



NLO QCD corrections to $pp/p\bar{p} \rightarrow t\bar{t} + \text{jet} + X$

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1 Introduction

Why is $pp/p\bar{p} \rightarrow t\bar{t} + \text{jet} + X$ interesting ?

- Top-quark properties and dynamics still not precisely known

↔ additional observables welcome, e.g.

FB charge asymmetry of heavy quarks in $q\bar{q} \rightarrow t\bar{t}(+\text{jets})$

origin: interference of C -odd / even parts initial and final states

$q\bar{q} \rightarrow t\bar{t}$: asymmetry appears first in NLO Kühn, Rodrigo '98

↔ $A_{\text{FB}}(\text{NLO})$ not feasible in near future (requires σ_{NNLO})

$q\bar{q} \rightarrow t\bar{t}+g$: asymmetry is LO effect Halzen, Hoyer, Kim '87

↔ $A_{\text{FB}}(\text{NLO})$ feasible and in progress

Note: asymmetry is larger for $t\bar{t}+\text{jet}$ than for $t\bar{t}$ inclusive process

Situation at Tevatron: asymmetry is measurable ! Bowen, S.D.Ellis, Rainwater '05

LHC: asymmetry requires preferred direction from partonic boost effects

- Important background process for Higgs search at the LHC via

$pp(\text{WW} \rightarrow \text{H}) \rightarrow \text{H} + 2\text{jets}$

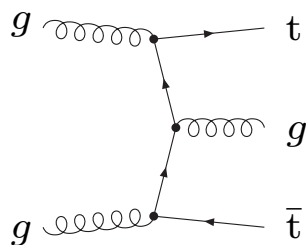
$pp \rightarrow t\bar{t}\text{H} + X$

2 Calculation of NLO corrections

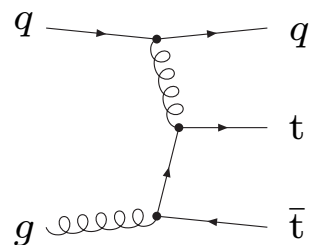
2.1 Lowest-order prediction

Some LO diagrams:

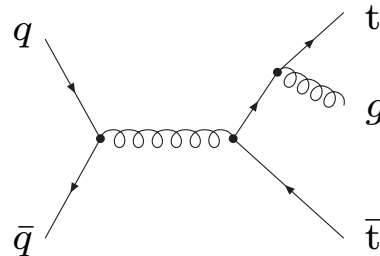
$gg \rightarrow t\bar{t}g$:
(16 diagrams)



$qg \rightarrow t\bar{t}q$:
(5 diagrams)



$q\bar{q} \rightarrow t\bar{t}g$:
(5 diagrams)



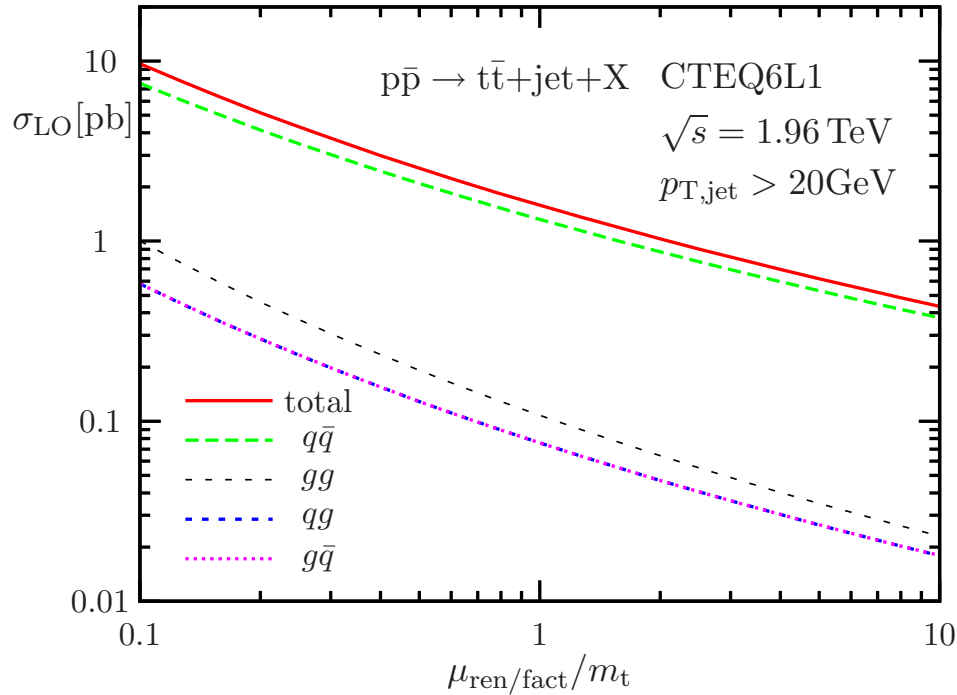
etc.

Features of the LO cross section:

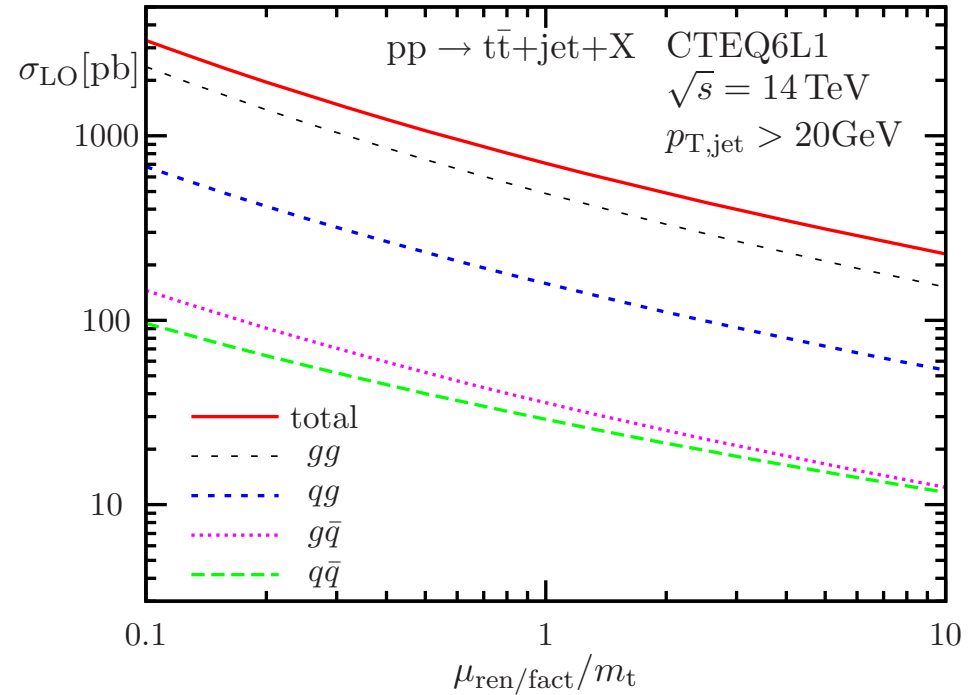
- IR safety requires at least lower cut on $p_{T,\text{jet}}$
 \hookrightarrow apply jet algorithm for NLO cross-section before cut on $p_{T,\text{jet}}$
- LO hadron cross section $\propto \alpha_s^3$
 \hookrightarrow strong dependence on renormalization and factorization scales

Scale dependence of LO cross sections:

Tevatron



LHC



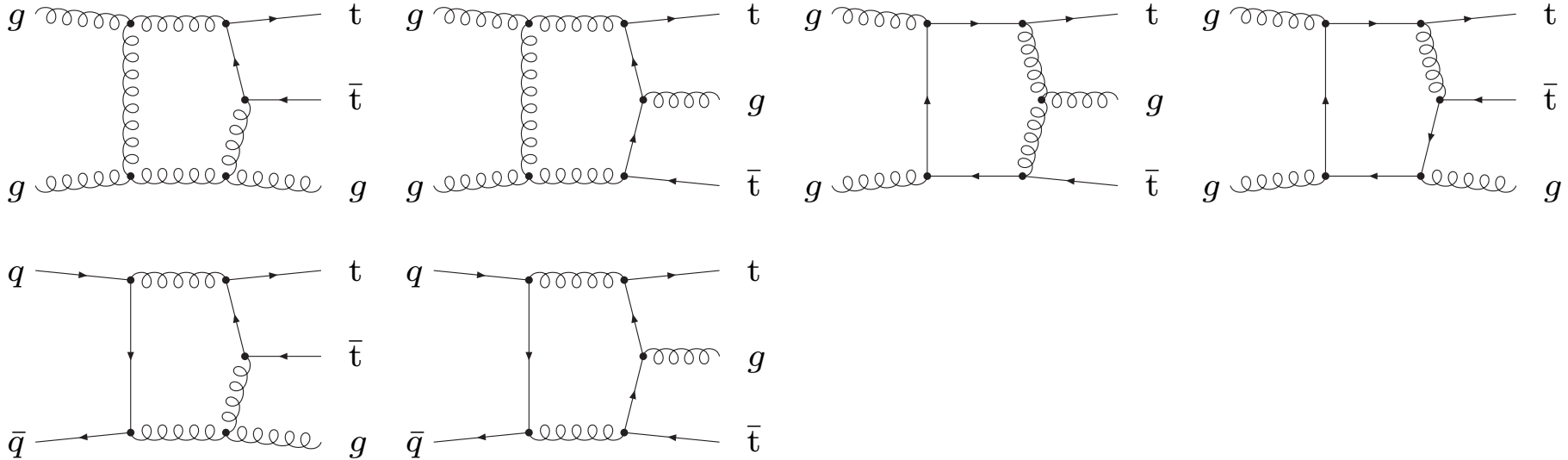
In LO: light final-state parton \equiv jet

\hookrightarrow no dependence on jet algorithm

2.2 Virtual corrections

1-loop diagrams $\sim 350(100)$ for $gg(q\bar{q}) \rightarrow t\bar{t}g$

Most complicated 1-loop diagrams—pentagons of the types:



Algebraic reduction of amplitudes to standard form, e.g.

$$\mathcal{M}_{g_a g_b \rightarrow t_i \bar{t}_j g_c} = \sum_{k=1}^{10} \sum_{l=1}^{144} \underbrace{F_{kl}(\{p_i \cdot p_j\})}_{\substack{\text{invariant functions} \\ \text{containing loop integrals}}} \underbrace{C_k}_{\text{standard colour structures}} \underbrace{\hat{M}_l(\{p_i\})}_{\substack{\text{standard spinor structures} \\ \hat{M}_1 = (\bar{v}_t u_{\bar{t}})(\varepsilon_a \varepsilon_b)(k_t \varepsilon_c^*), \text{ etc.}}}$$

$C_1 = (T^{c_a} T^{c_b} T^{c_c})_{ij}, \text{ etc.}$

Two independent strategies for evaluation of loop integrals

- Calculation analogous to NLO QCD calculation for $pp \rightarrow t\bar{t}H$ Beenakker, S.D., Krämer, Plümper, Spira, Zerwas '01,'02
 - ◇ diagrams generated with FEYNARTS 1.0 Küblbeck, Böhm, Denner '90 and reduced with in-house MATHEMATICA routines \rightarrow FORTRAN
 - ◇ analytical extraction of soft / collinear singularities Beenakker et al. '02; S.D. '03
 - ◇ reduction of 5-point to 4-point integrals according to Denner, S.D. '02
 \hookrightarrow no (leading) inverse Gram det's \rightarrow sufficient numerical stability
 - ◇ **outlook:** process will be used as further test ground for more sophisticated tensor reduction methods (seminumerical and/or expansion techniques) Denner, S.D. '05
used at NLO EW for $e^+e^- \rightarrow 4f$ Denner et al. '05
- Alternative calculation
 - ◇ diagrams generated with QGRAF (Nogueira '93) and reduced with FORM \rightarrow C++
 - ◇ reduction of 5-point to 4-point integrals according to Giele, Glover '04
 \hookrightarrow no (leading) inverse Gram det's \rightarrow still testing numerical stability
 - ◇ **outlook:** further numerical stabilization via expansion method suggested by Giele, Glover, Zanderighi '04

Strategy for extracting or translating IR (soft / collinear) singularities:

Idea: convert integrals $I^{(D)}$ in $D=4-2\epsilon$ dim.

→ 4-dim. integrals $I^{(\lambda)}$ with mass regulator λ

Procedure: consider finite and reg.-scheme-independent difference

$$\left[I^{(D)} - I_{\text{sing}}^{(D)} \right] \Big|_{D \rightarrow 4} = \left[I^{(\lambda)} - I_{\text{sing}}^{(\lambda)} \right] \Big|_{\lambda \rightarrow 0}$$

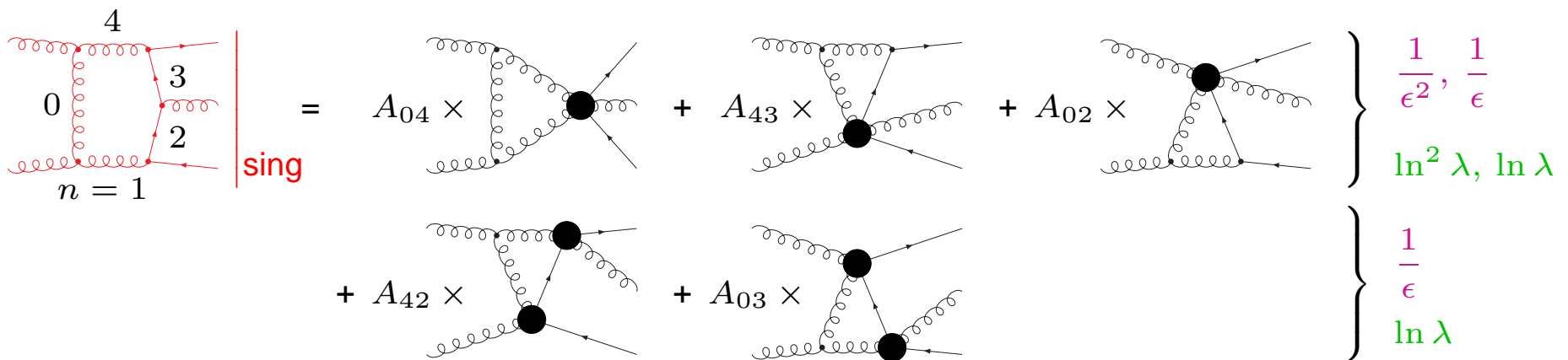
$$\Rightarrow I^{(D)} = I_{\text{sing}}^{(D)} + \left[I^{(\lambda)} - I_{\text{sing}}^{(\lambda)} \right] \Big|_{\lambda \rightarrow 0} + \mathcal{O}(\epsilon)$$

Note: mass-singular part can be universally constructed from 3-point integrals

↪ general result known explicitly S.D. '03

Beenakker et al. '01

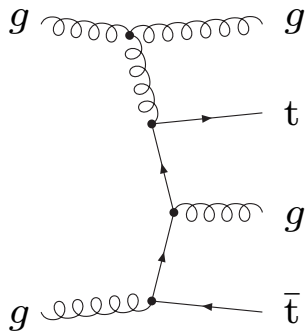
An example:



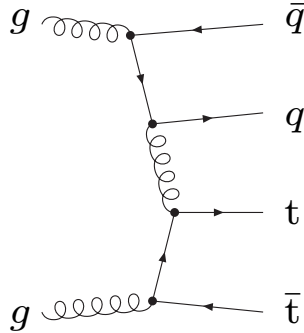
2.3 Real corrections

Some diagrams with 1-parton emission

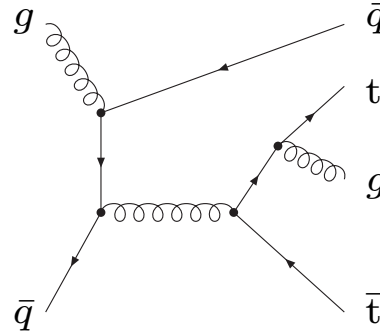
$gg \rightarrow t\bar{t}gg$:
(123 diagrams)



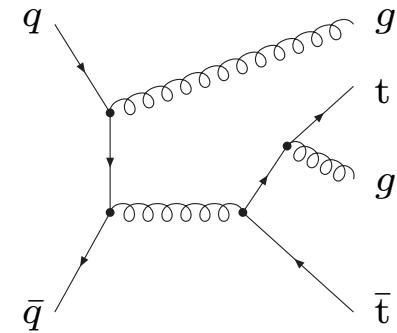
$gg \rightarrow t\bar{t}q\bar{q}$:
(36 diagrams)



$g\bar{q} \rightarrow t\bar{t}q\bar{q}$:
(36 diagrams)



$q\bar{q} \rightarrow t\bar{t}gg$:
(36 diagrams)



etc.

Features of the calculation

- evaluation of helicity amplitudes via
 - ◇ conventional (4-dimensional) spinor techniques
automated a la **Weinzierl '05**
 - ◇ **Berends/Giele** recurrence relations
- ↔ $|\mathcal{M}|^2$ checked against **MADGRAPH** **Stelzer, Long '94**
- extraction and integration of soft / collinear singularities via
dipole subtraction formalism **Catani, Seymour '96; S.D. '99; Phaf, Weinzierl '01**
Catani, S.D., Seymour, Trocsanyi '02

Dipole subtraction formalism

→ process-independent treatment of singularities in real NLO corrections

worked out for

- QCD with massless partons (Catani, Seymour '96)
 - γ radiation off massive fermions (S.D. '99)
- } QCD with massive partons
Phaf, Weinzierl '01
Catani, S.D., Seymour, Trócsányi '02

basic idea: NLO correction to process with m partons

$$\sigma^{\text{NLO}} = \underbrace{\int_{m+1} \left[d\sigma^{\text{real}} - d\sigma^{\text{sub}} \right]}_{\text{finite}} + \underbrace{\int_m \left[d\sigma^{\text{virtual}} + d\bar{\sigma}_1^{\text{sub}} \right]}_{\text{finite}} + \int_0^1 dx \underbrace{\int_m \left[d\sigma^{\text{fact}}(x) + \left(d\bar{\sigma}^{\text{sub}}(x) \right)_+ \right]}_{\text{finite}}$$

conditions on $d\sigma^{\text{sub}}$:

- sum rule: $-\int_{m+1} d\sigma^{\text{sub}} + \int_m d\bar{\sigma}_1^{\text{sub}} + \int_0^1 dx \int_m \left(d\bar{\sigma}^{\text{sub}}(x) \right)_+ = 0$
- asymptotics: $\sigma^{\text{sub}} \sim \sigma^{\text{real}}$ in all collinear/IR regions

2.4 Checks and status of the calculation

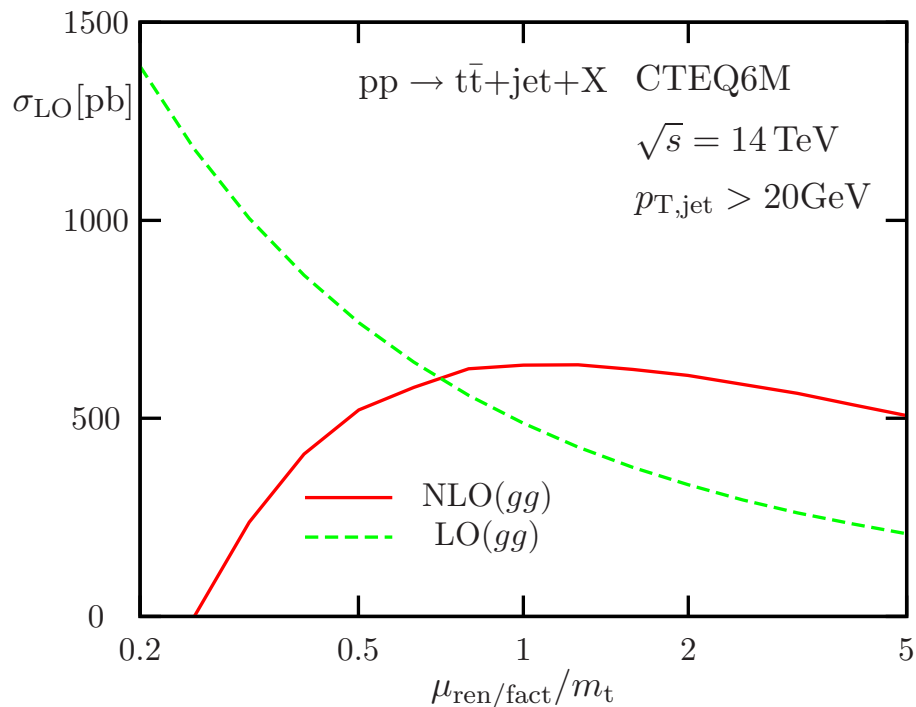
All ingredients available, but not all cross-checked:

- gg -fusion channel almost completely checked:
 - ◇ UV structure of virtual correction
 - ◇ soft and collinear structure in real and virtual corrections
 - ◇ crossing symmetries
 - ◇ all ingredients (but some parts of real corrections)
confirmed in second, independent calculation
- $q\bar{q}, gq, g\bar{q}$ channels:
 - ◇ UV structure of virtual correction
 - ◇ soft and collinear structure in real and virtual corrections
 - ◇ crossing symmetries
 - ◇ cross-checks by second calculation still missing

3 Numerical results

LO versus NLO contributions of $gg \rightarrow t\bar{t}g(g)$:

LHC



Preliminary !

Jet definition:

algorithm of S.D.Ellis, Soper '93

with $R = 1$

applied to jets other than t and \bar{t}

\hookrightarrow Scale dependence stabilizes at NLO

4 Conclusions

The process $pp/p\bar{p} \rightarrow t\bar{t} + \text{jet} + X$

- important background process for Higgs and other searches at the LHC
- interesting playground to investigate top-quark dynamics
 - ↪ measurement of FB charge asymmetry in $t\bar{t}(\text{+jets})$ already at Tevatron
- ↪ NLO prediction for $t\bar{t} + \text{jet}$ production desirable

$pp/p\bar{p} \rightarrow t\bar{t} + \text{jet} + X$ at NLO QCD

- calculation completed, but not yet fully cross-checked
- preliminary numerical results shown
- example is important test ground for NLO methods for many-particle processes
 - ↪ methods not yet exhausted,
more complicated applications (2 → 4) feasible !

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“It’s promising for the future. The important matches are coming.” (Zinedine Zidane)