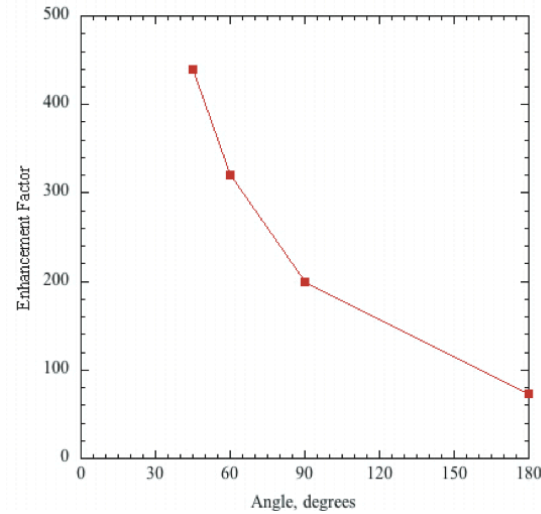
$n(\beta)$ from angle dependence could be derived if needed.

Sharp crack angles have high β s but are uncommon.

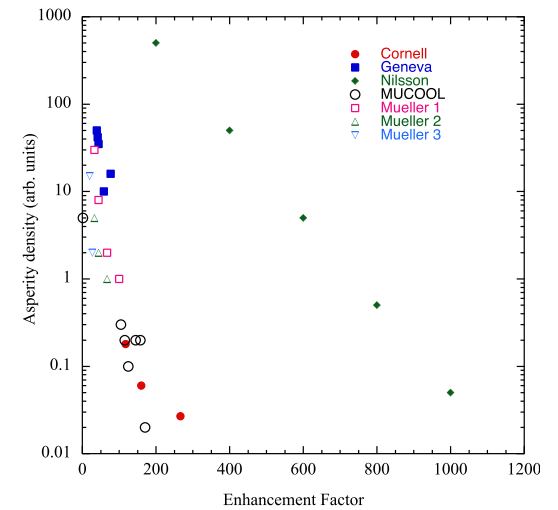
SEM showing cracks



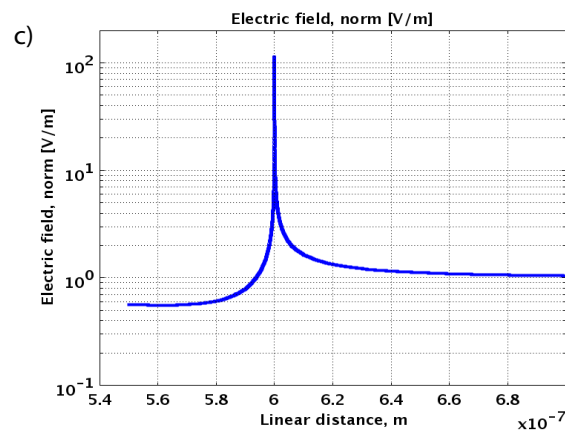
β (crack angle)



$n(\beta)$



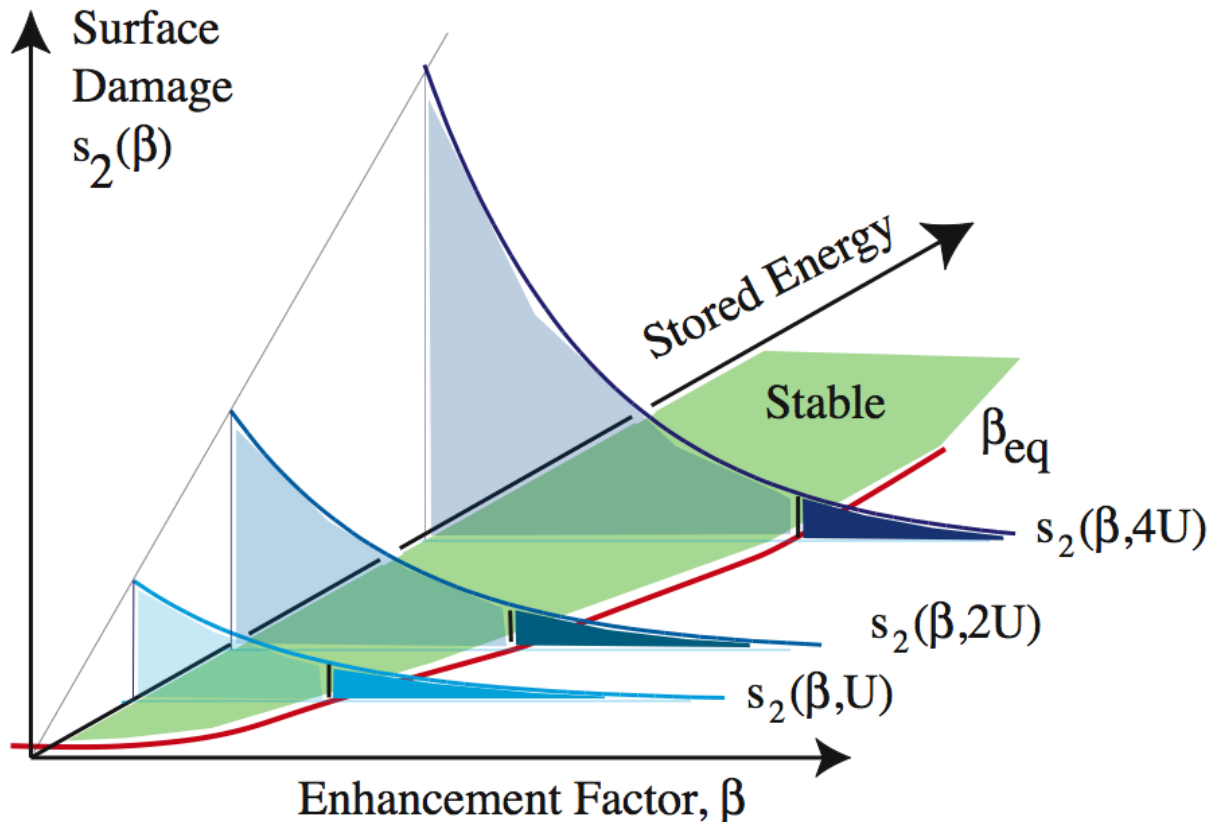
Particulates can also explain this effect.



Gradient limits

Gradient limits represents an equilibrium obtained from two effects:

- 1) More energy in the arc produces more damage. (They are proportional.)
- 2) Because of the spectrum of enhancements drops off exponentially, damage rises logarithmically with arc energy.



We published a paper in '05 describing how damage causes gradient limits, with a number of predictions and examples.

Conditioning

We can calculate the Kilpatric limit

Gap dependence: there is none

BDR(E) can be due to a number of causes

Ohmic

Electromigration

Fatigue

Are all indistinguishable, give $BDR \sim E^{30}$.

Gas pressure scaling: there is none

Temperature dependence: there is none

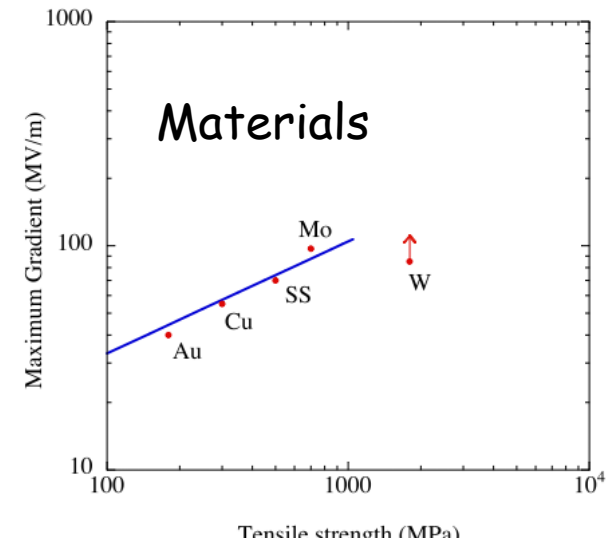
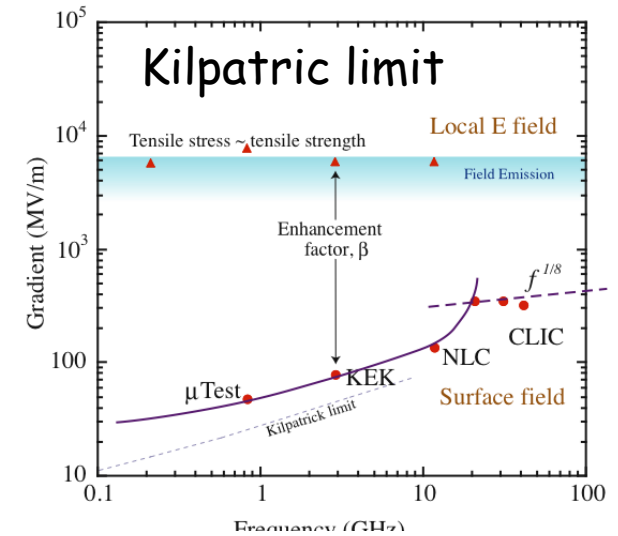
Pulse length dependence: depends on energy.

Material dependence: depends on tensile str.

noble metals only

Correlated breakdown rate: complex problem

And more . . .



Active Effort

Thermodynamics of breakdown sites

B field effects in vacuum arcing

Sputtering and self-sputtering dependence on many parameters

Differential erosion depends on grain orientation

Surface electric fields in arcs.

How does background plasma affect gradient limits?

Non-Debye, non-LTE plasma interactions with materials.

Conclusions

We feel we can calculate all aspects of arc properties in a realistic way.

We have produced many predictions

We are working on many aspects of the problem

Some problems are very difficult (3D PIC /w B fields)

We feel that we can predict any quantity with useful precision.

We believe our methods have very general applicability.

Vacuum arcs and gradient limits seem to be considered an insolvable problem.

Little interest, little support

History (SSC, NLC, ILC) implies we need to know how things work.

There seem to be different ideas, we should have vigorous debate.

If pulse heating is real, Vacuum arcs must be the damage mechanism