Simulations of pulsar winds: the story of binary system J0737 Anatoly Spitkovsky (Chandra Fellow, KIPAC)

Relativistic (extreme) astrophysics

- Pulsars + winds
- Plerions
- Extragalactic radio sources
- Superluminal expansion
- Black hole energy extraction
- Gamma ray bursts
- Magnetars / AXPs
- Galactic superluminal sources
- UHE CR











Magnetized rotators

Magnetized rotators are ubiquitous: pulsars, AGN, GRBs (?)

Rotation very efficient at long-term energy storage

Extraction of rotational energy is associated with relativistic outflows: pulsar winds, AGN jets, GRB jet flows.

Unipolar induction: $V \sim \Omega \Phi$; $P \sim V^2 / Z_0$

Crab Pulsar

B ~ 10¹² G, W ~ 200 rad s⁻¹, R ~ 10 km

V ~ 3 x 10¹⁶ V; I ~ 3 x 10¹⁴ A; P ~ 10³⁸erg/s

Magnetar

B ~ 10^{14} G; P ~ 10^{44} erg/s Massive Black Hole in AGN B ~ 10^{4} G; P ~ 10^{46} erg/s GRB B ~ 10^{16} G; P ~ 10^{49} erg/s

How is energy extracted, transported, collimated?

Best example -- pulsar winds

Pulsars and their winds







Best studied examples are pulsar winds in young supernova remnants. The wind shines in synchrotron light when collides with the nebula.

- Properties of pulsar winds:
 - Highly relativistic ($\gamma \sim 10^6$)
 - Kinetic energy dominated at the nebula (σ ~10⁻³)
 - Pole-equator asymmetry and collimation

How do they do this?

Pulsar spindown: wind energy loss





Corotation radius = "Light Cylinder" (v_{rot} =c) R_L = cP/2 π

Energy loss in the form of Poynting flux and particle flux:

$$\dot{E} = -I\Omega\dot{\Omega} = \int dA \left[\frac{c}{4\pi} (\mathbf{E} \times \mathbf{B})_{\mathbf{r}} + \rho(\gamma - 1)c^2 v_r\right] = cr^2 \frac{B_{\phi}^2}{4\pi} \left(1 + \frac{1}{\sigma}\right)$$
$$= cR_L^2 B_p(R_L) B_{\phi}(R_L) = c \frac{\mu^2}{R_L^4} = \frac{\mu^2 \Omega^4}{c^3}$$

Energy is in particle form at the nebula, but in Poynting flux at the source. How can we probe where the conversion happens?

PSR J7037 A&B - A Laboratory for Late Stellar Evolution, General Relativity and Relativistic Winds and Magnetospheres



$$P_{A} = 22.7 \text{ msec}, \dot{E}_{A} = 6 \times 10^{33} \text{ ergs/s}$$
$$P_{B} = 2773 \text{ msec}, \dot{E}_{B} = 2 \times 10^{30} \text{ ergs/s}$$
$$R_{LA} = 1098 \text{ km}, R_{LB} = 132,400 \text{ km}$$
$$a = 425,000 \text{ km}$$

Pulsar B



Pulsar B brightness as a function of rotation phase (vertical axis) and orbital phase (horizontal axis) - - Lyne et al (2004 - discovery paper)



Doppler shifts of pulse period allow Measurement of orbital parameters

Eclipses of Pulsar A

Lyne et al, Kaspi et al see brief eclipse of A when pulsar B moves in front of A



Pulsed flux of pulsar A for 4 minutes centered on superior conjunction (A behind B). Each data point is a 2 s integration. The eclipse lasts ~30 s, corresponding to a 18,600 km physical dimension of the obscuring material along the direction of orbital motion.



Cartoon of eclipse model magnetosheath absorption by PSR B

PSR J0737 A&B -- laboratory for relativistic winds

Binary orbit:



Ramachandran et al (2004)

600

Lyne et al (2004)

Binary pulsar: outstanding questions

- What is causing the eclipse of pulsar A?
- Why is eclipse asymmetric?
- Why does pulsar B change its pulse shape and brightness?
- What are the properties of pulsar wind in this system?

This system is unique in that mutual orientation of two pulsars changes on human timescale (~10yrs) due to GR. Predictions can be verified.

Strategy: investigate wind-magnetosphere interaction using numerical simulations

Numerical simulation of wind-magnetosphere interaction

Particle-in-cell method:

- Collect currents at the cell centers
- Find fields on the mesh (Maxwell's eqs)
- Interpolate fields to particles positions
- Move particles under Lorentz force

Computational cycle





Code "TRISTAN":

- 2D and 3D cartesian
- Algorithms for rotating magnet
- Handles magnetized flows
- Fully parallelized (128proc+)
- Resolution requirements: typical run 512^3 grid, 10⁷⁻⁸ particles, 50+ Gb. Run at NERSC.
- Diagnostics + visualization

Movie: magnetosheath of pulsar B



Similar to the interaction between Earth magnetosphere and solar wind.

3D magnetosphere: movies



3D magnetosphere: movies



Eclipse and synchrotron absorption

Propose synchrotron absorption in the hot postshock plasma as the eclipse mechanism:

$$\tau_{v} = \frac{\sqrt{3}cr_{e}n_{1\pm}R_{m}}{v_{g1}T_{2}\beta_{2}} \left(\frac{v_{g1}}{v}\right)^{2} I(z) = \frac{2}{v_{1400}^{2}} \left(\frac{\kappa}{10^{6}}\right)^{2} f(\sigma)$$

To explain observed eclipse require:

$$\kappa \sim 3 \times 10^6$$
, $\gamma \sim 10$, $\sigma \sim .1$

The wind is too dense and too slow according to conventional wisdom!

Eclipse duration:



Spindown of pulsar B -- wind torques

Reconnection over half a turn.



Acceleration on the dusk side, deceleration on the dawn side. Net torque:

$$(\dot{J}_B)_{rec} = \int dV(\vec{r} \times \frac{1}{c}\vec{j} \times \vec{B}) \cdot \hat{\Omega}_B \sim \left(\frac{\dot{E}_A}{a^2c}\right)^{1/3} \mu^{4/3} \frac{\Omega_B}{2c\beta_{\parallel}} = 2 \times 10^{30} ergs$$

The wind torque is comparable to the regular spindown torque:

$$(\dot{J}_B)_{spindown} = \frac{\dot{\varepsilon}}{\Omega} = \frac{\mu^2 \Omega^3}{c^3} \left(\frac{R_m}{R_{LB}}\right)^2 \sim 10^{29} ergs$$
 Leads to more accurate B field

Emission morphology of pulsar B



Conclusions

- First simulation of 3D wind-magnetosphere interaction for oblique magnetized rotators.
- Synchrotron absorption in the shocked wind is a likely origin of the eclipse phenomena in the system
- Wind exerts "propellor" torque -- opportunity to study spindown of a pulsar under stress.
- Inferred wind properties are unusual (high particle flux, low magnetization)
- 3D particle-in-cell simulation allows us to study formation of magnetized collisionless shocks. Kinetic nature of the simulation is essential to understanding collisionless shock acceleration.
 Magnetized shocks are ubiquitous in astrophysics (pulsars, AGNs, GRBs), yet acceleration is still not understood -- future direction.