POTENTIAL-WELL DISTORTION IN BARRIER RF

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• Head-tail asymmetry was observed at SLAC SLC damping ring, but has never been observed in proton machines.



Figure 1: Potential-well distortion of bunch shape for various beam intensities for the SLAC SLC damping ring.

• Haissinski equation

$$\begin{split} \rho(\tau) &= \rho_0 \left[-\frac{\tau^2}{2\sigma_\tau^2} - \alpha_R N \int_0^\tau \rho(\tau') d\tau' \right] \\ &\uparrow \\ U_{\rm rf}(\tau) & \stackrel{\uparrow}{\sim} \rho_0 \tau \end{split}$$

gives shift of peak:

τ	$\underline{\alpha_R N}$	$\underline{e^2\beta^2 ENR_s}$
$\overline{\sigma_{ au}}$	$-\sqrt{2\pi}$	$-\sqrt{2\pi}\eta T_0\sigma_E^2$



- For proton beam, $\alpha_R N/\sqrt{2\pi} \ll 1$, the rf potential.
- However, at the Fermilab Recycler Ring, the rf potential $U_{\rm rf}(\tau) = 0$, and head-tail asymmetry has been observed.

THE RECYCLER RF

• Barrier RF cavities are used to store \bar{p} .



- Merits:
 - 1. Beam spreads out with lower space-charge force.
 - 2. Can merge two batches together easily.
 - 3. Can compress by moving a barrier slowly.
 - 4. Can move batch from one location to another.



- - Four 50 Ω ferrite-loaded rf cavity stations.
 - Amplifiers: 3.5 kW and broad-band (10 kHz to 100 MHz) capable of supplying 2 kV.
 - Rf waveform determined by amplitude and phase of each of the 1113 revolution harmonics
- If baseline is not zero, rf potential will be head-tail asymmetric.



- Nonzero baseline can come from rf error or impedance.
- Here, we are talking about ~ 10 V out of 2 kV (0.5%).



Figure 2: Top: Beam intensity $\sim 1 \times 10^{11}$. Bottom: Nonlinear region of Recycler high-level RF. Require flatness from 90 kHz to 1 MHz 0.26 dB in amplitude and 1.8° in phase.



Figure 3: After linearization transfer function implemented.

• With the linearization compensation properly adjusted, the slant reappears: at higher beam intensity like $N\gtrsim 2\times 10^{11}$

with barriers separation $T_2 = 1.6 \ \mu s$.



- Now beam intensity shifts towards tail → appears to come from resistive impedance.
- Barrier cavities have total $\mathcal{R}e Z_0^{\parallel} = 200 \Omega$, visible to beam up to ~ 45 MHz (harmonics n = 1 to 500).
- $Z_0^{\parallel} = 7.6 \ \Omega$ at $n = 1 \rightarrow$ can be neglected.
- Beam loading voltage

$$V_b = I_{\text{local}} \approx eNZ_0^{\parallel} \rho(\tau) \approx \frac{eNZ_0^{\parallel}}{T_2} = 12.9 \text{ V}.$$

• The slant can be compensated by adding ~ 8.82 V to the region between the barriers.



- Assuming ΔE has Gaussian distribution, can solve Haissinski equation to obtain $\rho(\tau)$.
- Gives $\pm 17\%$ asymmetry, larger than measurement.
- To correct for asymmetry, needs 12.35 V, which is $\sim 40\%$ larger than actual compensation.



UNSOLVED PROBLEMS Problem 1

• The voltage compensation is smaller than actually used.



- Possible reasons:
 - 1. Linearization has been <u>over compensated</u>. It has been used to make slant zero at 1.1×10^{11} .
 - 2. Shunt impedance of cavities may be less than 200 Ω . However, loaded shunt impedance may not be much less because cavities are broad-band.

Theoretical determination of V_{comp} is very general. It depends only on

$$\psi(\Delta E,\tau)=\psi(H)$$
 .

• We tried a more realistic model with wake

$$W_0(t) = \sqrt{\frac{2}{\pi}} R_s \sigma_\omega e^{-\sigma_\omega^2 t^2/2} ,$$

so that the resistive impedance

$$\operatorname{Re} Z_0^{\parallel}(\omega) = R_s e^{-\omega^2/2\sigma_\omega^2}$$

rolls off around $\sigma_{\omega}/2\pi \approx 45$ MHz.

Negligible change in result because the beam does not have many high-frequency components.

Problem 2

The predicted slant is not linear.



• Haissinski equation is

$$\rho(\tau) = \rho_0 \exp\left[\frac{\beta^2 E f_0 e V_0 T_1}{-\eta \sigma_E^2} U_{\rm rf}(\tau) - \frac{\beta^2 e^2 N E R_s}{-\eta T_0 \sigma_E^2} \int_0^\tau \rho(\tau') d\tau'\right]$$

In the region where potential $U_{\rm rf}(\tau) = 0$,

$$\rho' = -\alpha_R N \rho^2$$

Or

$$\rho(\tau) = \frac{\rho_0}{1 + \alpha_R N \rho_0 \tau} \; .$$

• Elliptical-like distributions:

$$\psi(\tau, \Delta E) = A \left[\widehat{\Delta E}_0^2 - \Delta E^2 \right]^n \longrightarrow A \left[\widehat{\Delta E}^2(\tau) - \Delta E^2 \right]^n$$

where

$$\widehat{\Delta E}^2(\tau) = \widehat{\Delta E}_0^2 \left[1 + bU_{\rm rf}(\tau) - a \int_0^\tau \rho(\tau') d\tau' \right]$$

Then

$$\rho(\tau) = \int \psi d\Delta E = 2\gamma_n A \left[\widehat{\Delta E}^2(\tau)\right]^{n+1/2}$$

$$\rho' = -\left(n + \frac{1}{2}\right) a\rho_0^{\frac{2}{2n+1}} \rho^{\frac{4n}{2n+1}} ,$$

$$\rho(\tau) = \rho_0 \left[1 + \left(n - \frac{1}{2} \right) a \rho_0 \tau \right]^{-\frac{2n+1}{2n-1}}$$



PROBLEM 3

- Elliptical-like distribution with $n = \frac{1}{2}$ does not fit measured energy spread.
- 1.75 GHz Schottky signals in dBm, logarithmic scale.
- $\sigma_{\underline{E}} = 2.3 \text{ MeV}$ (3.9 kHz).
- $\Delta E_{90\%} = 7.8 \text{ MeV}$ (13 kHz).



• Convert to linear scale.



- Elliptical distribution does not fit.
- Gaussian distribution fits better.

CONCLUSION

- We have reported head-tail asymmetry in a proton beam due to potential-well distortion of the special barrier rf.
- This slant $\rho(\tau)$ may overflow the barrier bucket when beam intensity is too high.
- When barrier bucket is too full, instabilities will occur due to resonances driven by rf jitters.
- It will also affect other application of barrier rf, like doubling proton intensity in Fermilab Main Injector.
- The required compensating beam-loading voltages are smaller than theory predicted.
 When the Recycler is turned on again a month later, we must make sure linearization is properly compensated but not over-compensated.
- Elliptical dist. $(n \leq 0.5)$ fits $\rho(\tau)$ better but not ΔE dist. Gaussian dist. fits ΔE dist. better but not $\rho(\tau)$.
- We should investigate some other phase-space distributions like <u>cosine square</u>, etc.