

Lepton Flavor Violation

II. The New Physics Flavor Puzzle

SSI₂₀₁₀: Neutrinos – Nature's mysterious messengers

SLAC, 10 August 2010

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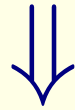
Plan of Lectures

1. The (extended) Standard Model flavor puzzle
2. The New Physics flavor puzzle
 - The NP flavor puzzle
 - The SUSY flavor puzzle
 - LFV at ATLAS/CMS
3. Leptogenesis

The NP Flavor Puzzle

The SM = Low energy effective theory

1. Gravity $\implies \Lambda_{\text{Planck}} \sim 10^{19}$ GeV
2. $m_\nu \neq 0 \implies \Lambda_{\text{Seesaw}} \leq 10^{15}$ GeV
3. m_H^2 -fine tuning; Dark matter = WIMP $\implies \Lambda_{\text{NP}} \sim$ TeV



- The SM = Low energy effective theory
- Must write non-renormalizable terms suppressed by $\Lambda_{\text{NP}}^{d-4}$
- $\mathcal{L}_{d=5} = \frac{Z_{ij}^\nu}{\Lambda_{\text{seesaw}}} L_i L_j \phi \phi$
- $\mathcal{L}_{d=6}$ contains many flavor changing operators

New Physics - I

- The effects of new physics at a high energy scale Λ_{NP} can be presented as higher dimension operators

- For example, we expect the following dimension-six operators:

$$\mathcal{L}_{d=6} \supset \sum_{i \neq j} \frac{z_{ij}^{ee}}{\Lambda_{\text{NP}}^2} (\overline{\ell_{Lj}} \gamma_\mu \ell_{Li}) (\overline{e_L} \gamma_\mu e_L)$$

- New contribution to $\ell_i \rightarrow \ell_j e^+ e^-$

$$B_{\ell_i \rightarrow \ell_j e^+ e^-} \equiv \frac{\text{BR}(\ell_i \rightarrow \ell_j e^+ e^-)}{\text{BR}(\ell_i \rightarrow \ell_j \nu_i \bar{\nu}_j)} = \frac{(|z_{ij}^{ee}| / \Lambda_{\text{NP}}^2)^2}{G_F^2 / 2}$$

- Given an upper bound on $B_{ij} \implies$ an upper bound on $z_{ij}^{ee} / \Lambda_{\text{NP}}^2$

- Generic flavor structure $\equiv z_{ij}^{ee} \sim 1$ or, perhaps, loop – factor

New Physics - II

- Another example:

$$\mathcal{L}_{d=6} \supset \frac{evz^{\gamma M}}{\Lambda_{\text{NP}}^2} \bar{\ell}_i \sigma_{\mu\nu} F^{\mu\nu} P_M \ell_j$$

- New contribution to $\ell_i \rightarrow \ell_j \gamma$

$$B_{\ell_i \rightarrow \ell_j \gamma} \equiv \frac{\text{BR}(\ell_i \rightarrow \ell_j \gamma)}{\text{BR}(\ell_i \rightarrow \ell_j \nu_i \bar{\nu}_j)} = \frac{48\pi^3 \alpha v^2}{G_F^2 m_{\ell_i}^2} \frac{|z_{ij}^{\gamma L}|^2 + |z_{ij}^{\gamma R}|^2}{\Lambda_{\text{NP}}^4}$$

- Given an upper bound on $B_{ij} \implies$ an upper bound on $z_{ij}^{\gamma M} / \Lambda_{\text{NP}}^2$
- Generic flavor structure $\equiv z_{ij}^{\gamma} \sim \frac{1}{16\pi^2}$

Some data

$$B_{\mu^- \rightarrow e^- e^+ e^-} < 1.0 \times 10^{-12}$$

$$B_{\tau^- \rightarrow e^- e^+ e^-} < 2.0 \times 10^{-7}$$

$$B_{\tau^- \rightarrow \mu^- e^+ e^-} < 1.6 \times 10^{-7}$$

$$B_{\mu^- \rightarrow e^- \gamma} < 1.2 \times 10^{-11}$$

$$B_{\tau^- \rightarrow e^- \gamma} < 6.2 \times 10^{-7}$$

$$B_{\tau^- \rightarrow \mu^- \gamma} < 2.6 \times 10^{-7}$$

High Scale?

- For $z_{ij}^{ee} \sim 1$: $\Lambda_{\text{NP}} \gtrsim B^{-1/4} \times 350 \text{ GeV}$
- For $z_{ij}^{\gamma M} \sim 1/16\pi^2$: $\Lambda_{\text{NP}} \gtrsim B^{-1/4} (m_\mu/m_{\ell_i})^{1/2} \times 1.7 \text{ TeV}$

	$\Lambda_{\text{NP}} \gtrsim$
$B_{\mu^- \rightarrow e^- e^+ e^-}$	350 TeV
$B_{\tau^- \rightarrow e^- e^+ e^-}$	17 TeV
$B_{\tau^- \rightarrow \mu^- e^+ e^-}$	17 TeV
$B_{\mu^- \rightarrow e^- \gamma}$	900 TeV
$B_{\tau^- \rightarrow e^- \gamma}$	15 TeV
$B_{\tau^- \rightarrow \mu^- \gamma}$	19 TeV

High Scale

- For generic lepton flavor parameters, $\Lambda_{\text{NP}} \gtrsim 10^3 \text{ TeV}$
- For generic quark flavor parameters, $\Lambda_{\text{NP}} \gtrsim 10^4 \text{ TeV}$



Did we misinterpret the Higgs fine tuning problem?

Did we misinterpret the dark matter puzzle (it is not WIMP)?

Small (hierarchical?) flavor parameters?

- For $\Lambda_{\text{NP}} \sim 1 \text{ TeV}$, $z_{ij}^{ee} \lesssim 10B_{ij}^{1/2}$; $16\pi^2 z_{ij}^{\gamma M} \lesssim 0.3B_{ij}^{1/2} (m_{\ell_i}/m_\mu)$

ij	Observable	$z_{ij}^{ee} \lesssim$
μe	$B_{\mu^- \rightarrow e^- e^+ e^-}$	8.2×10^{-6}
τe	$B_{\tau^- \rightarrow e^- e^+ e^-}$	3.7×10^{-3}
$\tau \mu$	$B_{\tau^- \rightarrow \mu^- e^+ e^-}$	3.3×10^{-3}
ij	Observable	$16\pi^2 z_{ij}^{\gamma M} \lesssim$
μe	$B_{\mu^- \rightarrow e^- \gamma}$	1.3×10^{-6}
τe	$B_{\tau^- \rightarrow e^- \gamma}$	4.4×10^{-3}
$\tau \mu$	$B_{\tau^- \rightarrow \mu^- \gamma}$	2.9×10^{-3}

Small (hierarchical?) flavor parameters

For $\Lambda_{\text{NP}} \sim \text{TeV}$,

- $z_{\mu e} < 10^{-5}$
- $z_{\tau \ell} < 10^{-2}$
- $z_{sd} < 10^{-7}$



The flavor structure of NP@TeV must be highly non-generic

How? Why? = The NP flavor puzzle

How does the SM ($\Lambda_{\text{SM}} \sim m_W$) do it?

		$z_{ij} \sim$	z_{ij}^{SM}
$\Delta m_K/m_K$	7.0×10^{-15}	5×10^{-9}	$\alpha_2^2 y_c^2 V_{cd} V_{cs} ^2$
$\Delta m_D/m_D$	8.7×10^{-15}	5×10^{-9}	Long Distance
$\Delta m_B/m_B$	6.3×10^{-14}	7×10^{-8}	$\alpha_2^2 y_t^2 V_{td} V_{tb} ^2$
$\Delta m_{B_s}/m_{B_s}$	2.1×10^{-12}	2×10^{-6}	$\alpha_2^2 y_t^2 V_{ts} V_{tb} ^2$
		$\frac{\text{Im}(z_{ij})}{ z_{ij} } \sim$	$\frac{\text{Im}(z_{ij}^{\text{SM}})}{ z_{ij}^{\text{SM}} }$
ϵ_K	2.3×10^{-3}	$\mathcal{O}(0.01)$	$\text{Im} \frac{y_t^2 (V_{td}^* V_{ts})^2}{y_c^2 (V_{cd}^* V_{cs})^2} \sim 0.01$
A_Γ	≤ 0.004	≤ 0.2	0
$S_{\psi K_S}$	0.67 ± 0.02	$\mathcal{O}(1)$	$\text{Im} \frac{V_{tb} V_{td}^*}{V_{tb}^* V_{td}} \frac{V_{cb}^* V_{cd}}{V_{cb} V_{cd}^*} \sim 0.7$
$S_{\psi\phi}$	≤ 1	≤ 1	$\text{Im} \frac{V_{tb} V_{ts}^*}{V_{tb}^* V_{ts}} \frac{V_{cb}^* V_{cs}}{V_{cb} V_{cs}^*} \sim 0.02$

- Does the new physics know the SM Yukawa structure? (MFV)

What does the ESM do?

	$z_{ij}^{\max} (\Lambda \sim m_W)$	z_{ij}^{ESM}	With $U_{e3} \sim 0.1$
$z_{\mu e}^{ee}$	8×10^{-8}	$\frac{\alpha_2^2}{m_W^2} (U_{e2} U_{\mu 2}^* \Delta m_{21}^2 + U_{e3} U_{\mu 3}^* \Delta m_{32}^2)$	3×10^{-29}
$z_{\tau e}^{ee}$	4×10^{-5}	$\frac{\alpha_2^2}{m_W^2} (U_{e2} U_{\tau 2}^* \Delta m_{21}^2 + U_{e3} U_{\tau 3}^* \Delta m_{32}^2)$	3×10^{-29}
$z_{\tau \mu}^{ee}$	3×10^{-5}	$\frac{\alpha_2^2}{m_W^2} (U_{\mu 2} U_{\tau 2}^* \Delta m_{21}^2 + U_{\mu 3} U_{\tau 3}^* \Delta m_{32}^2)$	4×10^{-28}
$z_{\mu e}^{\gamma L}$	8×10^{-11}	$\frac{\alpha_2 y_e}{4\pi m_W^2} (U_{e2} U_{\mu 2}^* \Delta m_{21}^2 + U_{e3} U_{\mu 3}^* \Delta m_{32}^2)$	2×10^{-34}
$z_{\tau e}^{\gamma L}$	3×10^{-7}	$\frac{\alpha_2 y_e}{4\pi m_W^2} (U_{e2} U_{\tau 2}^* \Delta m_{21}^2 + U_{e3} U_{\tau 3}^* \Delta m_{32}^2)$	2×10^{-34}
$z_{\tau \mu}^{\gamma L}$	2×10^{-7}	$\frac{\alpha_2 y_\mu}{4\pi m_W^2} (U_{\mu 2} U_{\tau 2}^* \Delta m_{21}^2 + U_{\mu 3} U_{\tau 3}^* \Delta m_{32}^2)$	6×10^{-31}
$z_{\mu e}^{\gamma R}$	8×10^{-11}	$\frac{\alpha_2 y_\mu}{4\pi m_W^2} (U_{e2} U_{\mu 2}^* \Delta m_{21}^2 + U_{e3} U_{\mu 3}^* \Delta m_{32}^2)$	4×10^{-32}
$z_{\tau e}^{\gamma R}$	3×10^{-7}	$\frac{\alpha_2 y_\tau}{4\pi m_W^2} (U_{e2} U_{\tau 2}^* \Delta m_{21}^2 + U_{e3} U_{\tau 3}^* \Delta m_{32}^2)$	7×10^{-31}
$z_{\tau \mu}^{\gamma R}$	2×10^{-7}	$\frac{\alpha_2 y_\tau}{4\pi m_W^2} (U_{\mu 2} U_{\tau 2}^* \Delta m_{21}^2 + U_{\mu 3} U_{\tau 3}^* \Delta m_{32}^2)$	1×10^{-29}

- The ESM LFV rates are negligibly tiny

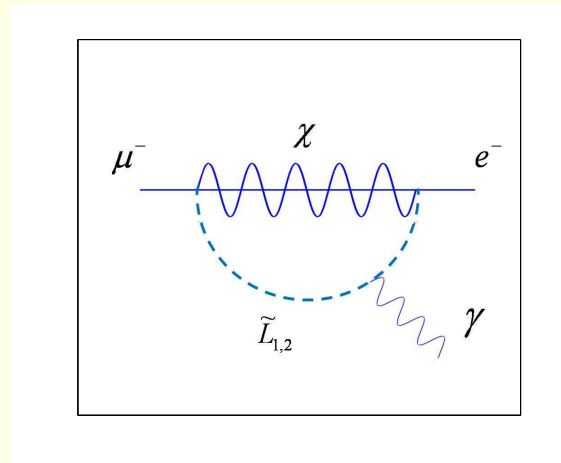
Intermediate summary

- The SM predicts no LFV
- The ESM predicts unobservably small LFV
- Any signal of LFV in charged lepton decays - an unambiguous signal of new physics beyond the ESM ($\Lambda_{\text{NP}} \lesssim 10^3 \text{ TeV}$)
- The flavor structure of new physics at the TeV scale must be highly non-generic (\implies The NP flavor puzzle)
- The NP flavor puzzle guarantees that such an LFV signal will provide important information about the NP

The Supersymmetric Flavor Puzzle

The $\mu \rightarrow e\gamma$ challenge

Take, for example, the contribution from the first two generations of slepton doublets to $\mu \rightarrow e\gamma$:



$$\Lambda_{\text{NP}} = m_{\tilde{L}}; \quad z_{\mu e}^{\gamma L} \simeq \left(\frac{\alpha_2}{60\pi} \frac{m_\mu \tan \beta}{v} \right) \frac{\Delta m_{21}^2}{m_{\tilde{L}}^2} K_{e2}^L K_{\mu 2}^{L*}$$

$$\Rightarrow \left(\frac{300 \text{ GeV}}{m_{\tilde{L}}} \right)^2 \times \frac{\Delta m_{21}^2}{m_{\tilde{L}}^2} \times \sin 2\theta \lesssim 10^{-3}$$

How can Supersymmetry do it?

$$\frac{\text{TeV}}{\tilde{m}} \times \frac{\Delta\tilde{m}_{ij}^2}{\tilde{m}^2} \times K_{ij} \ll 1$$

Why? = The SUSY flavor puzzle

How can Supersymmetry do it?

$$\frac{\text{TeV}}{\tilde{m}} \times \frac{\Delta\tilde{m}_{ij}^2}{\tilde{m}^2} \times K_{ij} \ll 1$$

Why? = The SUSY flavor puzzle

- Solutions:

- Heaviness: $\tilde{m} \gg 1 \text{ TeV}$
- Degeneracy: $\Delta\tilde{m}_{ij}^2 \ll \tilde{m}^2$
- Alignment: $K_{ij} \ll 1$
- Split Supersymmetry
- Gauge-mediation
- Horizontal symmetries

Gauge Mediation – Squarks

- $\Delta m_K, \Delta m_D$:
 - The first two squark generations are quasi-degenerate
 - Could be the result of alignment ($\sim \sin \theta_C$) and RGE ($\frac{\Delta m^2}{m^2} \sim 0.1$)
 - Natural with gauge mediation ($\frac{\Delta m^2}{m^2} \sim y_c^2 \sim 10^{-4}$)
- Gauge mediation:
 - $\widetilde{M}_{\tilde{q}_L}^2 = \tilde{m}^2 \mathbf{1} + D_{q_L} \mathbf{1} + v_q^2 Y_q Y_q^\dagger$
 - RGE: $\tilde{m}_{\tilde{Q}_L}^2(m_Z) = \tilde{m}^2(r_3 \mathbf{1} + c_u Y_u Y_u^\dagger + c_d Y_d Y_d^\dagger)$
 - The only source of flavor violation = The SM Yukawa couplings
 - An example of minimal flavor violation (MFV)
 - MFV solves all SUSY flavor problems

The seesaw mechanism - a reminder

- Add $N_i(1, 1)_0$
- $\mathcal{L}_{\text{leptons}} = Y_{ij}^e \bar{L}_i \phi_d E_j + Y_{ij}^\nu \bar{L}_i \phi_u N_j + M_j N_j N_j$
- Gives light neutrino masses $(m_\nu)_{ij} = \langle \phi_u \rangle^2 Y_{ik}^\nu Y_{jk}^\nu M_k^{-1}$
- Flavor: $SU(3)_L \times SU(3)_E \times SU(3)_N$ completely broken
- $15_R + 6_I$ lepton flavor parameters
(compared to $9_R + 3_I$ in the ESM)

Gauge Mediation – Sleptons

- $\mu \rightarrow e\gamma, \mu \rightarrow eee$:
 - Suggestive: The first two slepton gen's – quasi-degenerate
 - Alignment of $\mathcal{O}(|U_{e2}|)$ and RGE do not help
 - Natural with gauge mediation
- Gauge mediation (GM) with $\Lambda_{\text{seesaw}} < \Lambda_{\text{GM}}$:
 - $\widetilde{M}_{\widetilde{\ell}_L}^2 = \widetilde{m}^2 \mathbf{1} + D_{\ell_L} \mathbf{1} + v_d^2 Y^e Y^{e\dagger}$
 - LFV from RGE: $(\widetilde{m}_{\widetilde{L}_L}^2)_{ij} \simeq -\frac{3\widetilde{m}^2}{8\pi^2} Y_{ik}^\nu Y_{jk}^{\nu*} \ln\left(\frac{\Lambda_{\text{GM}}}{M_k}\right)$
 - $SU(3)_L \times SU(3)_N$ broken by Y^e, Y^ν, M_N
 - Observable LFV effects – possible
 - Even with minimal lepton flavor violation (MLFV), we can obtain new information on flavor

MFV + SUSY

- Squarks:
 - Spectrum: $2 + 1$
 - Decays: $2 \rightarrow u, d, s, c, \quad 1 \rightarrow t, b$
- Sleptons, $\Lambda_{\text{seesaw}} > \Lambda_{\text{mediation}}$:
 - spectrum: 3
 - Decays: flavor diagonal
- Sleptons, $\Lambda_{\text{seesaw}} < \Lambda_{\text{mediation}}$:
 - Y^ν, M_N may leave a footprint on the slepton spectrum and flavor decomposition

LFV@ATLAS/CMS

Flavor Physics at the LHC era

- If ATLAS/CMS observe no NP...
- and flavor factories observe no NP...

Flavor Physics at the LHC era

- If ATLAS/CMS observe no NP...
- but flavor factories observe NP...
- We may have misinterpreted the fine-tuning problem
- We may have misinterpreted the dark matter puzzle
- Flavor will provide the only clue for an accessible scale of NP

Flavor Physics at the LHC era

- ATLAS/CMS will, hopefully, observe NP at $\Lambda_{\text{NP}} \lesssim \text{TeV}$;
- If the NP couple to SM quarks and/or leptons –

There are new flavor parameters that can, in principle, be measured
- In combination with flavor factories, we may...
 - Understand how the NP flavor puzzle is (not) solved
 - Probe NP at $\Lambda_{\text{NP}} \gg \text{TeV}$
 - Get hints about the solution to the SM flavor puzzle

Gauge+Gravity Mediation

- Example: High (but not too high) scale gauge mediation
 - Gravity mediation sub-dominant but non-negligible
 - $r = \frac{\text{gravity-med}}{\text{gauge-med}} \sim \left(\frac{m_M}{m_P}\right)^2 \left(\frac{4\pi}{\alpha_2(m_M)}\right)^2 \frac{3}{8n_M}$
 - $\widetilde{M}_{\widetilde{L}_L}^2(m_M) = \widetilde{m}_{\widetilde{L}_L}^2 (\mathbf{1} + r X_{\widetilde{L}_L})$
 - Degeneracy depends on r

Assume: The flavor structure of X determined by FN:

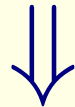
- $X_{\widetilde{L}_L} \sim \begin{pmatrix} 1 & U_{e2} & U_{e3} \\ \cdot & 1 & U_{\mu 3} \\ \cdot & \cdot & 1 \end{pmatrix}; \quad X_{\widetilde{E}_R} \sim \begin{pmatrix} 1 & \frac{m_e/m_\mu}{U_{e2}} & \frac{m_e/m_\tau}{U_{e3}} \\ \cdot & 1 & \frac{m_\mu/m_\tau}{U_{\mu 3}} \\ \cdot & \cdot & 1 \end{pmatrix}$

- Mixing depends only on X which is related to the SM flavor

Solving the NP Flavor Puzzle

If ATLAS/CMS observe sleptons...

- Determine the slepton mass scale (\tilde{m})
- Determine the slepton mass splitting ($m_{\tilde{\ell}_j} - m_{\tilde{\ell}_i}$)
- Determine the sfermion flavor decomposition (K_{ij}^e)



Learn how the SUSY flavor suppression is obtained

Physics at $\Lambda_{\text{NP}} \gg \Lambda_{\text{LHC}}$

If ATLAS/CMS determine slepton mass splittings...

- Find the ratio between gravity- and gauge-mediated contributions (r)
- Determine the messenger scale of gauge mediation (m_M)
- Find the hierarchy between the GMSB and see-saw scales



Probe physics at $m_M \sim 10^{15}$ GeV

Solving the SM Flavor Puzzle?

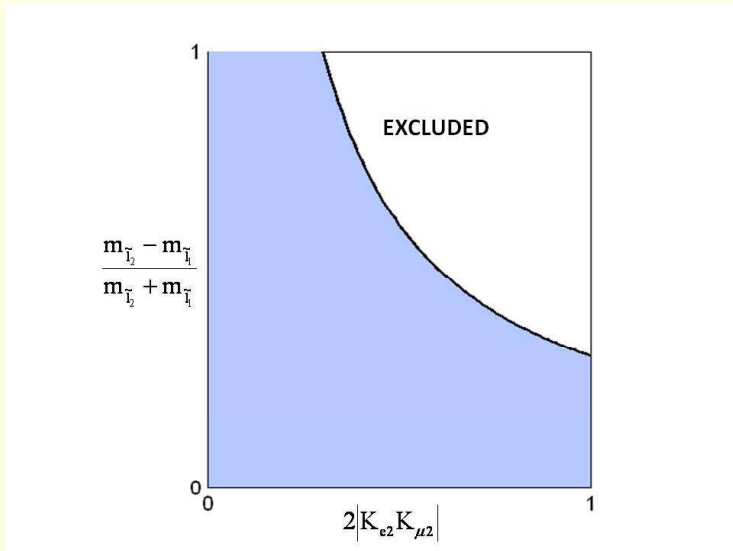
If ATLAS/CMS determine slepton flavor decomposition...

- Determine X of $\tilde{M}^2 = \tilde{m}^2(\mathbf{1} + rX)$
- Does X have the FN-predicted structure?



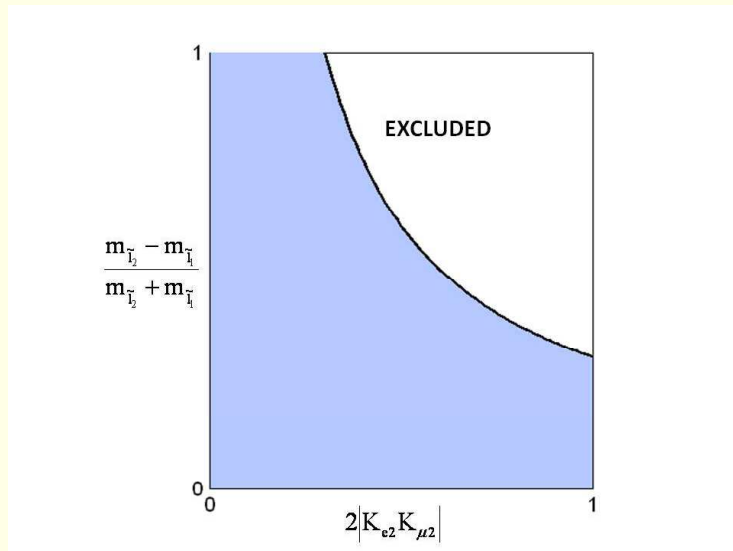
Test theories that explain the SM flavor structure

The SUSY flavor plane

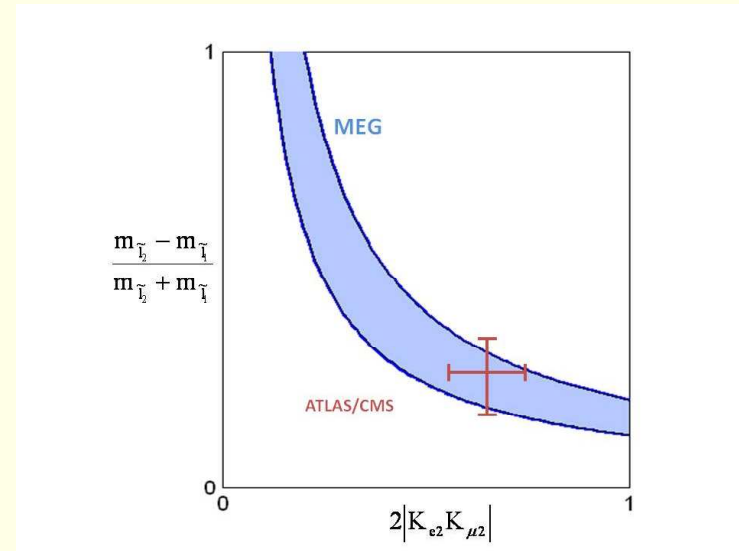


Flavor Factories

The SUSY flavor plane

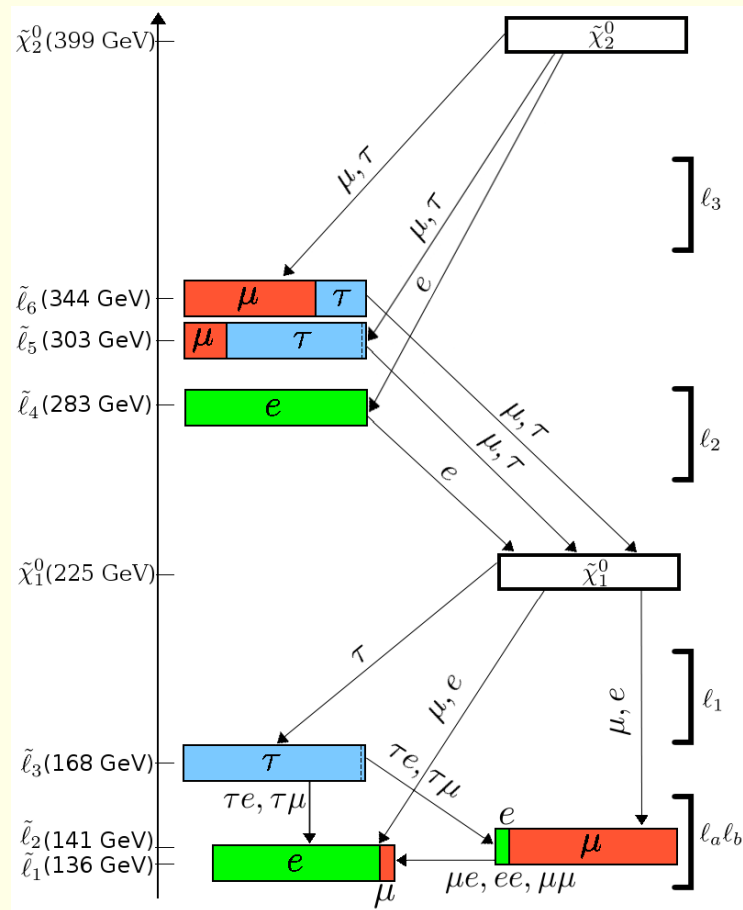


Flavor Factories



FF+ATLAS/CMS

Gauge+Gravity Mediation



$$\chi_1^0 \rightarrow \tilde{\ell}_1^\pm \tilde{\ell}_1^\mp$$

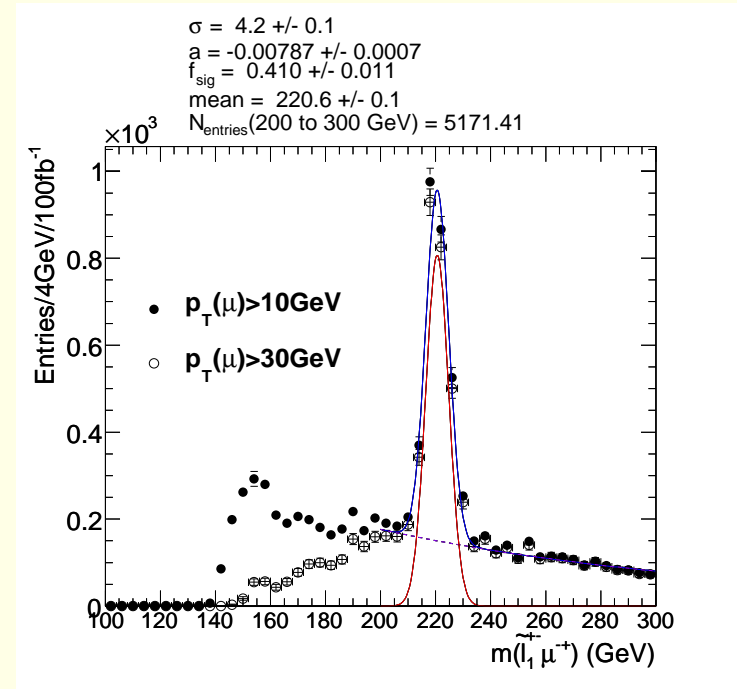
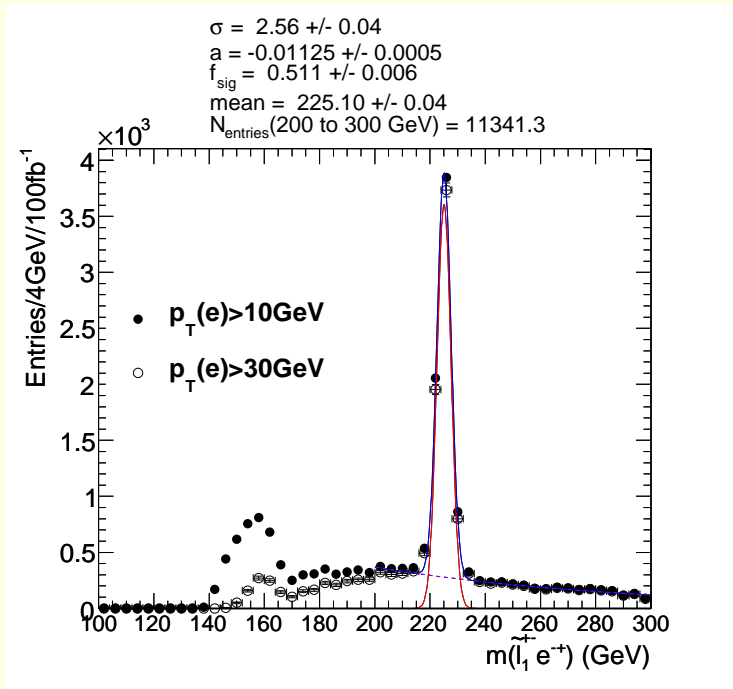
$$\chi_1^0 \rightarrow \tilde{\ell}_2^\pm \tilde{\ell}_2^\mp$$

$$\tilde{\ell}_2^\pm \rightarrow \tilde{\ell}_1^\pm X^{\pm\mp}$$

or
$$\tilde{\ell}_2^\pm \rightarrow \tilde{\ell}_1^\mp X^{\pm\pm}$$

$$m_{\tilde{G}} \ll m_{\tilde{\ell}_{1,2,3}} \ll m_{\chi_1^0}$$

Measuring slepton mass splitting

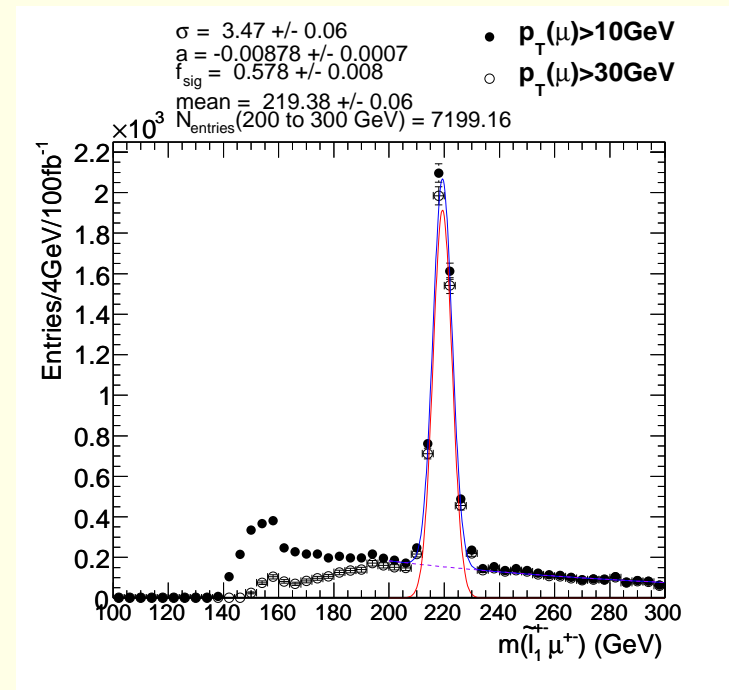
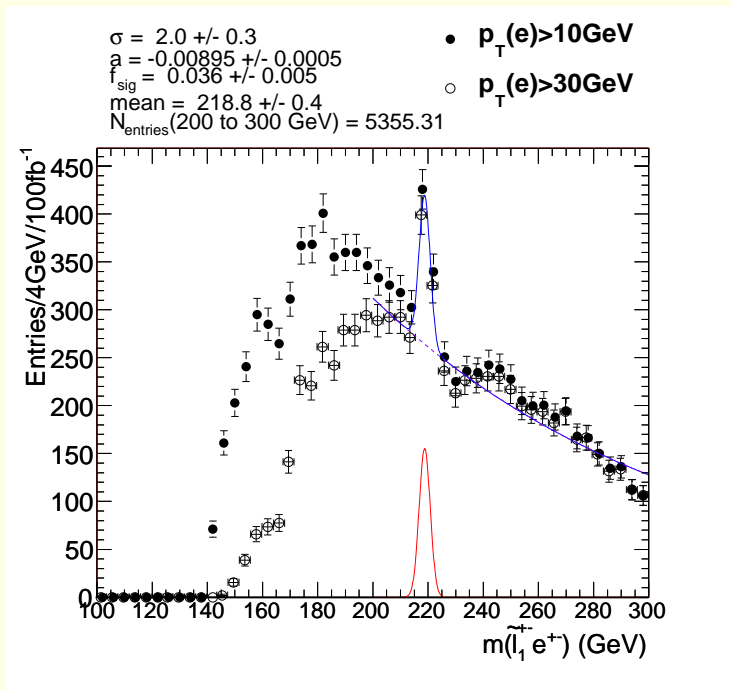


$$\sqrt{m_{\tilde{\ell}_1 e}^2} = M_{\chi_1^0}$$

$$\sqrt{m_{\tilde{\ell}_1 \mu}^2} = M_{\chi_1^0} - E_{\text{shift}}$$

$$\Delta m_{21} = \frac{2m_{\chi_1^0} m_{\tilde{\ell}_1}}{m_{\chi_1^0}^2 + m_{\tilde{\ell}_1}^2} E_{\text{shift}}$$

Measuring slepton flavor decomposition



$$\chi_1^0 \rightarrow \tilde{l}_1^\pm e^\pm X_{\text{invisible}}$$

$$\chi_1^0 \rightarrow \tilde{l}_1^\pm \mu^\pm X_{\text{invisible}}$$

$$|K_{e2}/K_{\mu2}|^2 = N_{SSe}/N_{SS\mu}$$

[Feng, Lester, Nir, Shadmi *et al.*, PRD77(2008)076002; PRD80(2009)114004; JHEP01(2010)047]

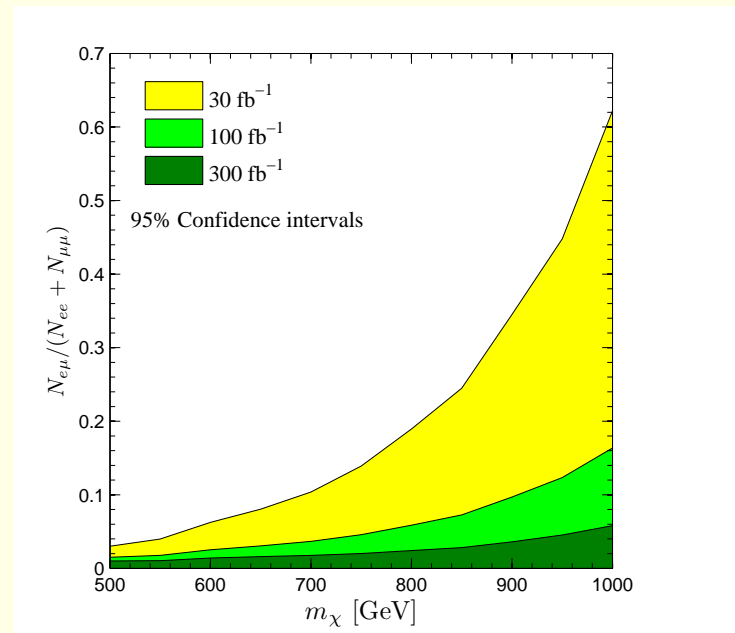
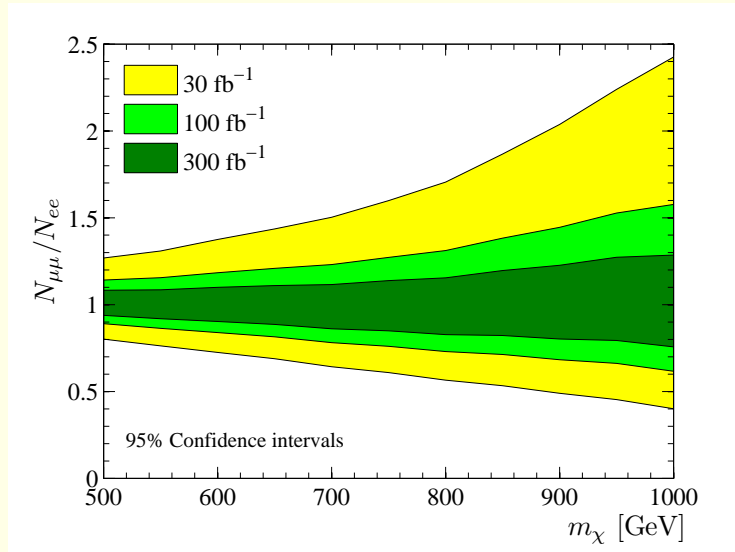
Lessons from $\tilde{\ell}_1, e, \mu$

- Determine Δm_{21} and $\sin 2\theta_{12}$:
Is it consistent with $\mu \rightarrow e\gamma$?
How the SUSY flavor problem is solved
- Determine $\Delta m_{21}, \Delta m_{54}, \dots$:
What is messenger scale of gauge mediation (M_m)?
Probe physics at $M_m \sim 10^{15}$ GeV
- Determine $|K_{e2}/K_{\mu2}|$:
Is the FN mechanism at work?
How the SM flavor puzzle is solved

Vector-like leptons and MLFV

- Imagine: Vector-like lepton doublets with $m \lesssim TeV$
 - Avoid large FCNC by MLFV
 - The only LFV comes from $Y^E = \text{diag}(y_e, y_\mu, y_\tau)$
 - The heavy mass spectrum:
quasi-degeneracy or hierarchy $\propto Y^E$
 - The heavy-to-light couplings:
universal or hierarchical (affects the lifetimes)
 - The heavy-to-light couplings:
flavor-diagonal

Vector-like leptons and MLFV



- $N_{ee} \neq N_{\mu\mu}$ and/or $N_{e\mu} \neq 0$:
Either MLFV with ν -related spurions or non-MLFV
- $N_{ee} = N_{\mu\mu}$ and $N_{e\mu} = 0$: Approximate $U(1)_e \times U(1)_\mu$
Plus $m_{\chi_e} \approx m_{\chi_\mu}$: Approximate $U(2)_{e\mu}$

[Gross, Grossman, Nir, Vitells, PRD81(10)055013 [1001.2883]]

The role of flavor factories (FF)

ATLAS/CMS and flavor factories give complementary information

- In the absence of NP at ATLAS/CMS, flavor factories will be crucial to find Λ_{NP}
- Consistency between ATLAS/CMS and FF is necessary to understand the NP flavor puzzle
- NP in $c \rightarrow u?$ $s \rightarrow d?$ $b \rightarrow d?$ $b \rightarrow s?$ $t \rightarrow c?$ $t \rightarrow u?$
 $\mu \rightarrow e?$ $\tau \rightarrow \mu?$ $\tau \rightarrow e?$
 - MFV?
 - Structure related to SM?
 - Structure unrelated to SM?
 - Anarchy?

[Hiller, Hochberg, Nir, JHEP0903(09)115; JHEP1003(10)079 [1001.1513]]

Thanks to my high- p_t -physics collaborators:

Yuval Grossman, YN, Jesse Thaler, Tomer Volansky, Jure Zupan
Phys. Rev. D76 (2007) 096006 [arXiv:0706.1845]

Jonathan Feng, Christopher Lester, YN, Yael Shadmi
Phys. Rev. D77 (2008) 076002 [arXiv:0712.0674]

Jonathan Feng, Sky French, Christopher Lester, YN, Yael Shadmi
Phys. Rev. D80 (2009) 114004 [arXiv:0906.4215]

Feng, French, Galon, Lester, YN, Shadmi, Sanford, Yu
JHEP 1001 (2010) 047 [arXiv:0910.1618]

Eilam Gross, Daniel Grossman, YN, Ofer Vitells
Phys. Rev. D81 (2010) 055013 [arXiv:1001.2883]

Thanks to my students:

Kfir Blum, Daniel Grossman, Yonit Hochberg