

Theory of Resonant x-ray spectroscopy

M. W. Haverkort

Max Planck Institute for Solid State Research, Stuttgart, Germany

University of British Columbia, Vancouver, Canada

maurits@physics.ubc.ca

Spectroscopy: interaction of photons with matter

Photons: x-rays

Matter: Transition metal compounds

Group	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
Period																			
1	1 H																		2 He
2	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne	
3	11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar	
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr	
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe	
6	55 Cs	56 Ba	* 71 Lu	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn	
7	87 Fr	88 Ra	** 103 Lr	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Uub	113 Uut	114 Uuq	115 Uup	116 Uuh	117 Uus	118 Uuo	
*Lanthanoids			* 57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb			
**Actinoids			** 89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No			

Spectroscopy: interaction of photons with matter

Photons: x-rays

Matter: Transition metal compounds

Group	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
Period																		
1	1 H																	2 He
2	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
3	11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
6	55 Cs	56 Ba	* 71 Lu	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
7	87 Fr	88 Ra	** 103 Lr	104 Rf	105 Db	106 Sg	107 Bh	108 Hs										



Important for technology

Important in biochemistry (active centers of enzymes)

Often open shell systems, involved to understand

*Lanthanoids	57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb
**Actinoids	89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No

Transition metal compounds show a large variety of properties

Exchange bias

PHYSICAL REVIEW

VOLUME 102, NUMBER 5

JUNE 1, 1956

Letters to the Editor

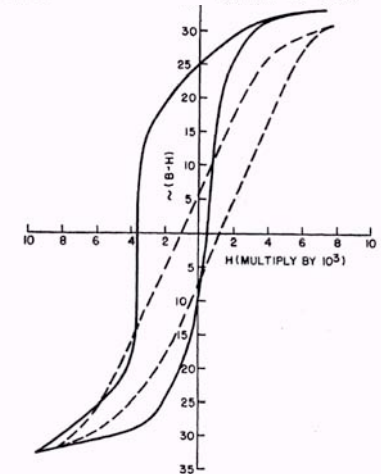
PUBLICATION of brief reports of important discoveries in physics may be secured by addressing them to this department. The closing date for this department is five weeks prior to the date of issue. No proof will be sent to the authors. The Board of Editors does not hold itself responsible for the opinions expressed by the correspondents. Communications should not exceed 600 words in length and should be submitted in duplicate.

New Magnetic Anisotropy

W. H. MEIKLEJOHN AND C. P. BEAN

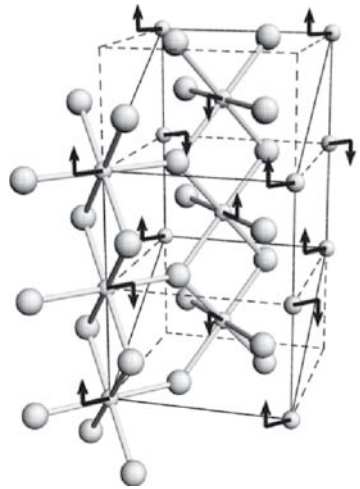
General Electric Research Laboratory, Schenectady, New York

(Received March 7, 1956)

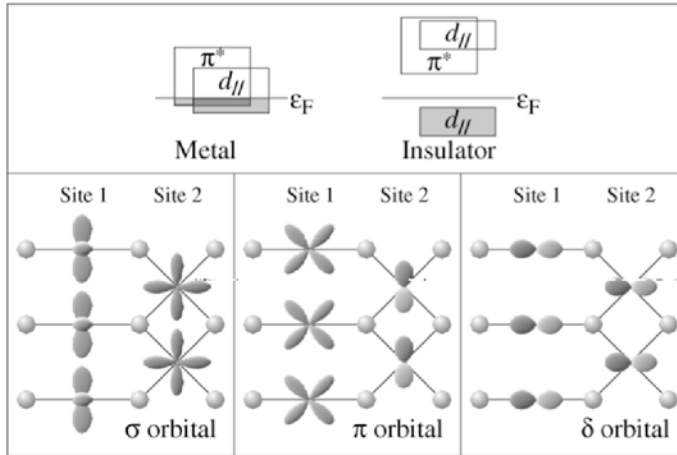


Transition metal compounds show a large variety of properties

Metal Insulator transitions



Phys. Rev. Lett. 95, 196404 (2005)



Exchange bias

PHYSICAL REVIEW

VOLUME 102, NUMBER 5

JUNE 1, 1956

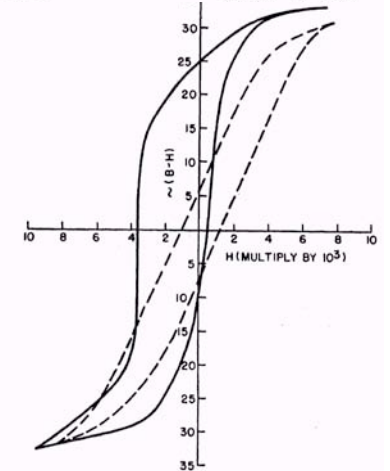
Letters to the Editor

PUBLICATION of brief reports of important "discoveries" in physics may be secured by addressing them to this department. The closing date for this department is five weeks prior to the date of issue. No proof will be sent to the authors. The Board of Editors does not hold itself responsible for the opinions expressed by the correspondents. Communications should not exceed 600 words in length and should be submitted in duplicate.

New Magnetic Anisotropy

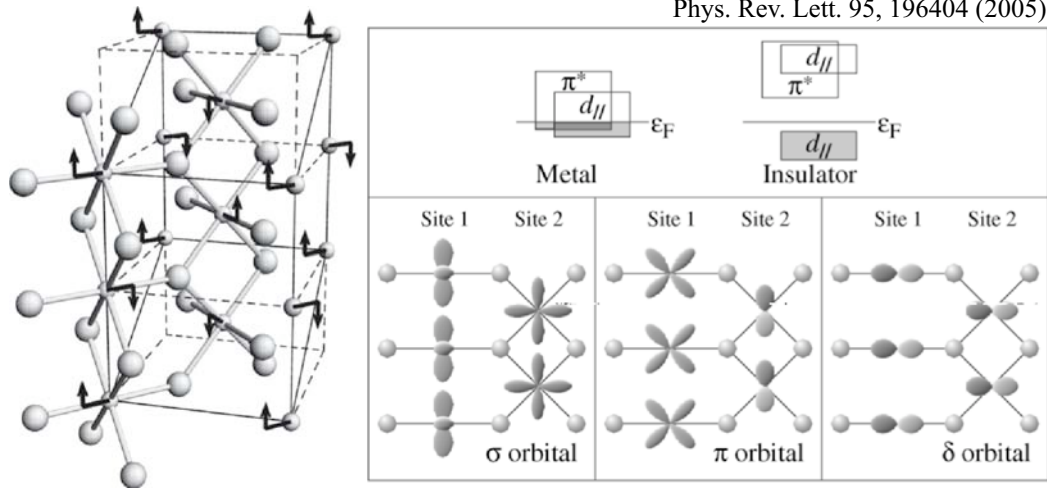
W. H. MEIKLEJOHN AND C. P. BEAN

General Electric Research Laboratory, Schenectady, New York
(Received March 7, 1956)



Transition metal compounds show a large variety of properties

Metal Insulator transitions



Exchange bias

PHYSICAL REVIEW VOLUME 102, NUMBER 5 JUNE 1, 1956

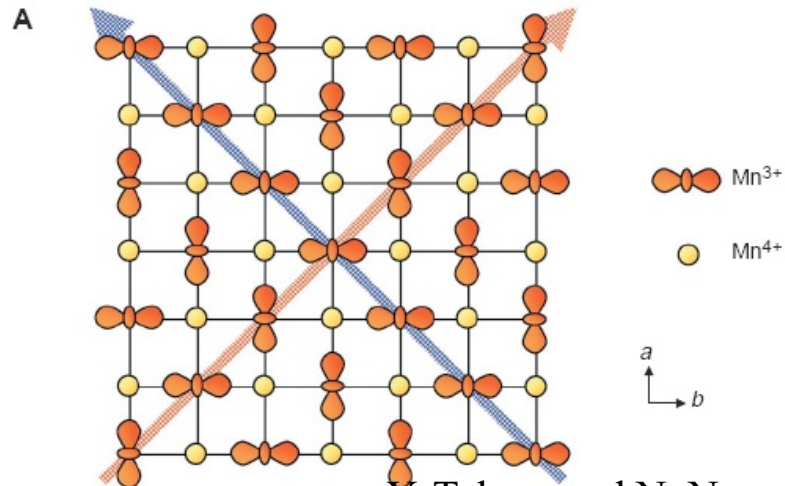
Letters to the Editor

PUBLICATION of brief reports of important discoveries in physics may be secured by addressing them to this department. The closing date for this department is five weeks prior to the date of issue. No proof will be sent to the authors. The Board of Editors does not hold itself responsible for the opinions expressed by the correspondents. Communications should not exceed 600 words in length and should be submitted in duplicate.

New Magnetic Anisotropy

W. H. MEIKLEJOHN AND C. P. BEAN
 General Electric Research Laboratory, Schenectady, New York
 (Received March 7, 1956)

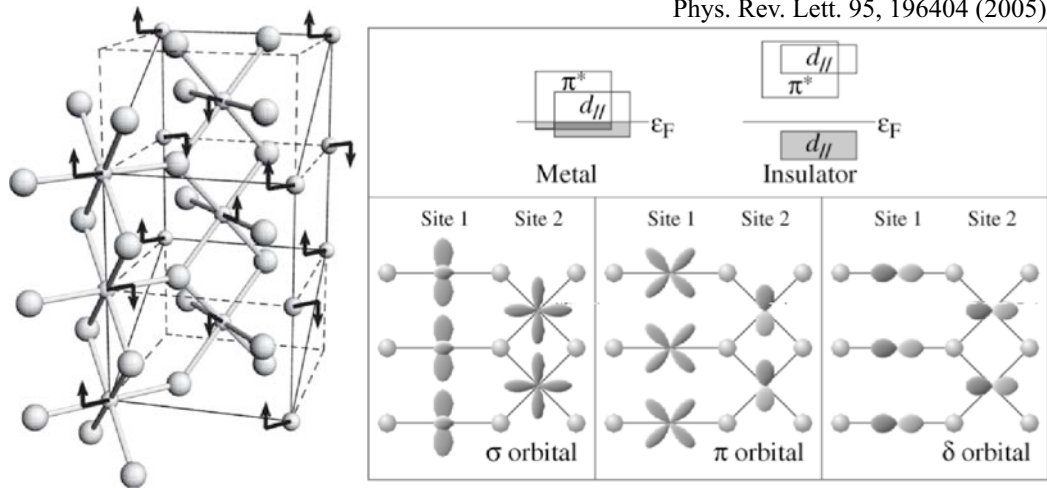
Colossal magneto resistance



Y. Tokura and N. Nagaosa
 Science 288, 462 (2000).

Transition metal compounds show a large variety of properties

Metal Insulator transitions



Exchange bias

PHYSICAL REVIEW VOLUME 102, NUMBER 5 JUNE 1, 1956

Letters to the Editor

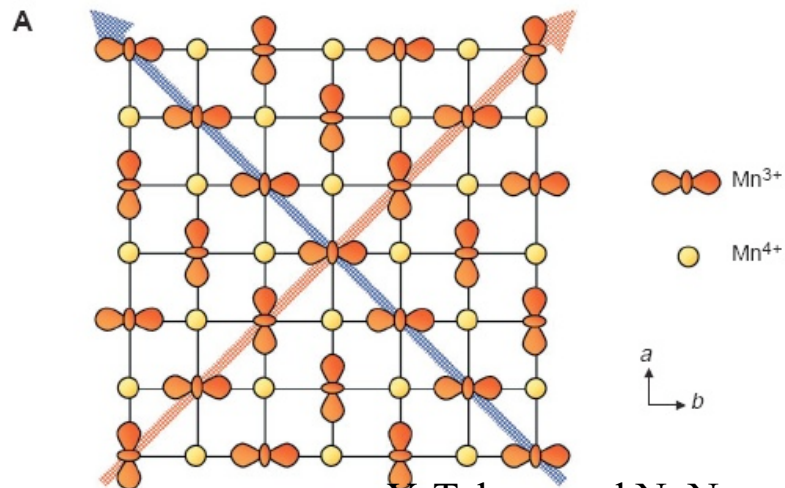
PUBLICATION of brief reports of important discoveries in physics may be secured by addressing them to this department. The closing date for this department is five weeks prior to the date of issue. No proof will be sent to the authors. The Board of Editors does not hold itself responsible for the opinions expressed by the correspondents. Communications should not exceed 600 words in length and should be submitted in duplicate.

New Magnetic Anisotropy

W. H. MEIRLEJOHN AND C. P. BEAN
 General Electric Research Laboratory, Schenectady, New York
 (Received March 7, 1956)

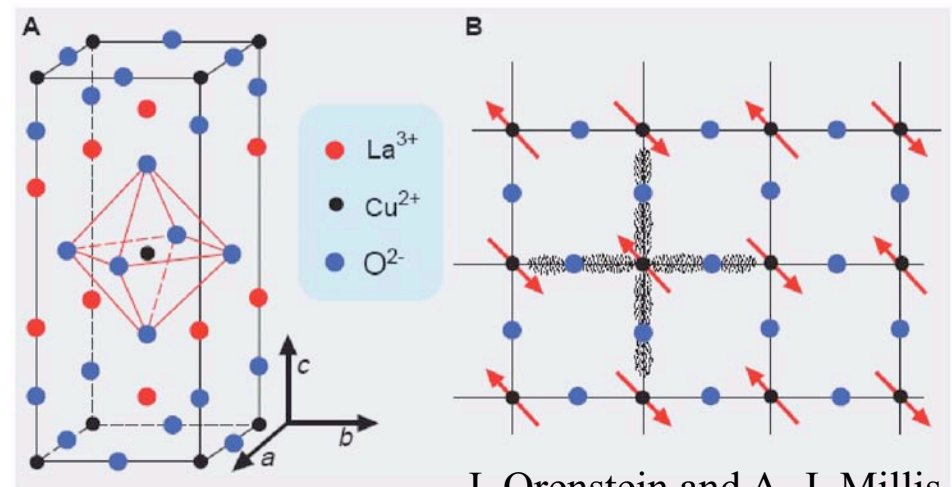
The graph shows magnetization M versus magnetic field H (multiplied by 10^3). The curves exhibit hysteresis loops that are shifted along the field axis, characteristic of exchange bias. The y-axis ranges from -35 to 30, and the x-axis ranges from -10 to 10.

Colossal magneto resistance



Y. Tokura and N. Nagaosa
 Science 288, 462 (2000).

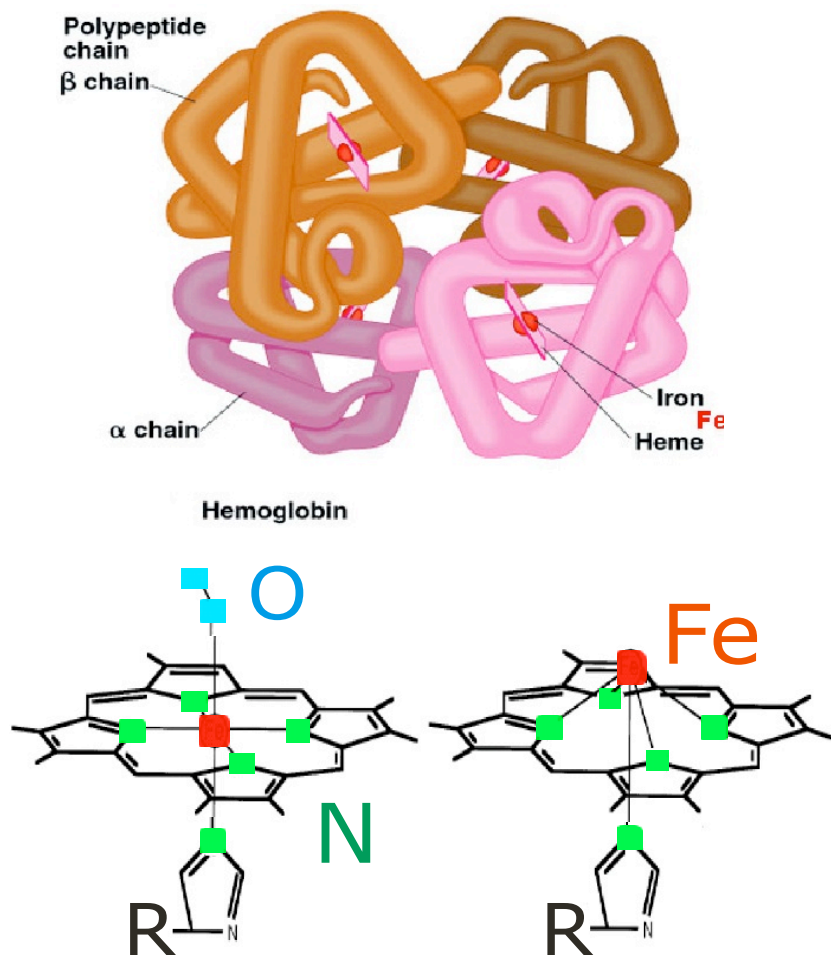
High-Tc Superconductivity



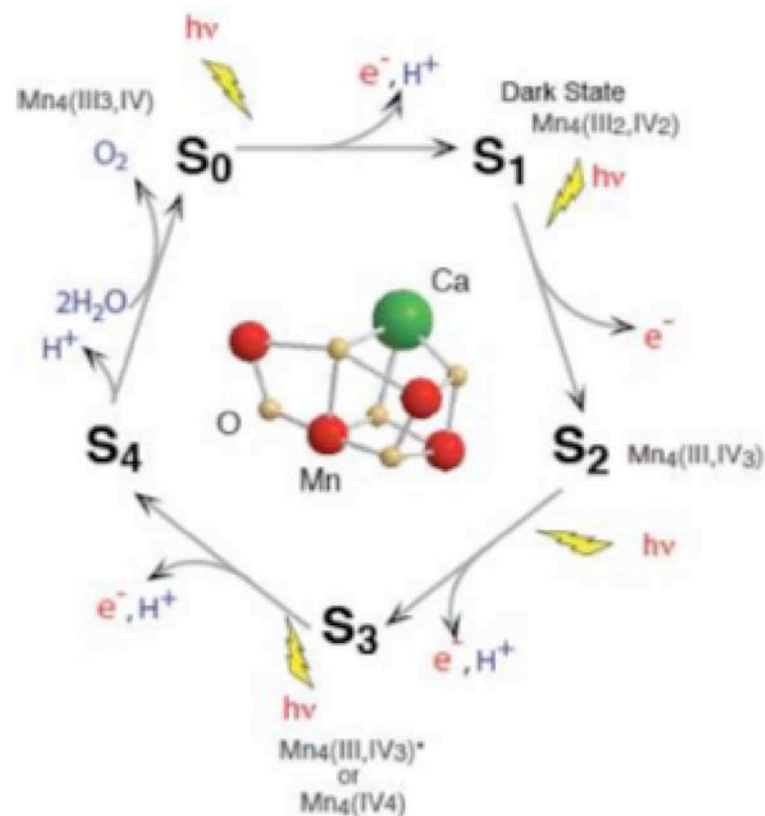
J. Orenstein and A. J. Millis,
 Science 288, 468 (2000).

Transition metal compounds show a large variety of properties

Active centers in enzymes



Fe in Heme

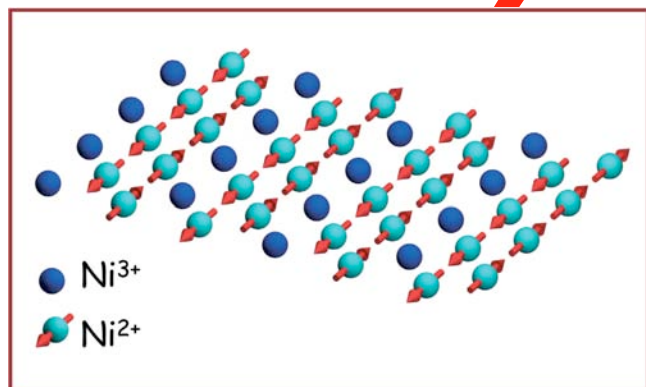


Kok Cycle for water splitting reaction in PSII.

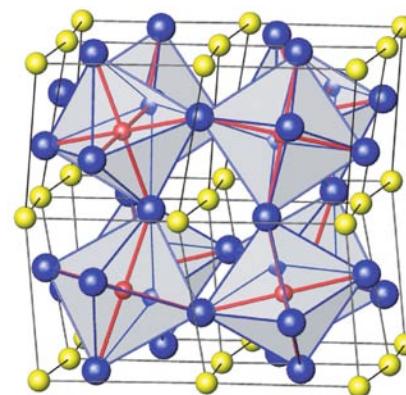
Mn in Photo-system II

Interactions between the charge, lattice, orbitals and spins

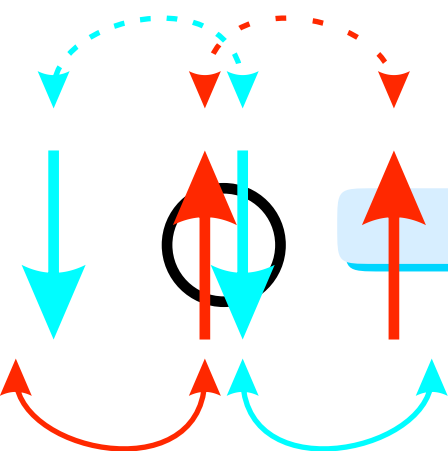
Charge



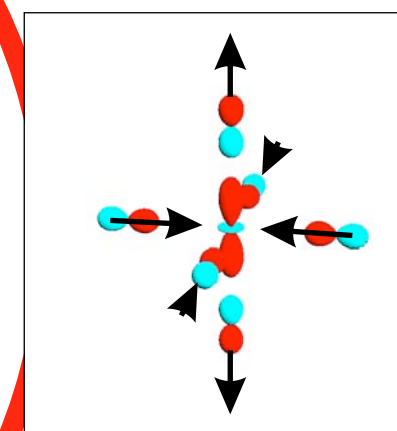
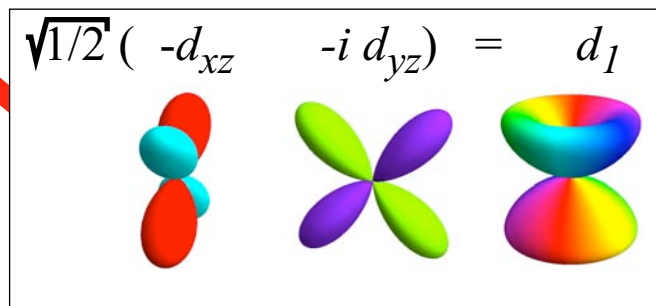
Lattice



Spin



Orbital



Need tools to measure

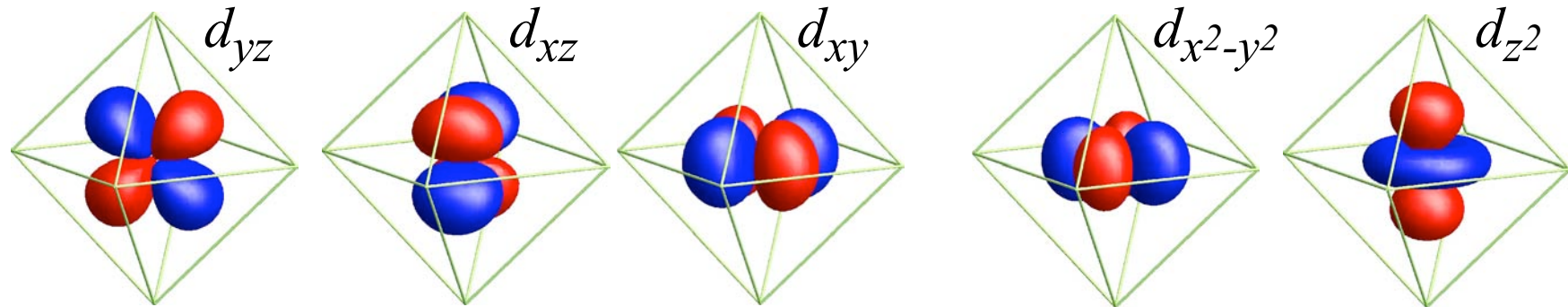
- Spin arrangements (element resolved) of TM multilayers

Need tools to measure

- Spin arrangements (element resolved) of TM multilayers
 - Valence and local symmetry of TM compounds
-

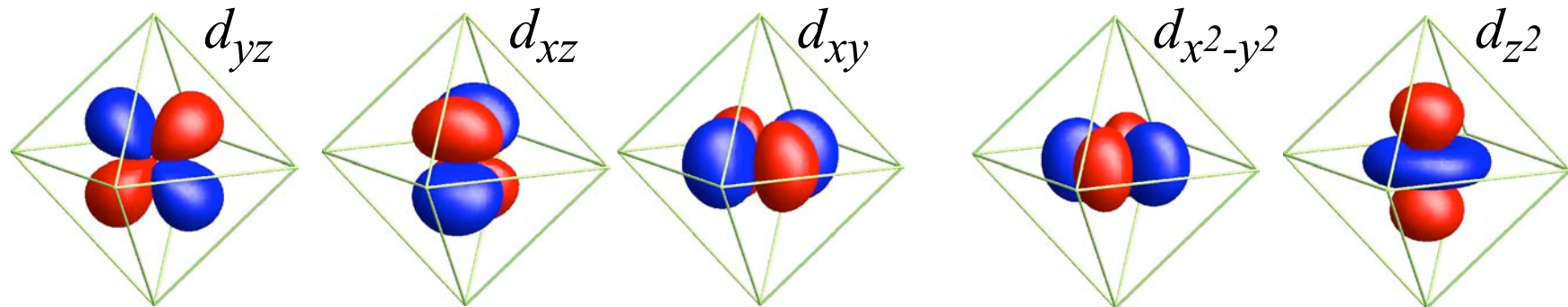
Need tools to measure

- ❑ Spin arrangements (element resolved) of TM multilayers
- ❑ Valence and local symmetry of TM compounds
- ❑ Orbital occupation (The d-shell is 5-fold degenerate)



Need tools to measure

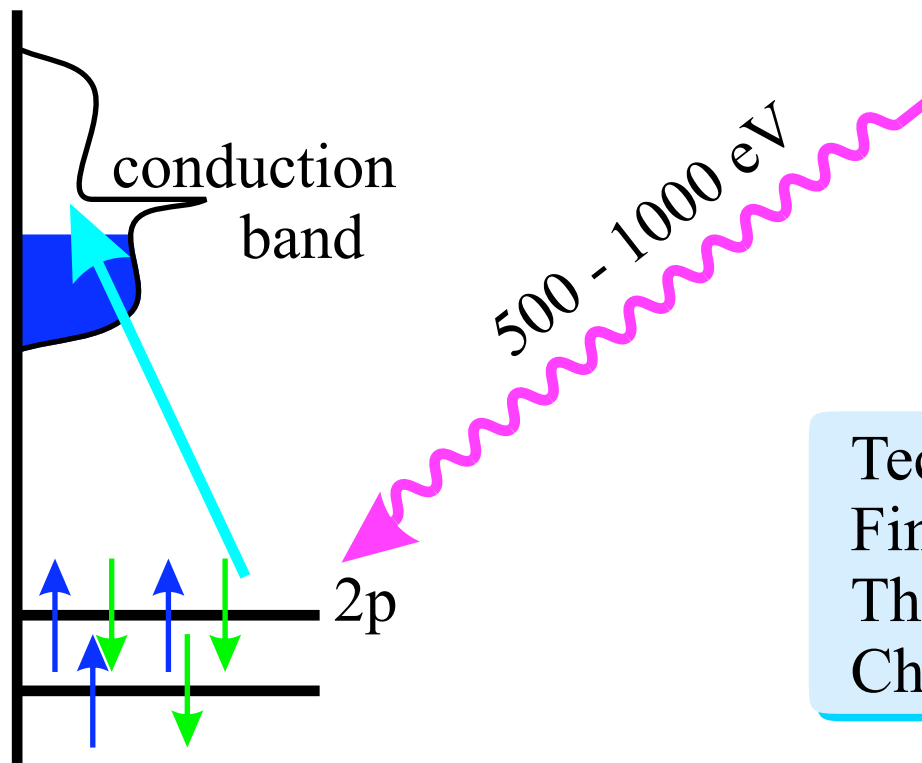
- ❑ Spin arrangements (element resolved) of TM multilayers
- ❑ Valence and local symmetry of TM compounds
- ❑ Orbital occupation (The d-shell is 5-fold degenerate)



- ❑ And excitations (spin-waves - magnons - orbital excitations)

Need tools to measure: Orbital, Spin, Charge

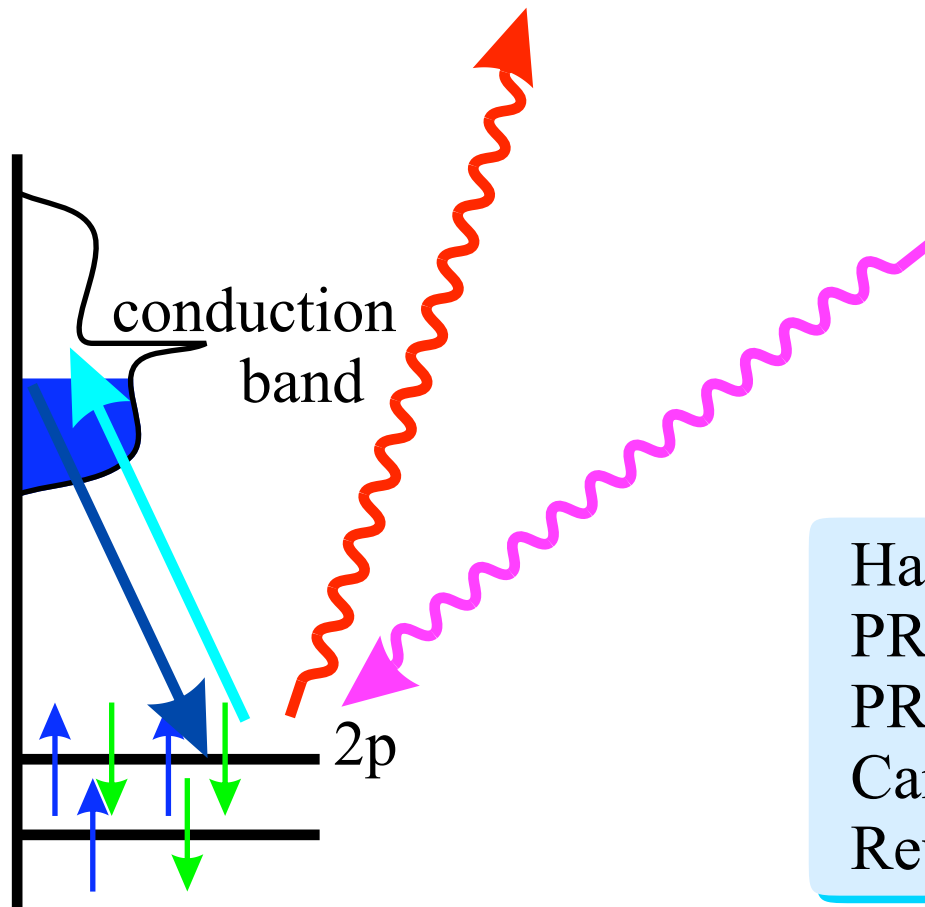
- X-ray Absorption Spectroscopy (XAS) at the TM $2p$ to $3d$ edge



Technique developed in late 1980's:
Fink, Sawatzky, Fuggle
Thole, van der Laan
Chen, Sette

Need tools to measure: Orbital, Spin, Charge

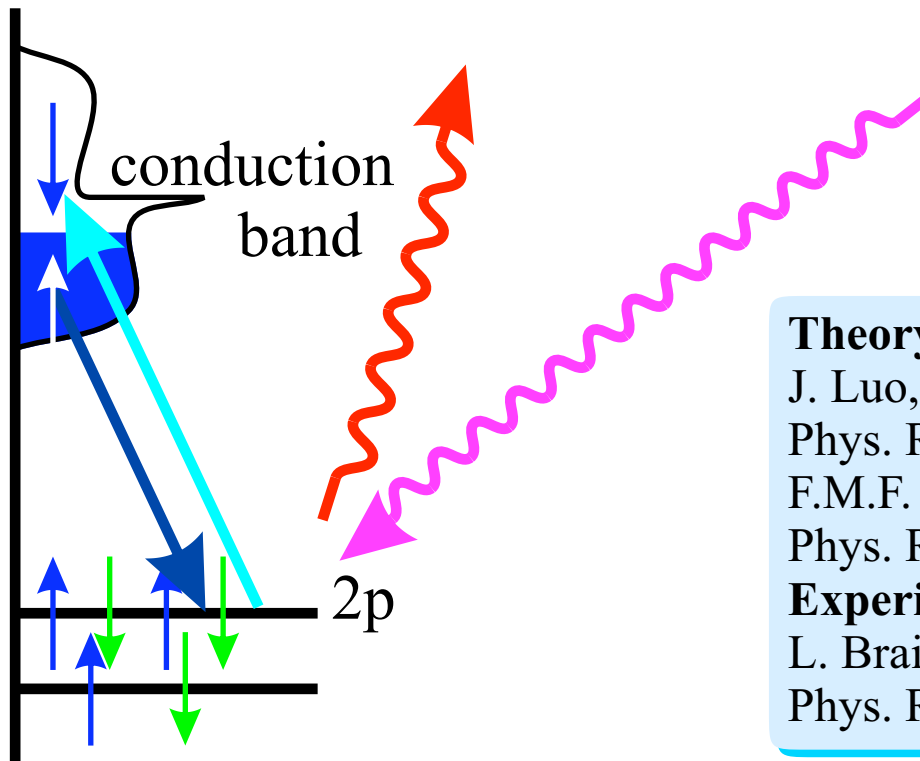
- ❑ X-ray Absorption Spectroscopy (XAS) at the TM $2p$ to $3d$ edge
- ❑ Resonant X-ray Diffraction (RXD)



Hannon, Trammel, Bloom, Gibbs
PRL **61**, 1245 (1988)
PRB **43**, 5663 (1991)
Carra, Thole
Rev. Mod. Phys. **66**, 1509 (1994).

Need tools to measure: Orbital, Spin, Charge

- ❑ X-ray Absorption Spectroscopy (XAS) at the TM $2p$ to $3d$ edge
- ❑ Resonant X-ray Diffraction (RXD)
- ❑ Resonant Inelastic X-ray Scattering (RIXS)



Theory prediction of magnetic excitations

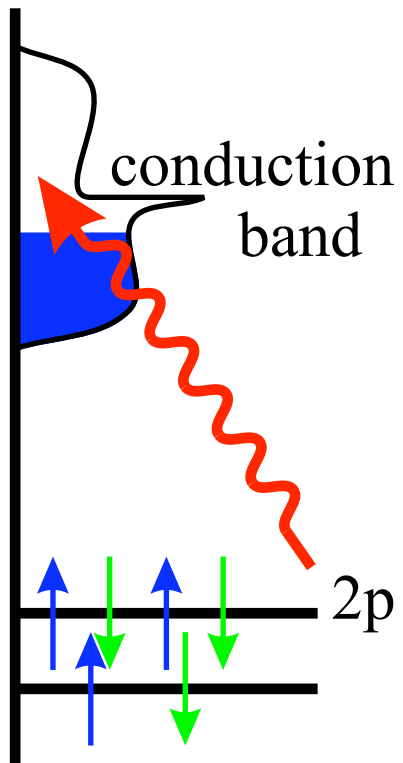
J. Luo, G.T. Trammell, and J.P. Hannon,
Phys. Rev. Lett. **71**, 287 (1993).

F.M.F. de Groot, P. Kuiper, and G.A. Sawatzky,
Phys. Rev. B **57**, 14584 (1998).

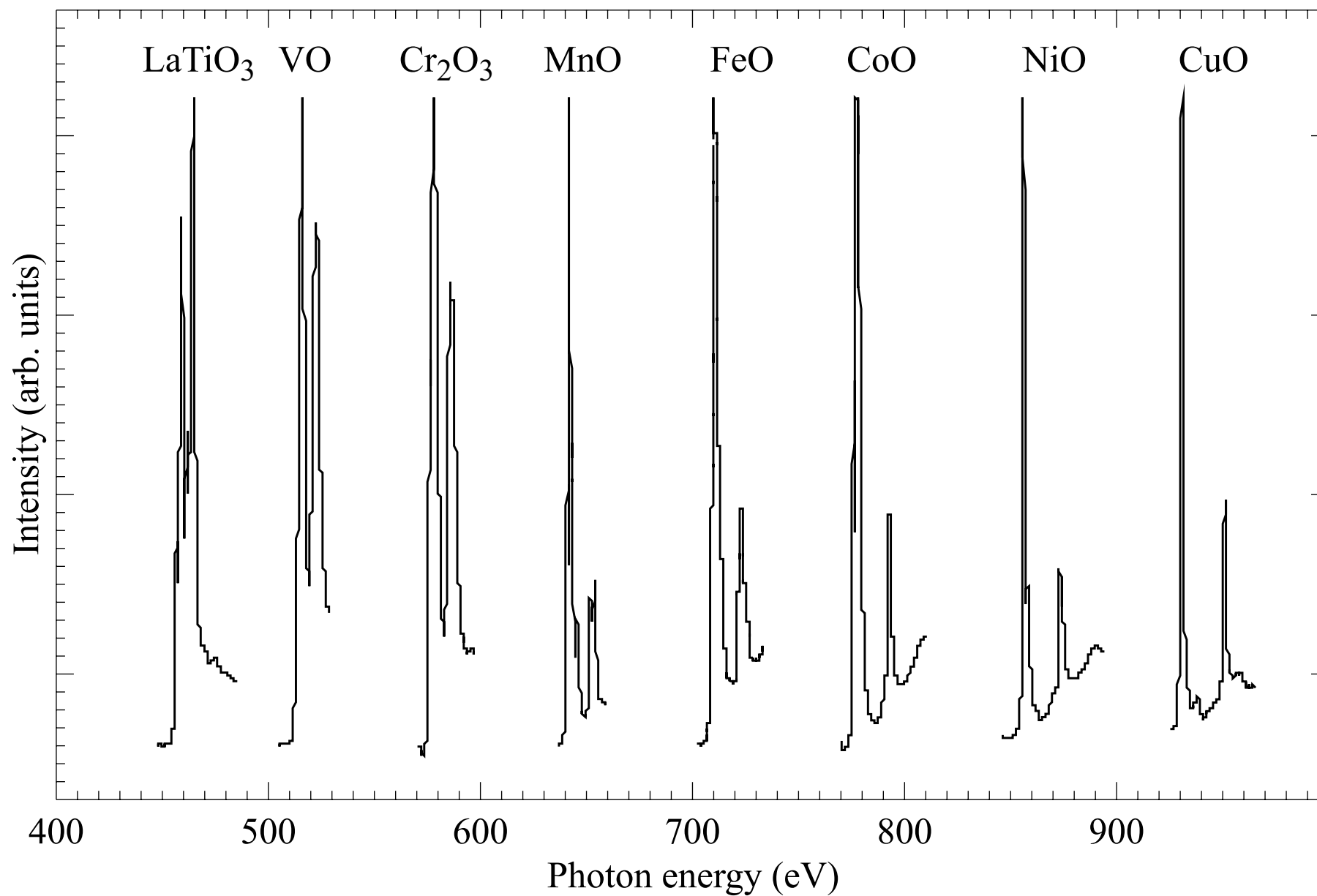
Experimental realization

L. Braicovich, G. Ghiringhelli, N. B. Brookes *et al.*
Phys. Rev. Lett. **102**, 167401 (2009).

X-ray Absorption Spectroscopy

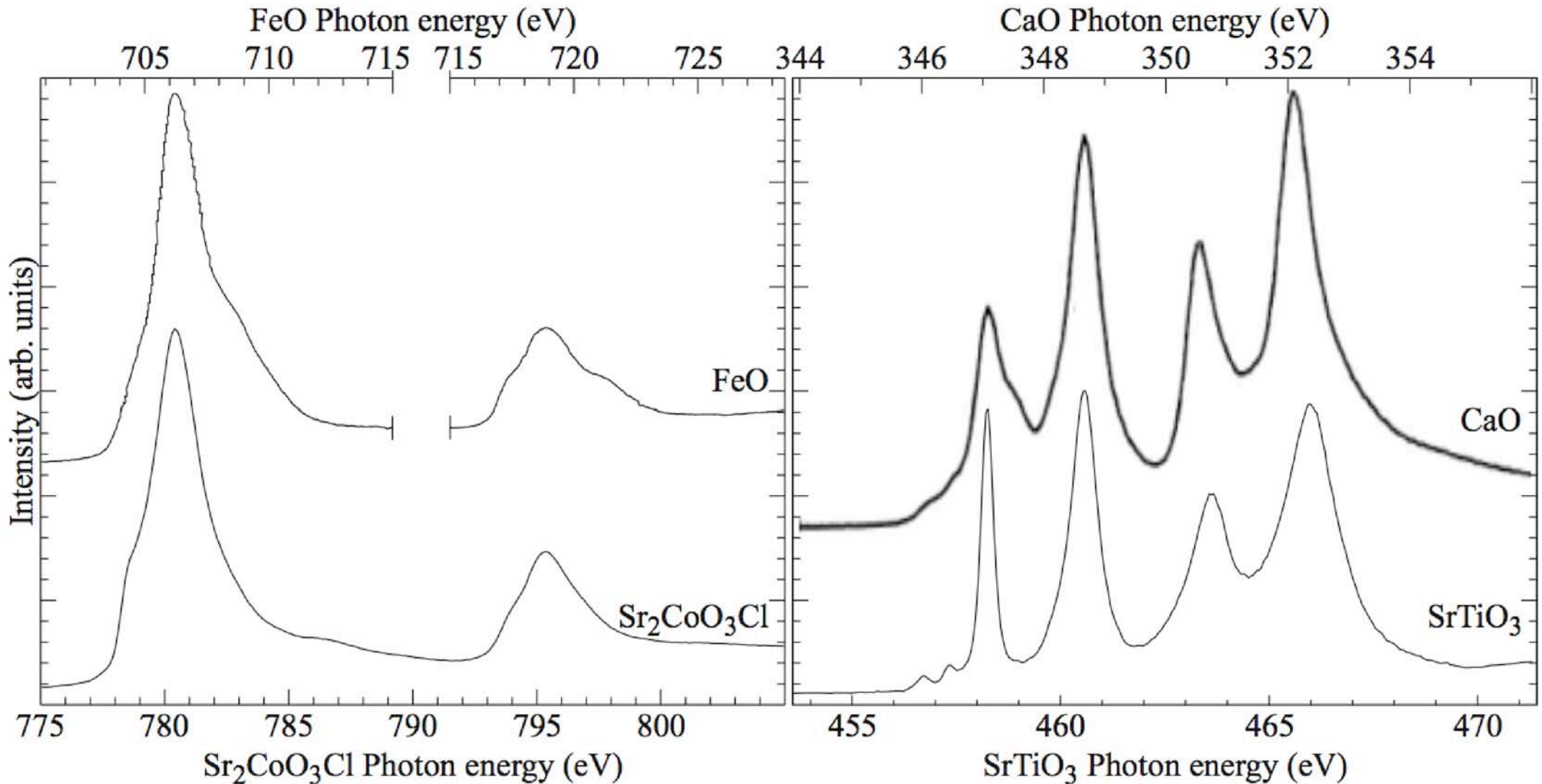


X-ray Absorption Spectroscopy: Some observations



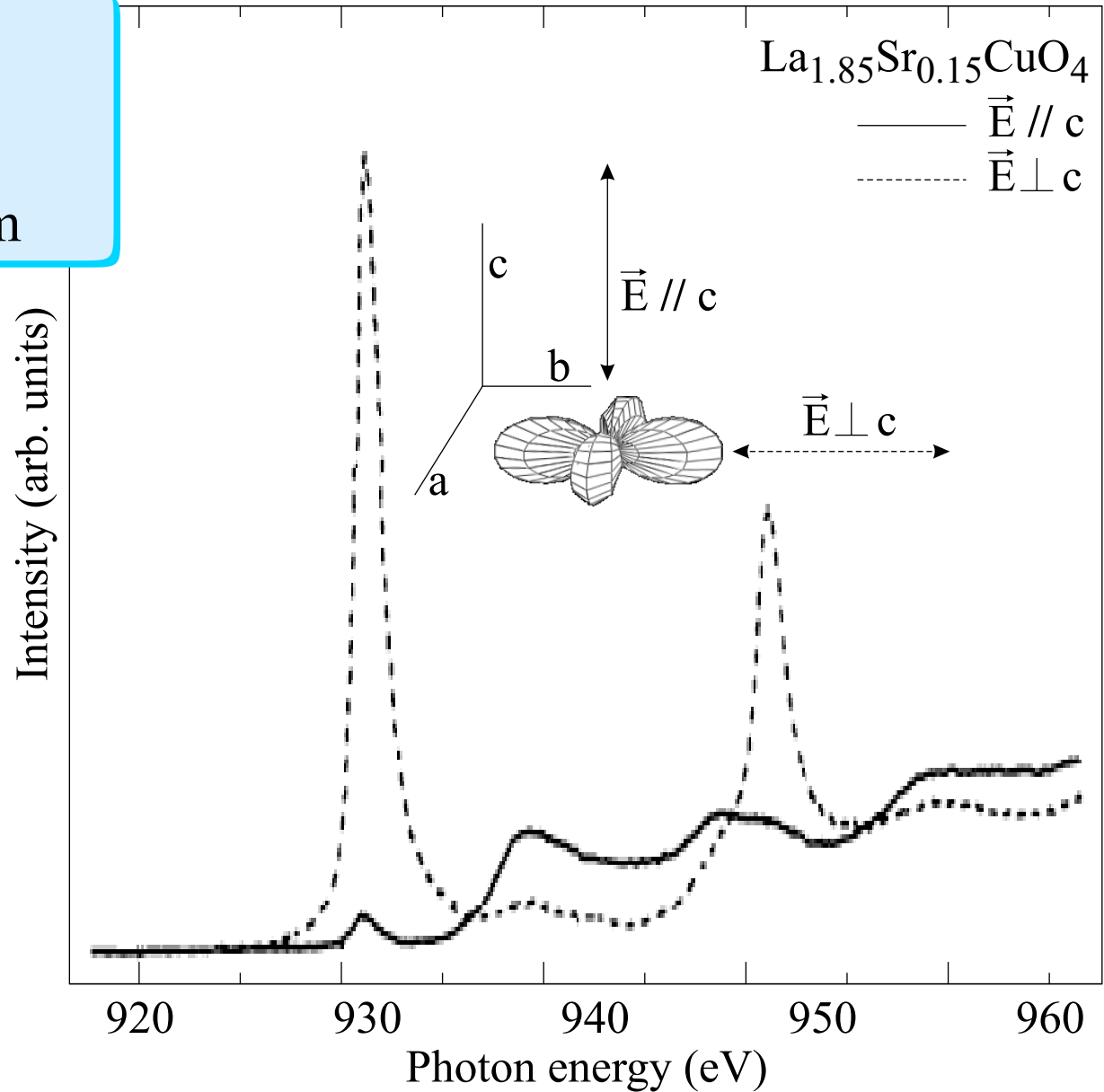
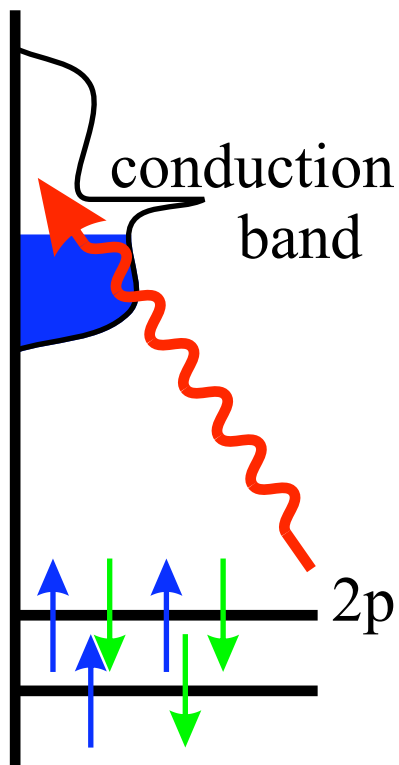
X-ray Absorption Spectroscopy: Some observations

Valence dependence - ground state symmetry determines the line shape



X-ray Absorption Spectroscopy: Some observations

Polarized XAS
Sensitive to
Orbital occupation
natural linear dichroism



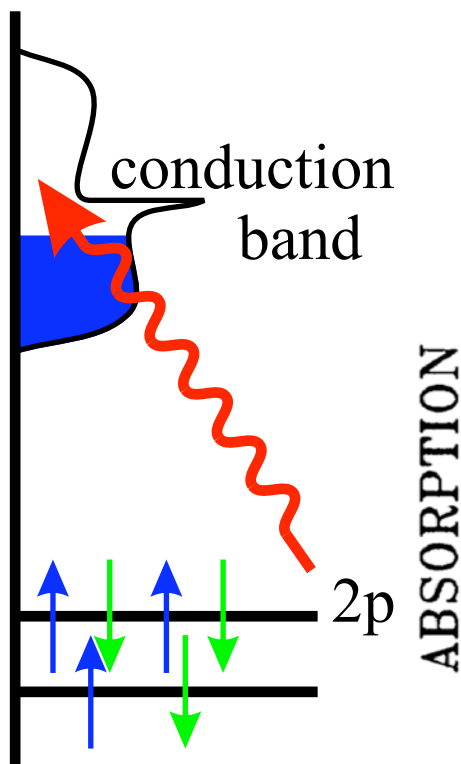
X-ray Absorption Spectroscopy: Some observations

Polarized XAS
Sensitive to
Magnetic moments
magnetic circular dichroism

VOLUME 75, NUMBER 1

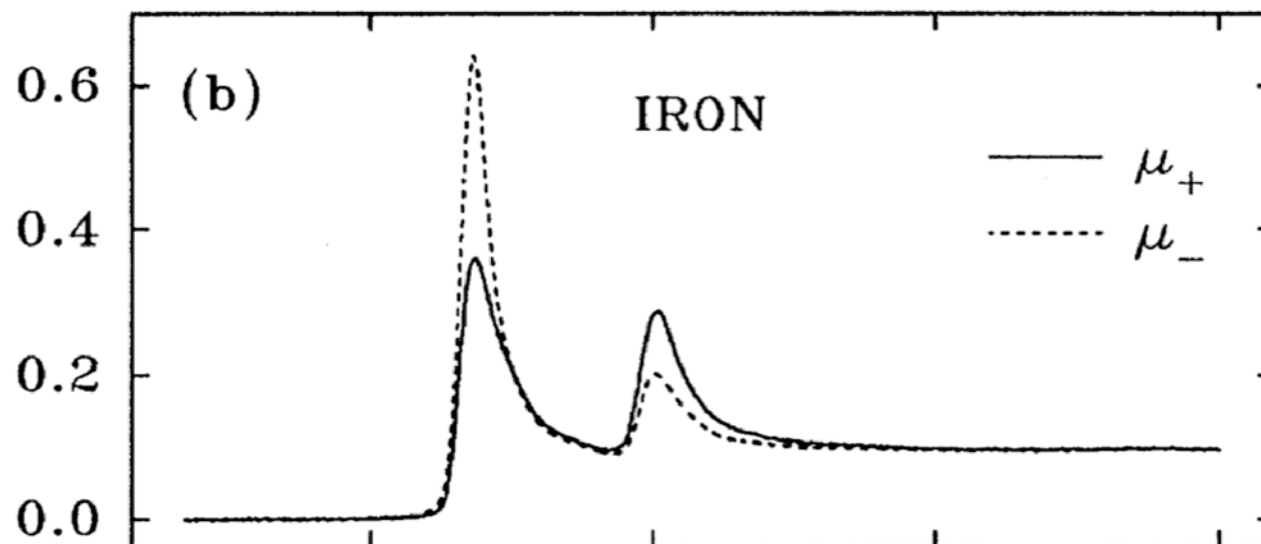
PHYSICAL REVIEW LETTERS

3 JULY 1995



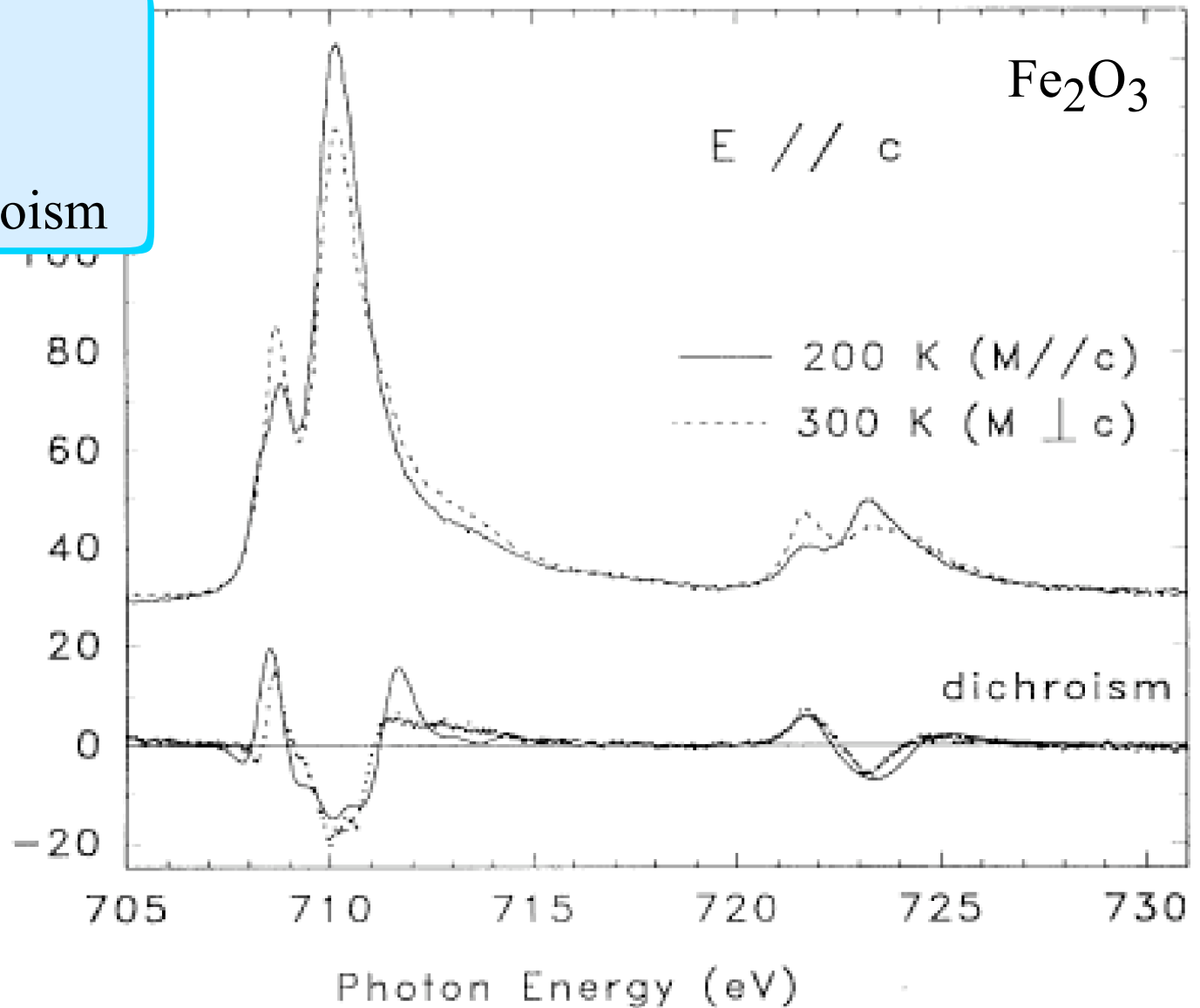
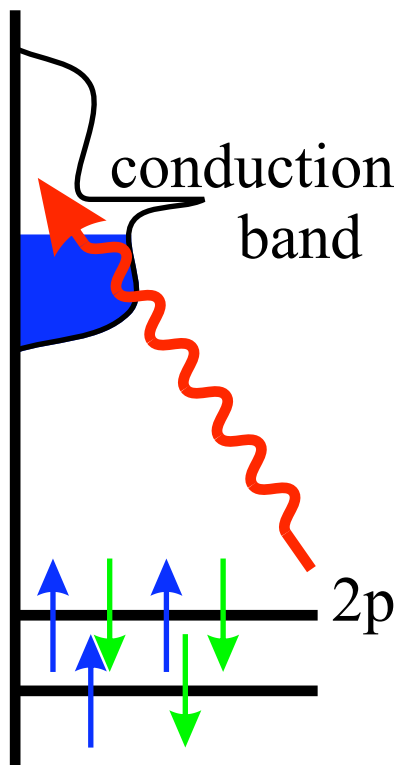
Experimental Confirmation of the X-Ray Magnetic Circular Dichroism Sum Rules for Iron and Cobalt

C. T. Chen,¹ Y. U. Idzerda,² H.-J. Lin,^{1,*} N. V. Smith,^{1,†} G. Meigs,¹ E. Chaban,¹
G. H. Ho,^{3,*} E. Pellegrin,¹ and F. Sette^{1,‡}



X-ray Absorption Spectroscopy: Some observations

Polarized XAS
Sensitive to
Spin orientation
magnetic linear dichroism



X-ray Absorption Spectroscopy: A band picture

RAPID COMMUNICATIONS

PHYSICAL REVIEW B

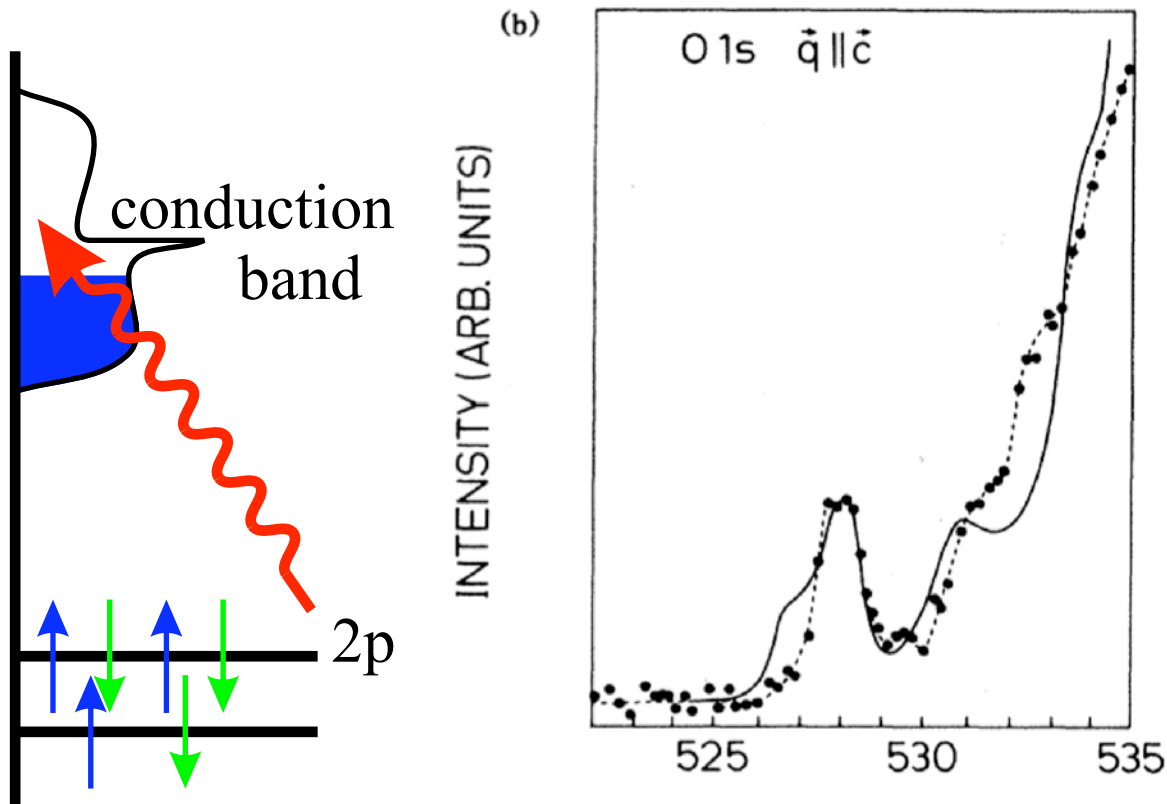
VOLUME 40, NUMBER 1

1 JULY 1989

X-ray absorption of $\text{YBa}_2\text{Cu}_3\text{O}_7$: A band picture

J. Zaanen, M. Alouani, and O. Jepsen

Max-Planck-Institut für Festkörperforschung, D-7000 Stuttgart 80, Heisenbergstrasse 1, Federal Republic of Germany



X-ray Absorption Spectroscopy: A band picture

RAPID COMMUNICATIONS

PHYSICAL REVIEW B

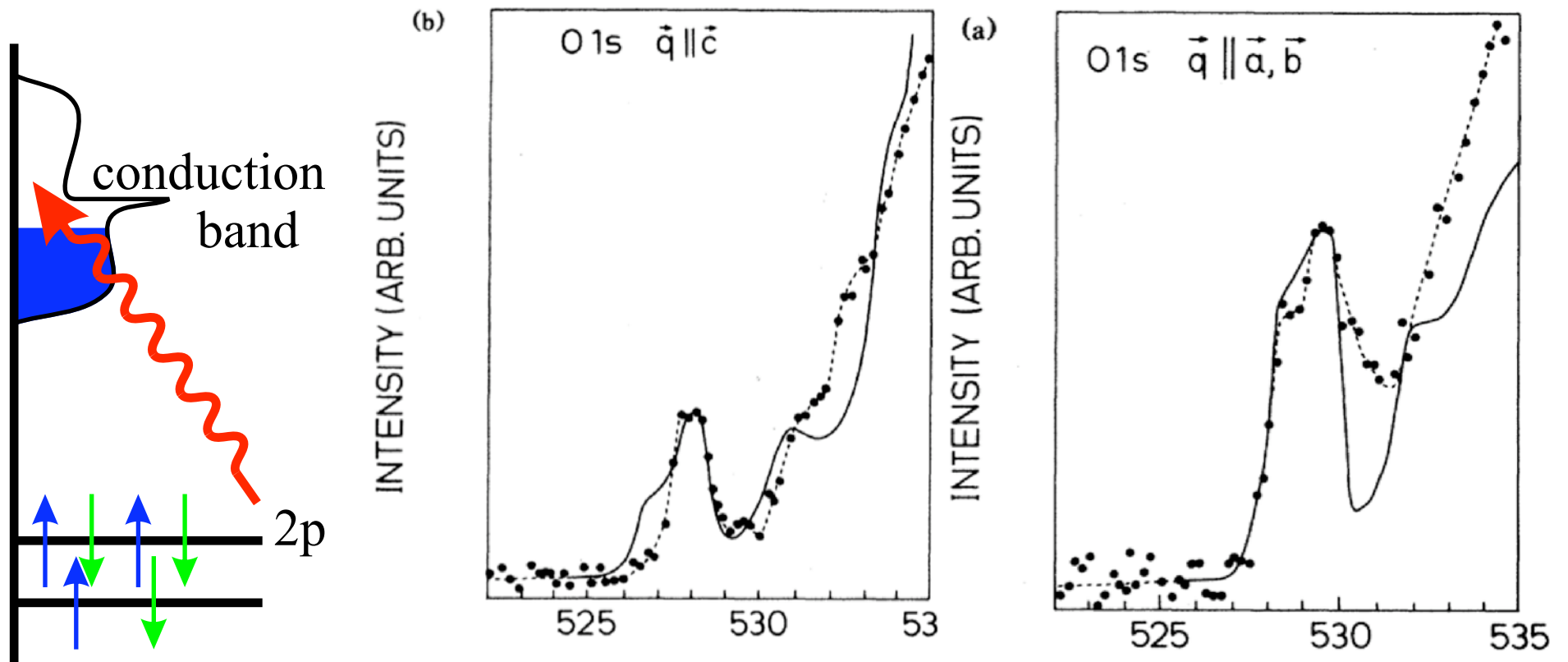
VOLUME 40, NUMBER 1

1 JULY 1989

X-ray absorption of $\text{YBa}_2\text{Cu}_3\text{O}_7$: A band picture

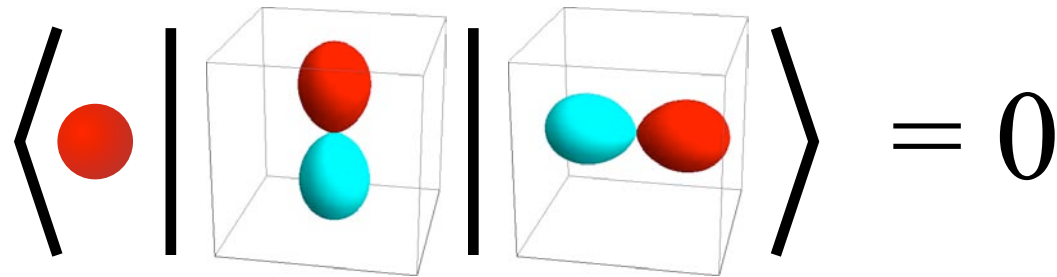
J. Zaanen, M. Alouani, and O. Jepsen

Max-Planck-Institut für Festkörperforschung, D-7000 Stuttgart 80, Heisenbergstrasse 1, Federal Republic of Germany

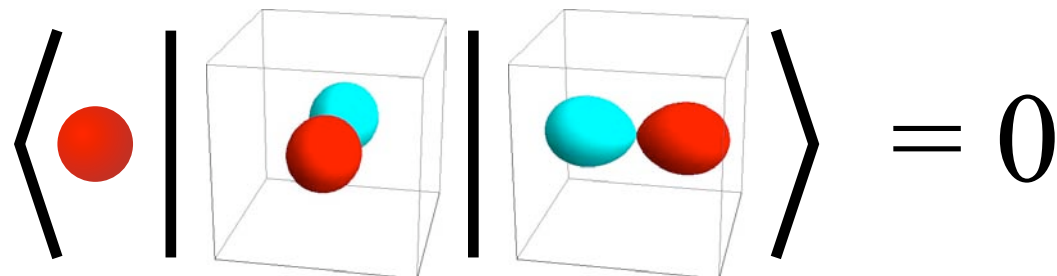


X-ray Absorption Spectroscopy: One electron selection rules

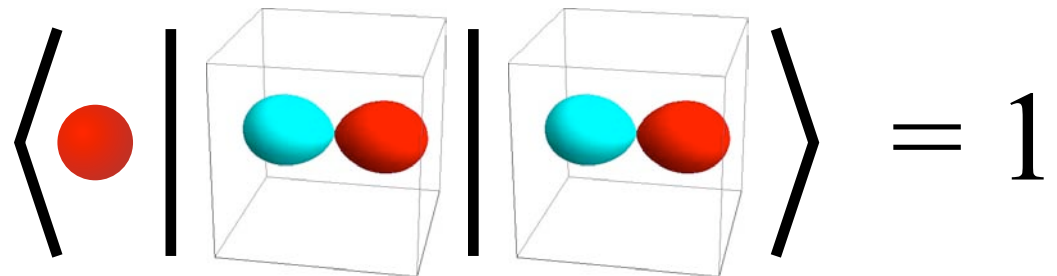
$$\mu_q(h\nu) = \sum_{\Phi_f} |\langle \Phi_f | r C_q^{(1)} | \Phi_i \rangle|^2 \delta(E_i - E_f + h\nu)$$



$\langle \text{red sphere} | \text{box with red and cyan spheres} | \text{box with red and cyan spheres} \rangle = 0$



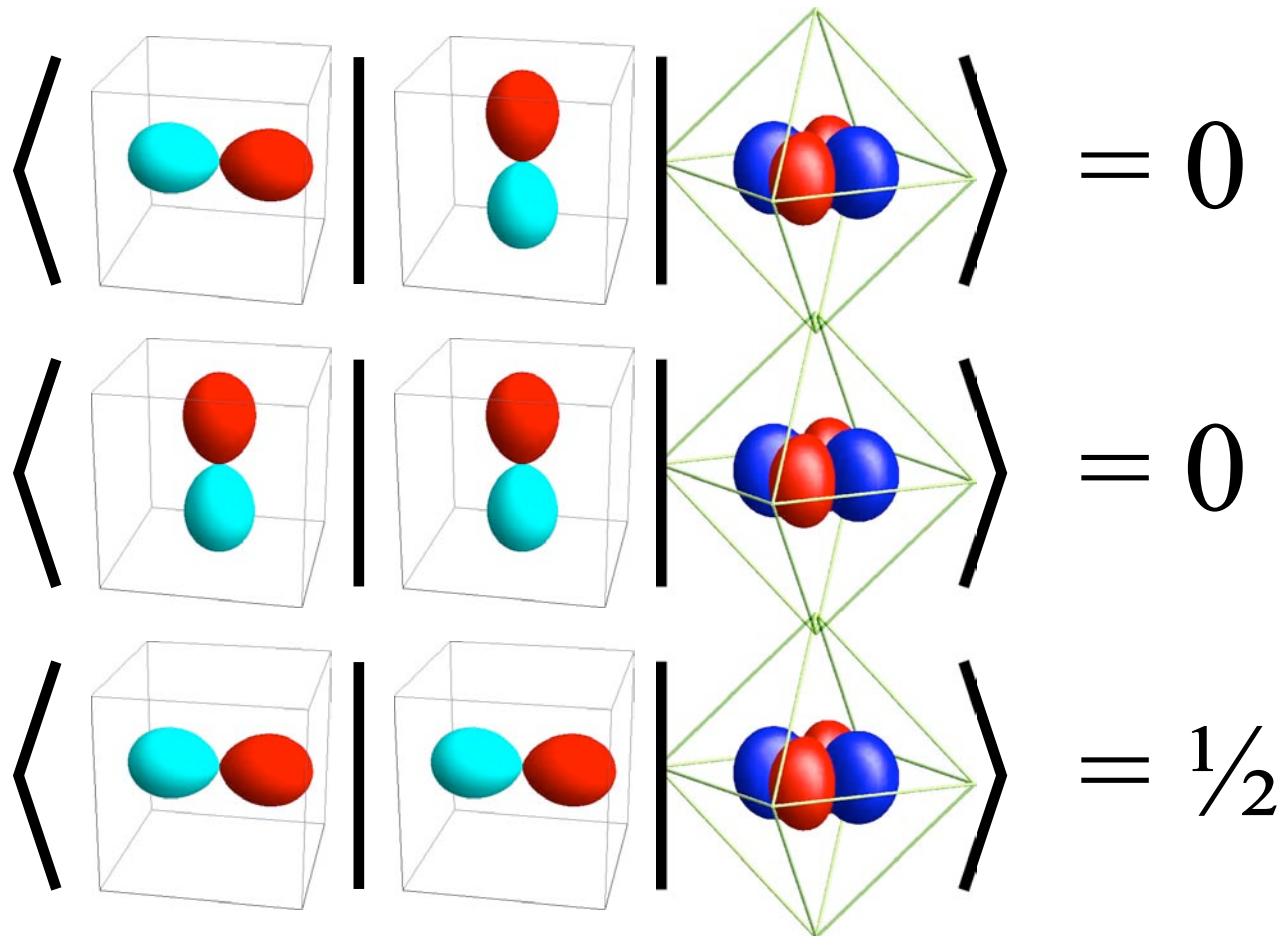
$\langle \text{red sphere} | \text{box with red and cyan spheres} | \text{box with red and cyan spheres} \rangle = 0$



$\langle \text{red sphere} | \text{box with red and cyan spheres} | \text{box with red and cyan spheres} \rangle = 1$

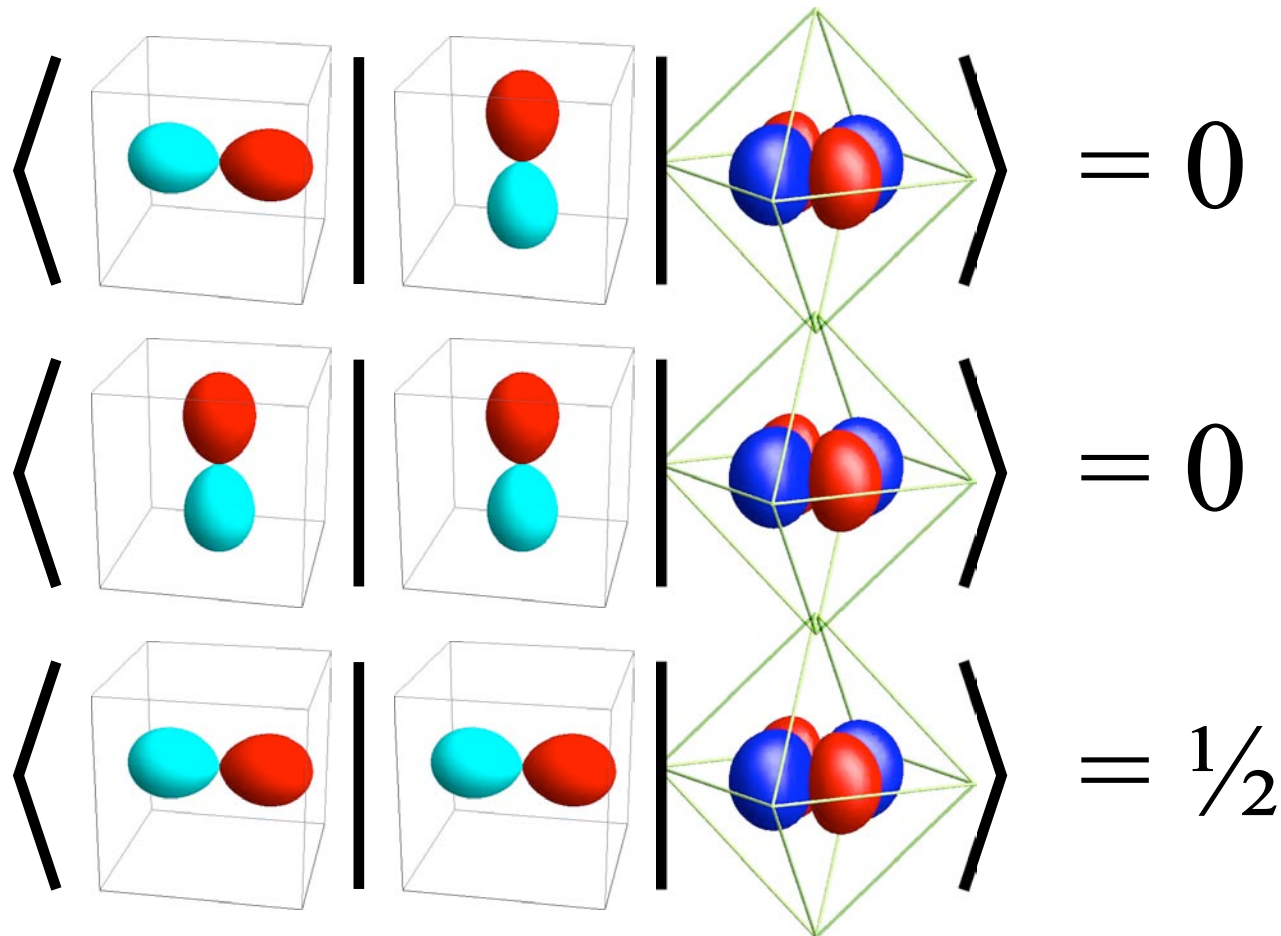
X-ray Absorption Spectroscopy: One electron selection rules

$$\mu_q(h\nu) = \sum_{\Phi_f} |\langle \Phi_f | r C_q^{(1)} | \Phi_i \rangle|^2 \delta(E_i - E_f + h\nu)$$

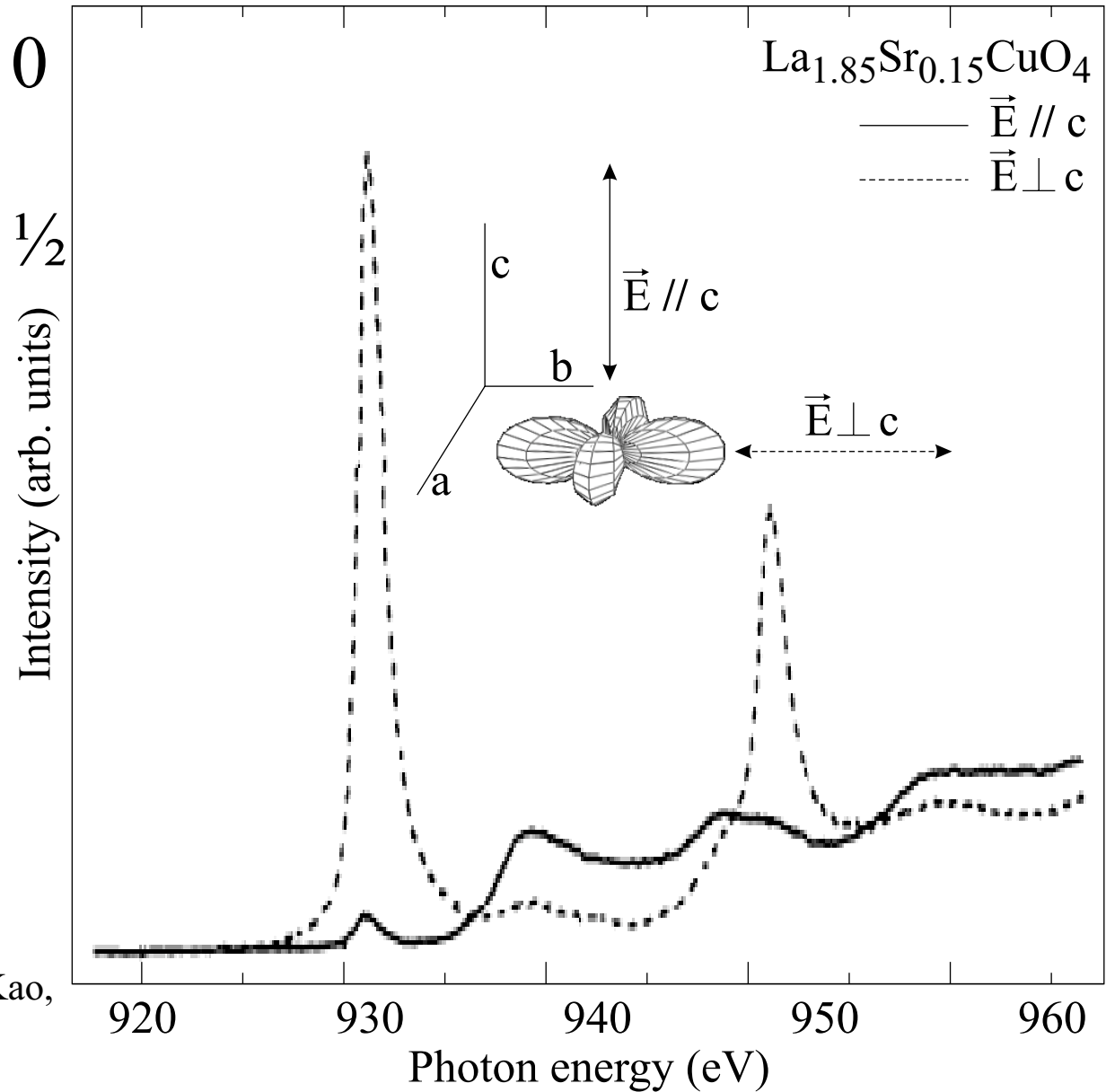
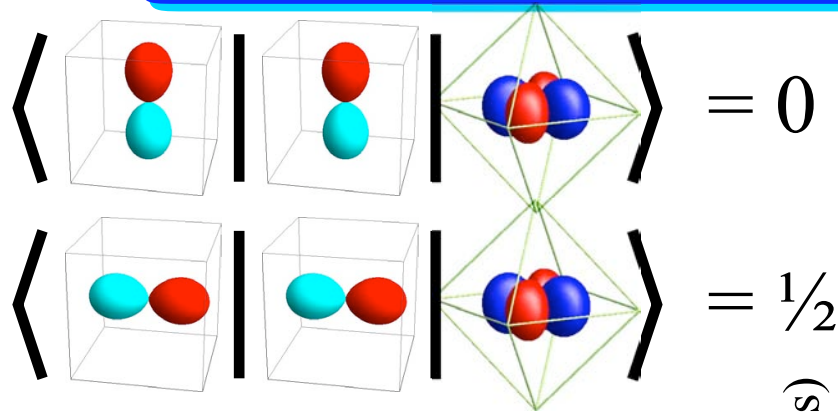


X-ray Absorption Spectroscopy: One electron selection rules

$$\mu_q(h\nu) = \sum_{\Phi_f} |\langle \Phi_f | r C_q^{(1)} | \Phi_i \rangle|^2 \delta(E_i - E_f + h\nu)$$

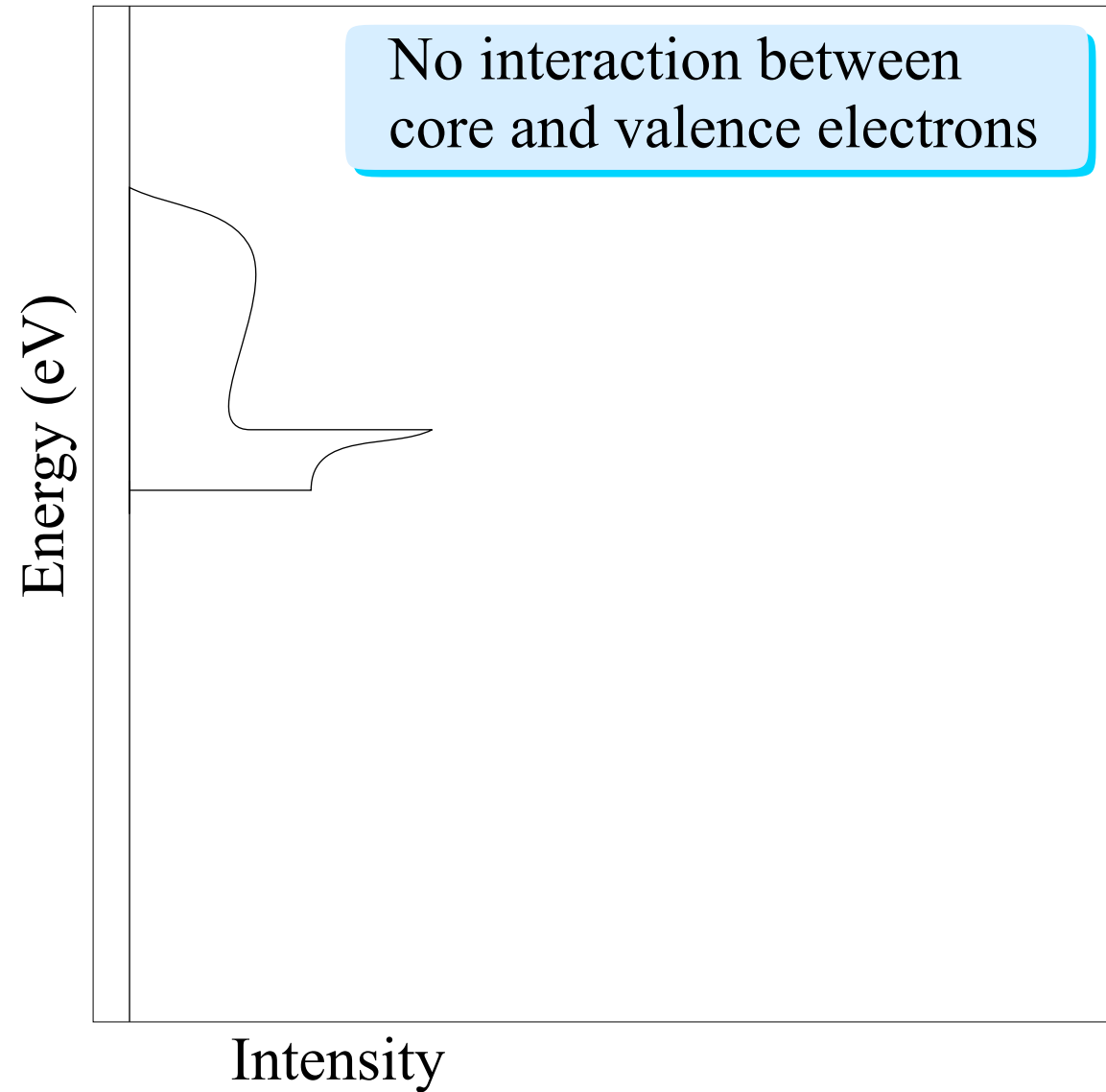
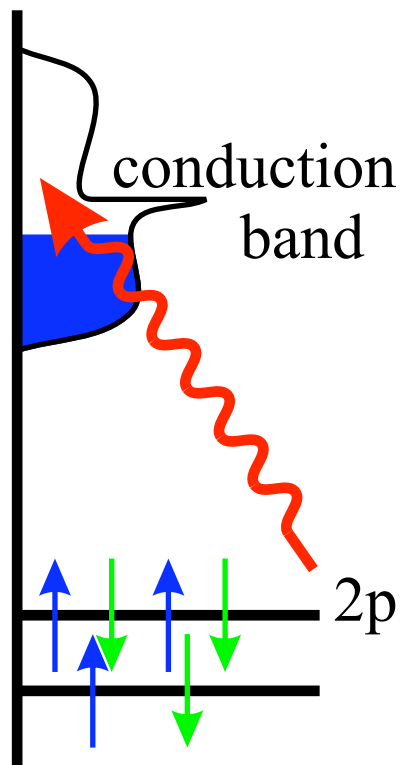


X-ray Absorption Spectroscopy: One electron selection rules

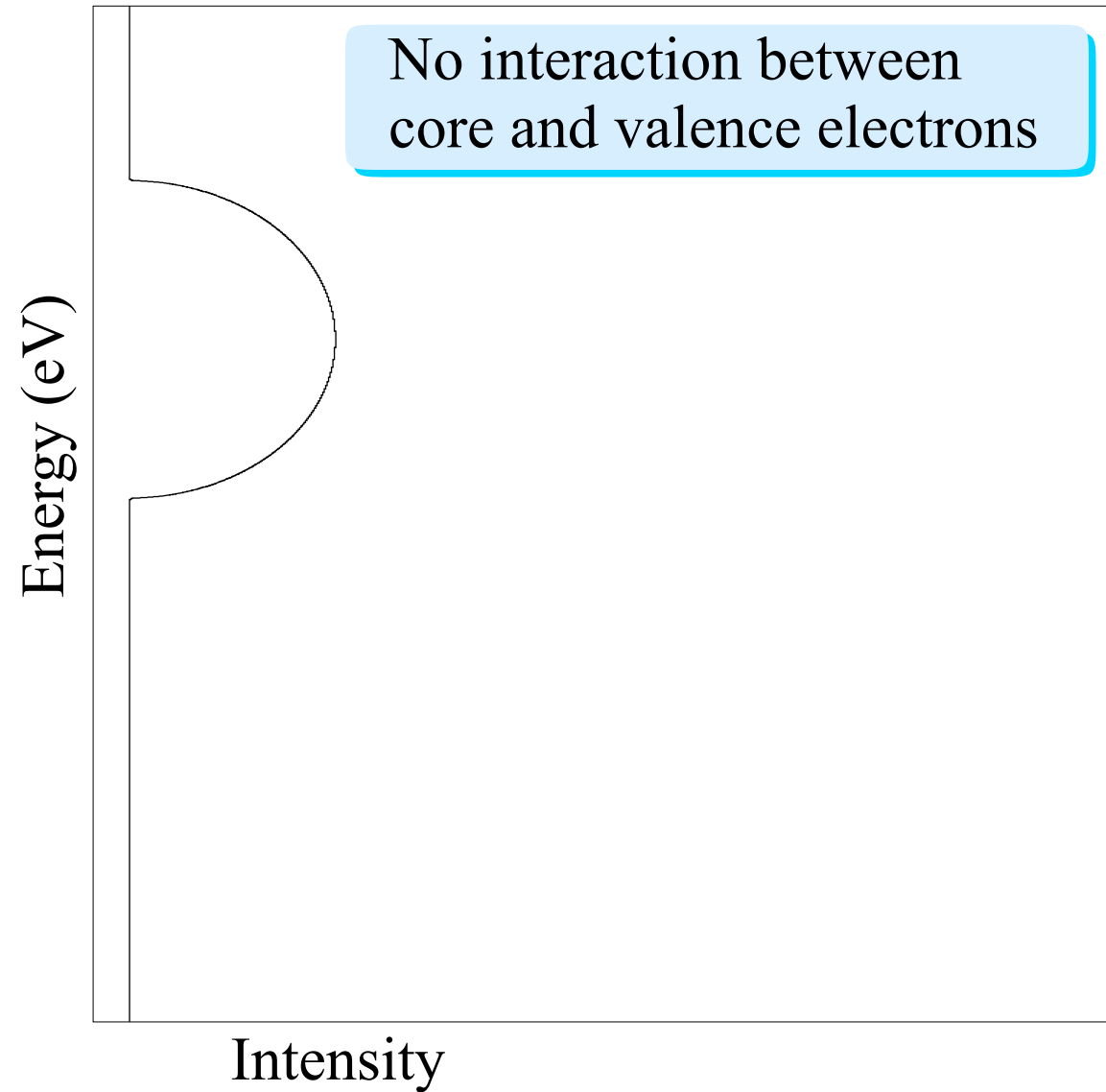
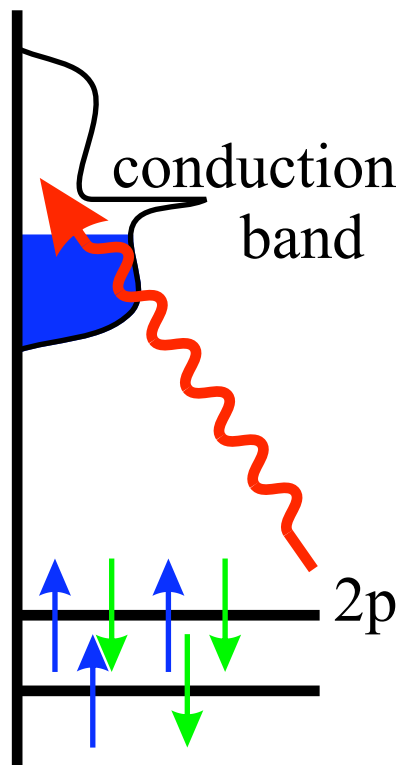


C. T. Chen, L. H. Tjeng, J. Kwo, H. L. Kao,
 P. Rudolf, F. Sette, and R. M. Fleming,
 Phys. Rev. Lett. 68, 2543 (1992).

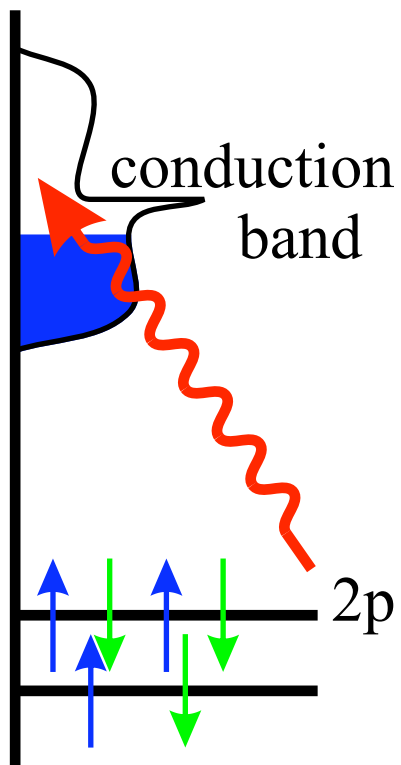
X-ray Absorption Spectroscopy: Excitons



X-ray Absorption Spectroscopy: Excitons



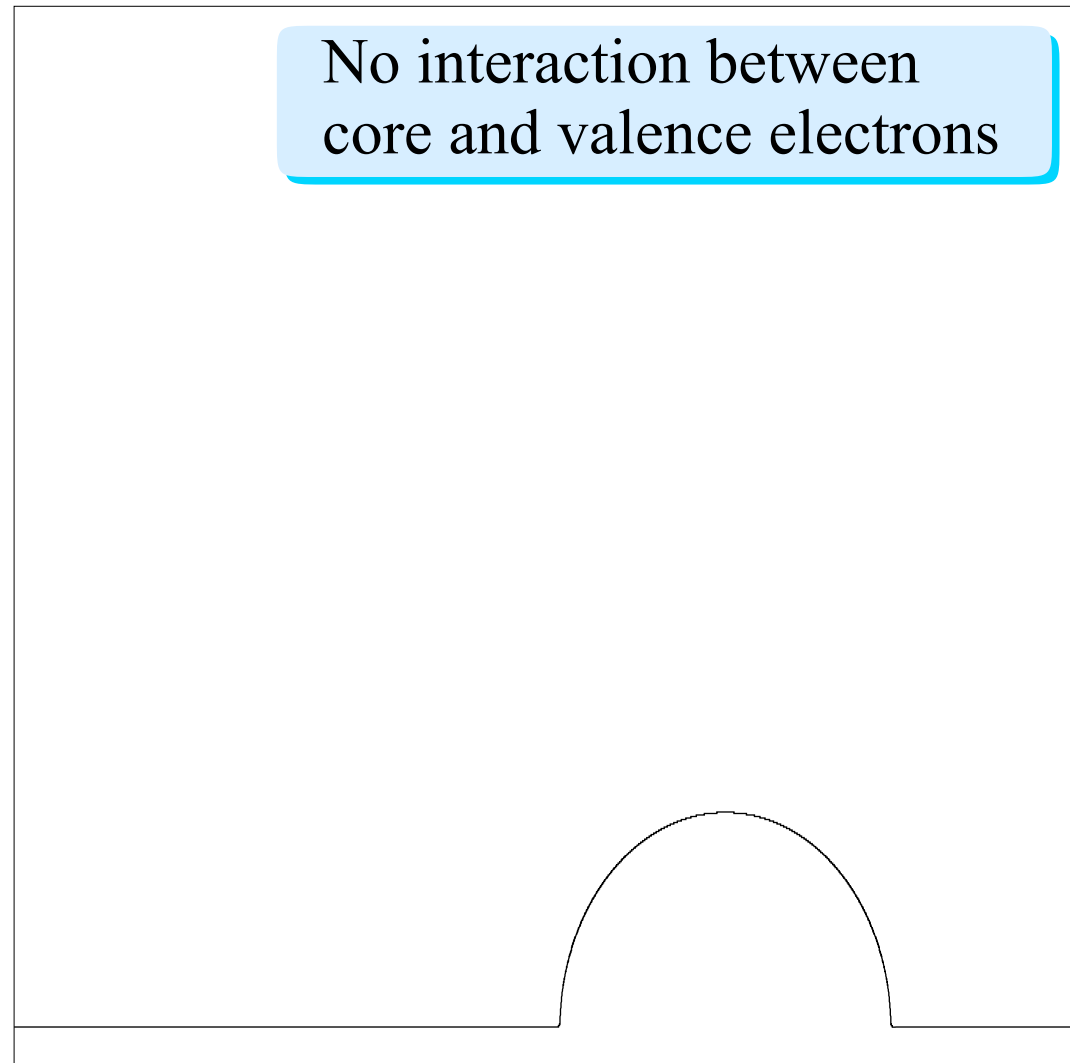
X-ray Absorption Spectroscopy: Excitons



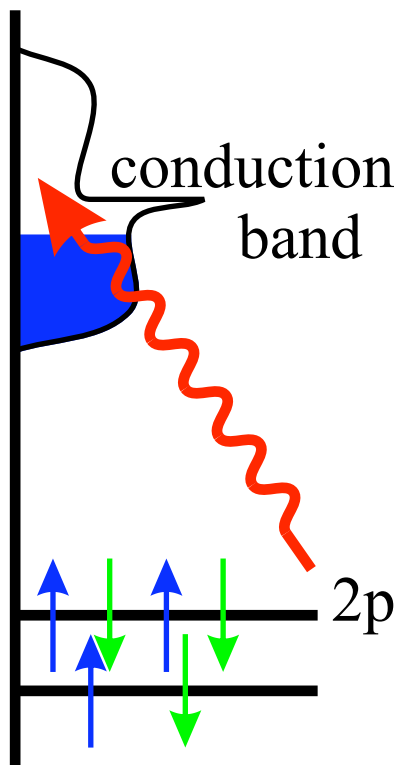
Intensity

No interaction between
core and valence electrons

Energy (eV)



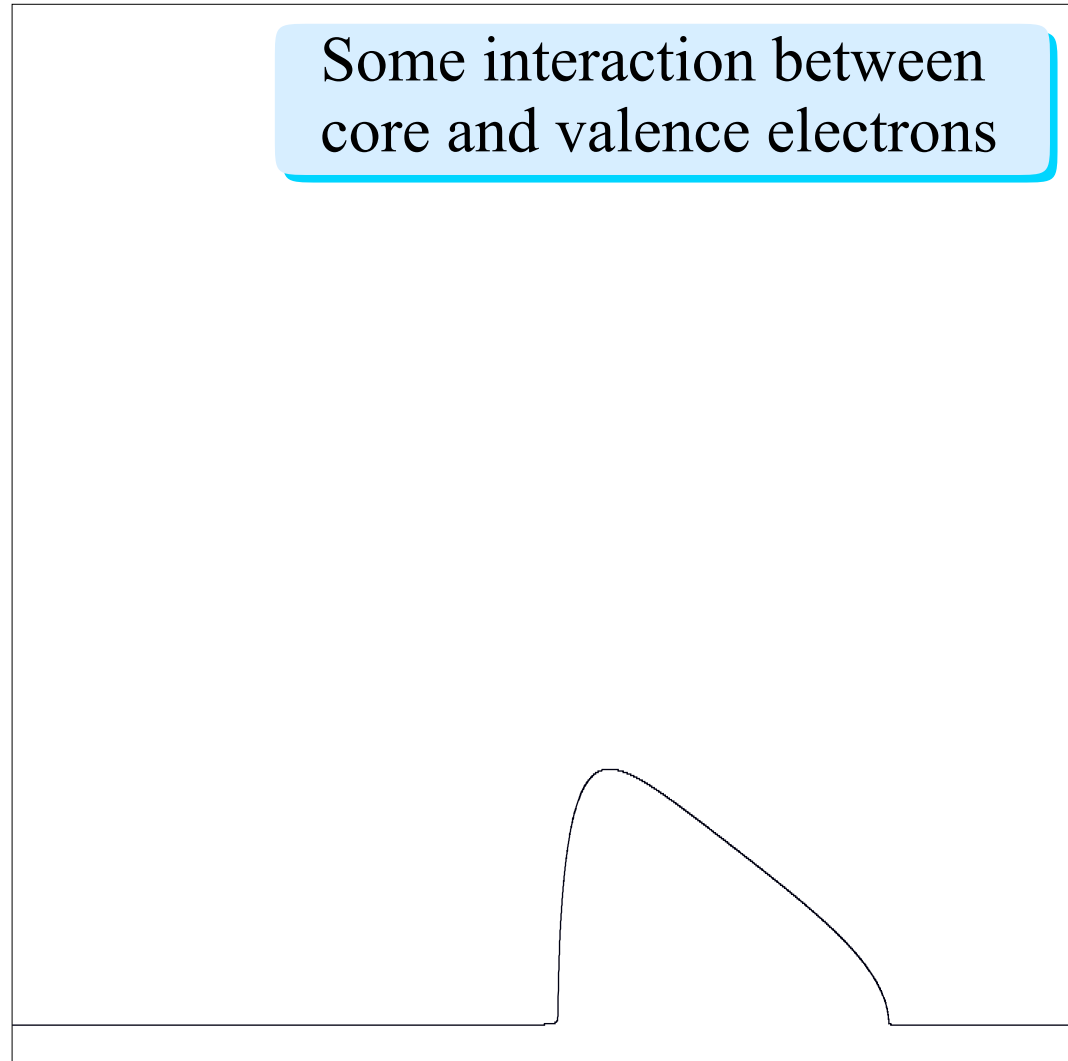
X-ray Absorption Spectroscopy: Excitons



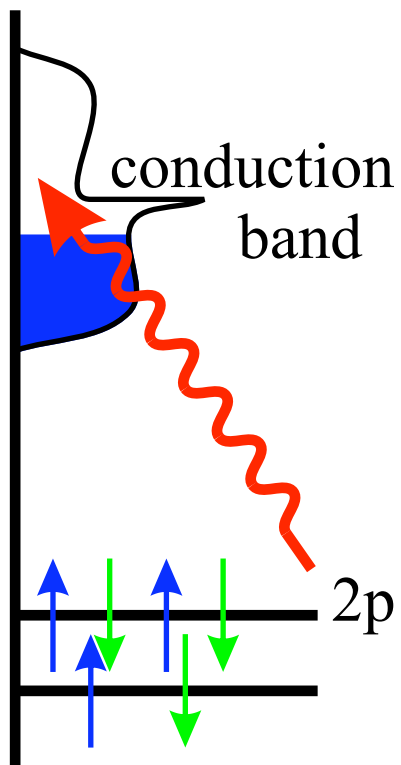
Some interaction between core and valence electrons

Intensity

Energy (eV)



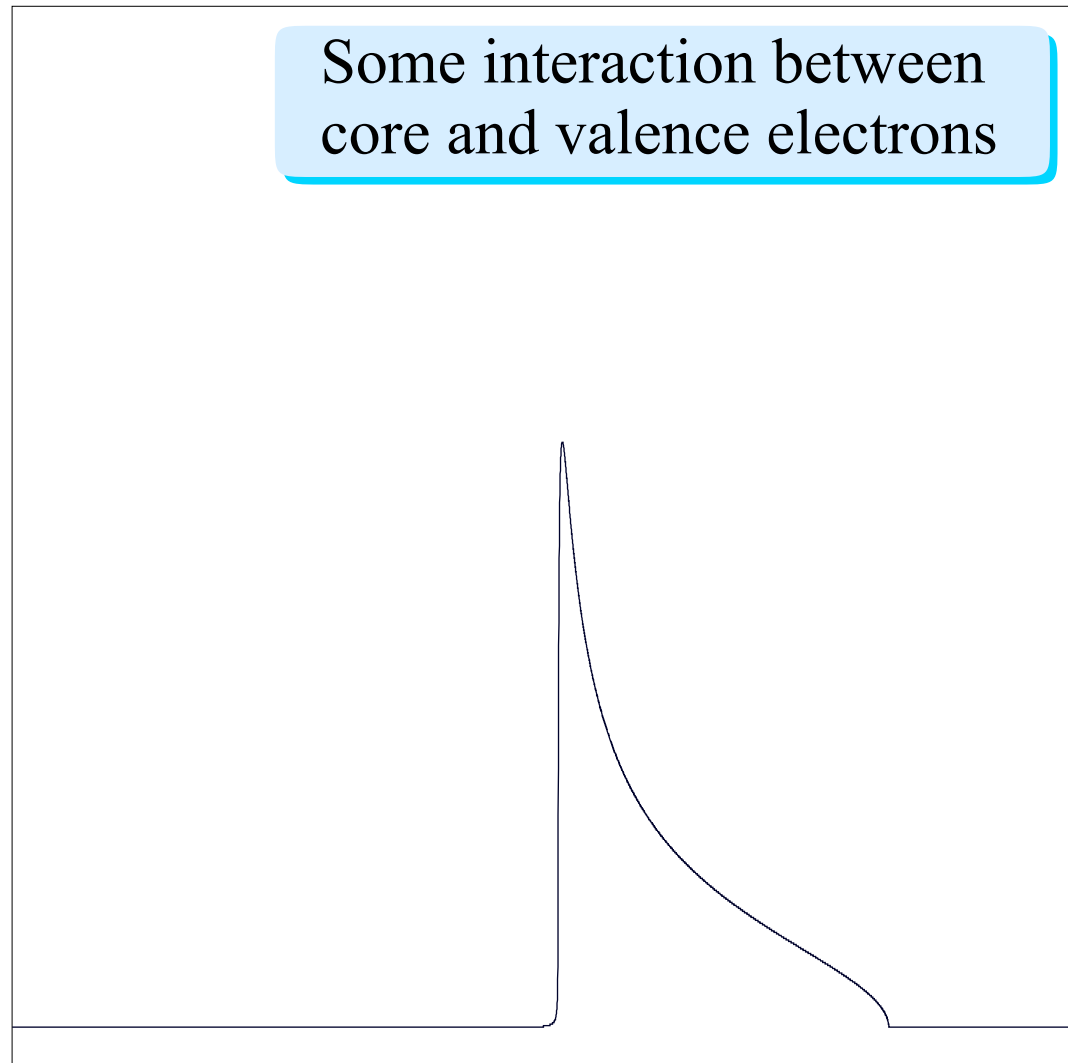
X-ray Absorption Spectroscopy: Excitons



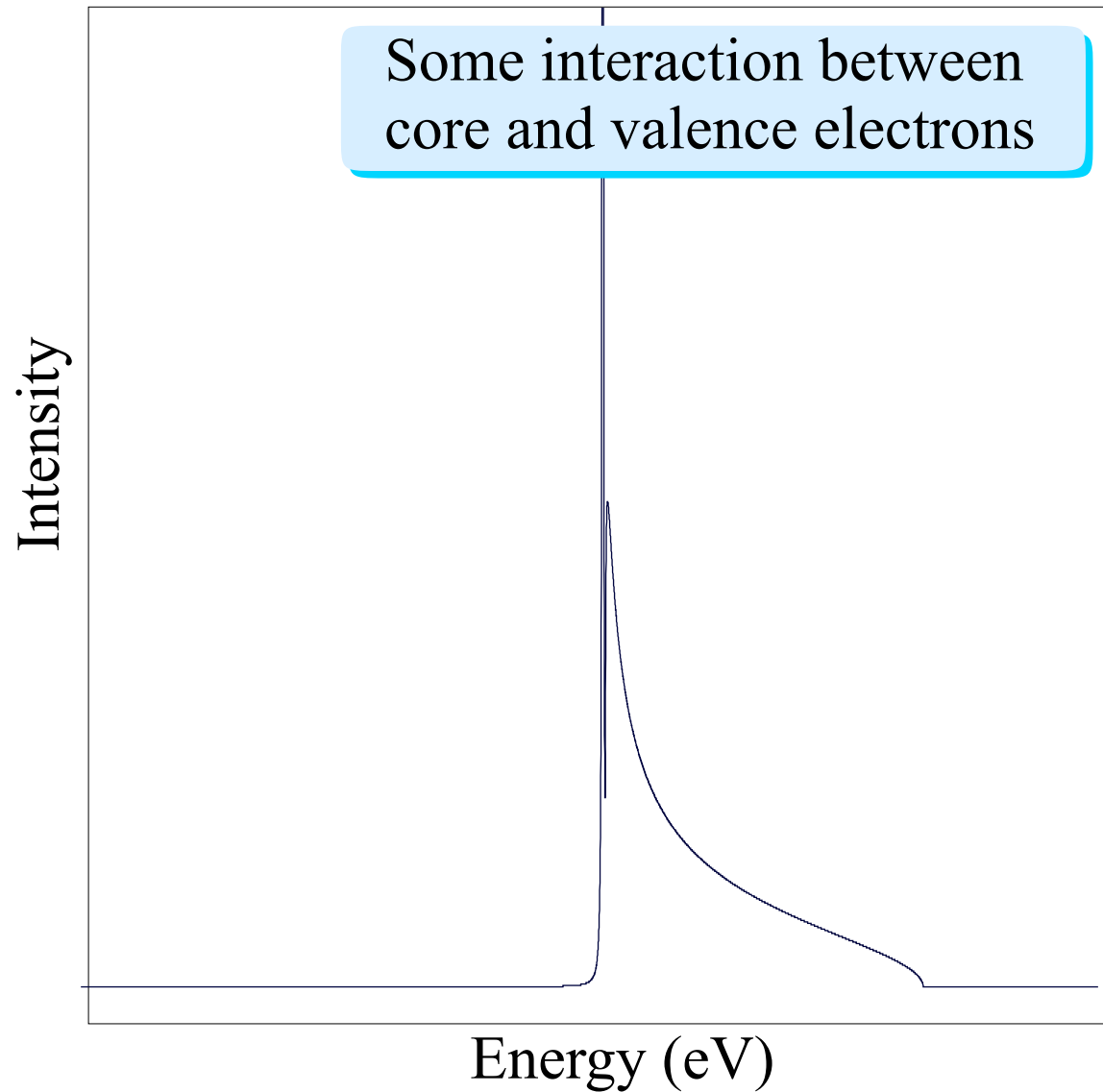
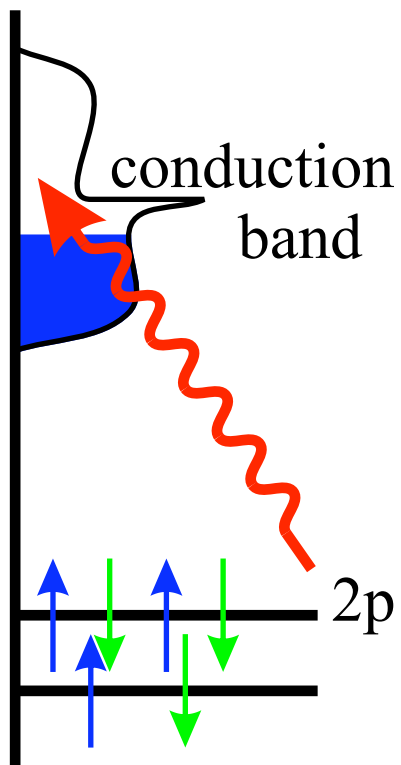
Some interaction between core and valence electrons

Intensity

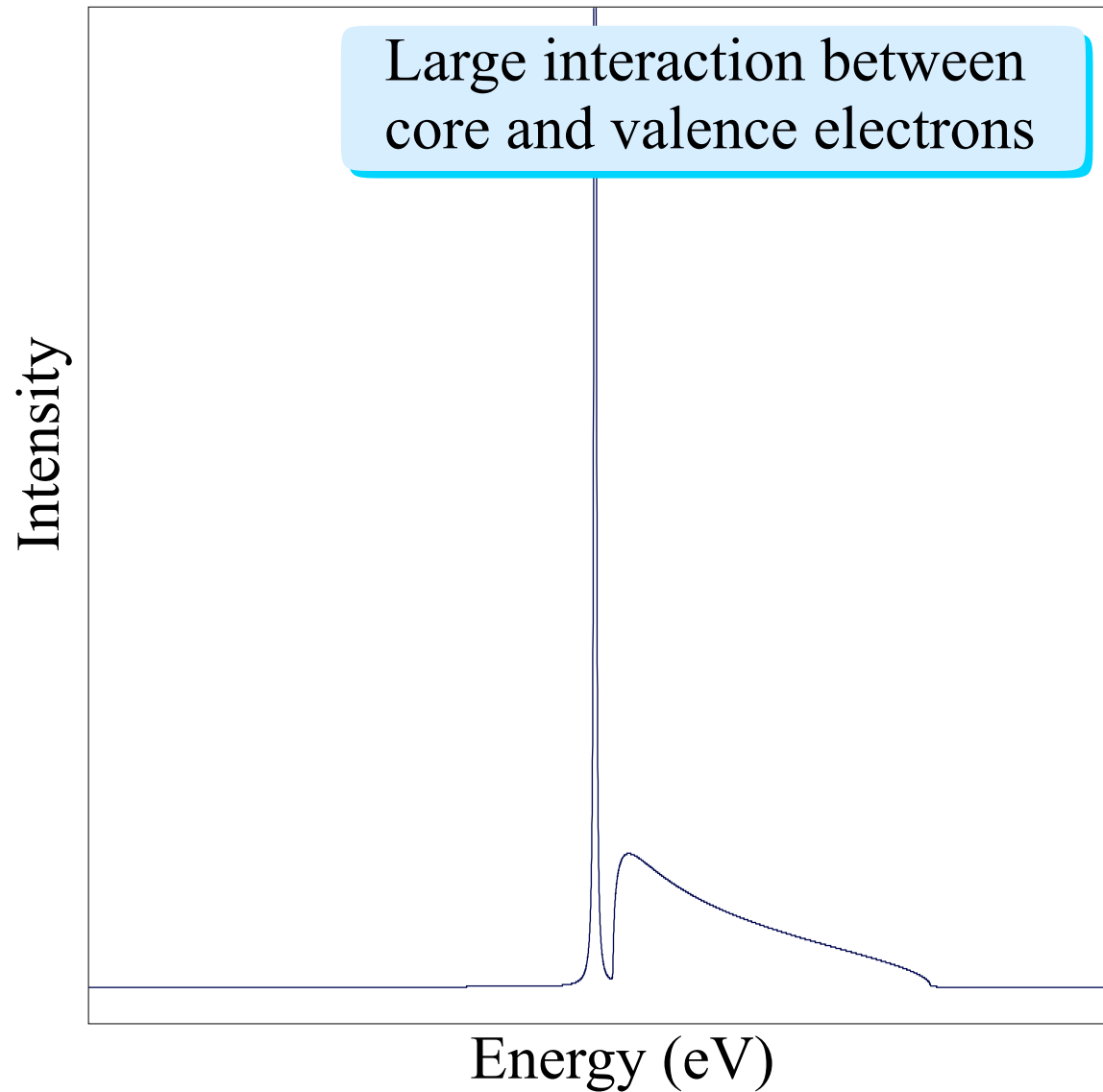
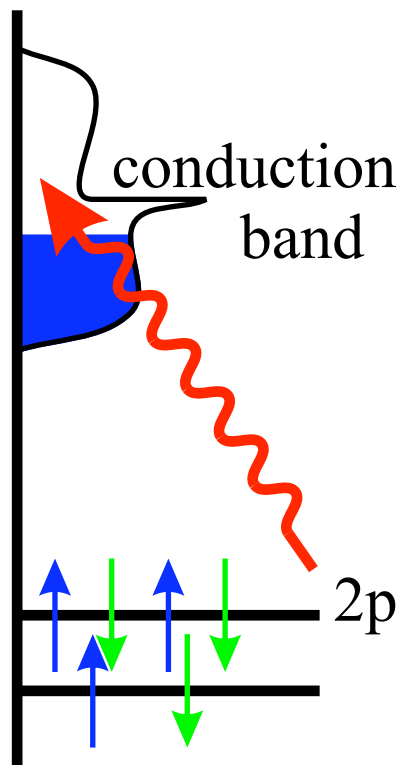
Energy (eV)



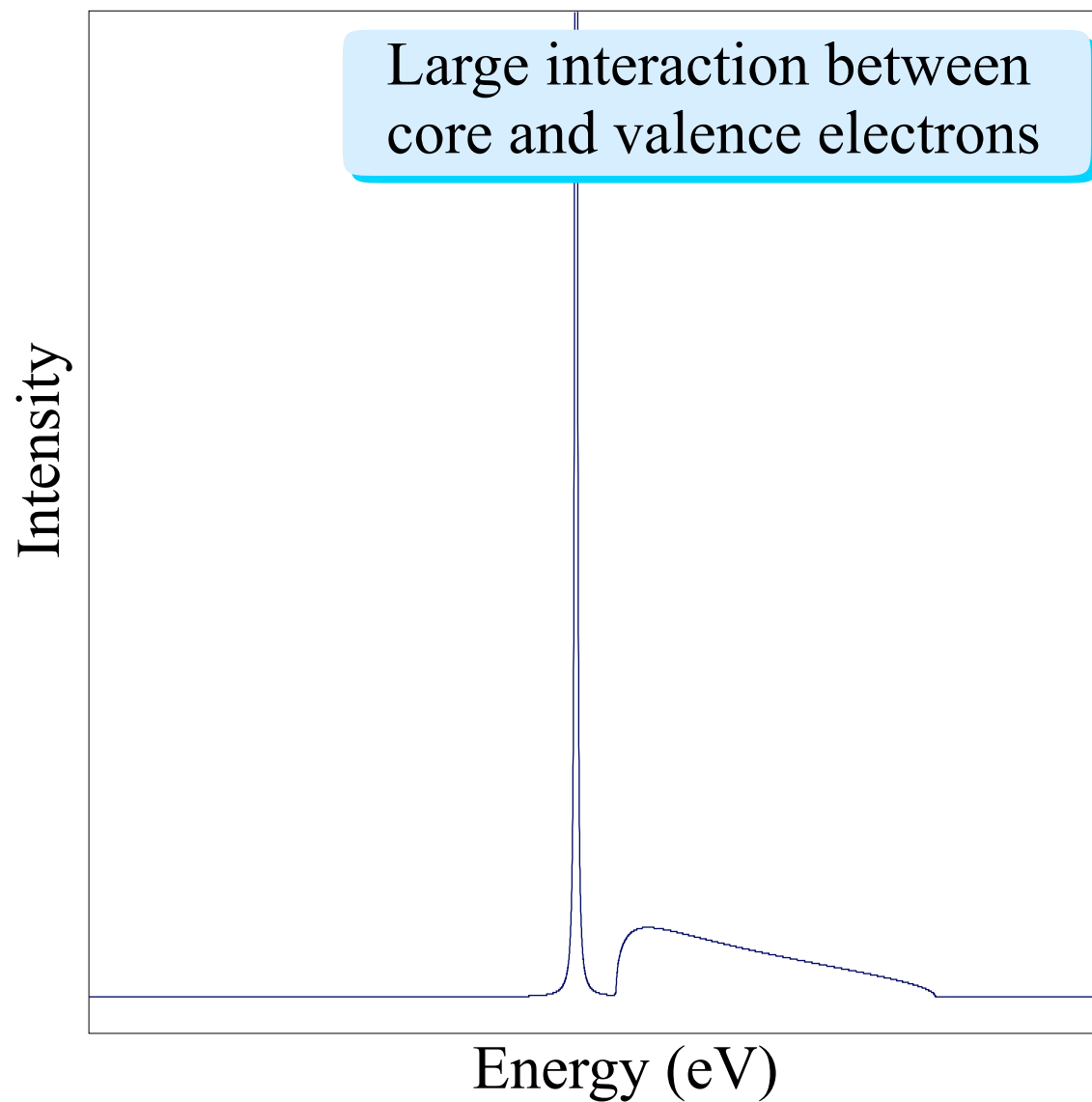
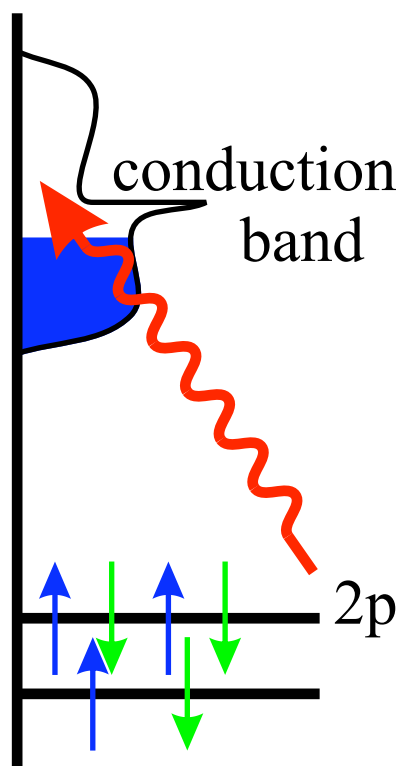
X-ray Absorption Spectroscopy: Excitons



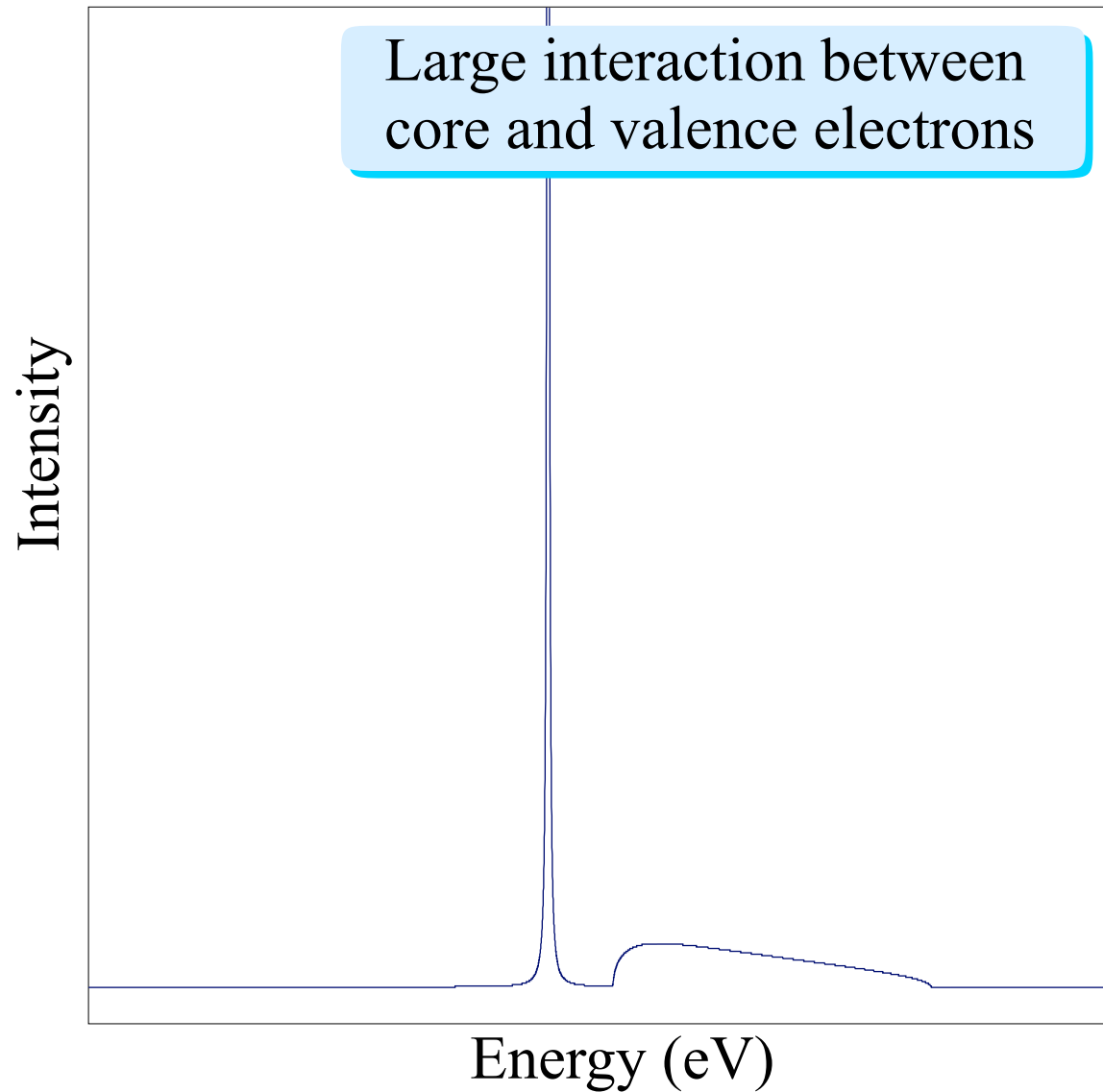
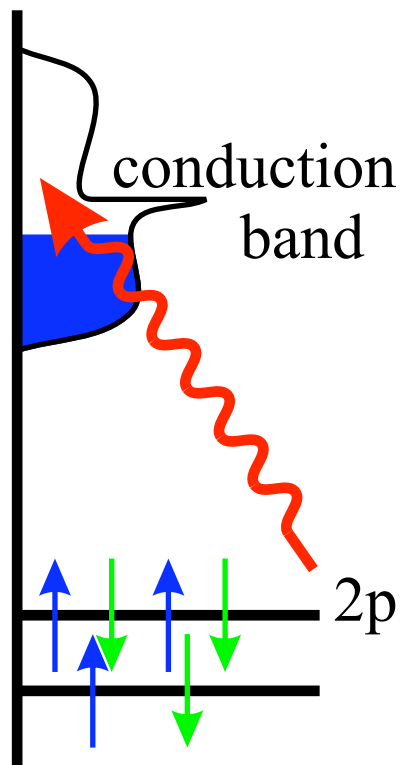
X-ray Absorption Spectroscopy: Excitons



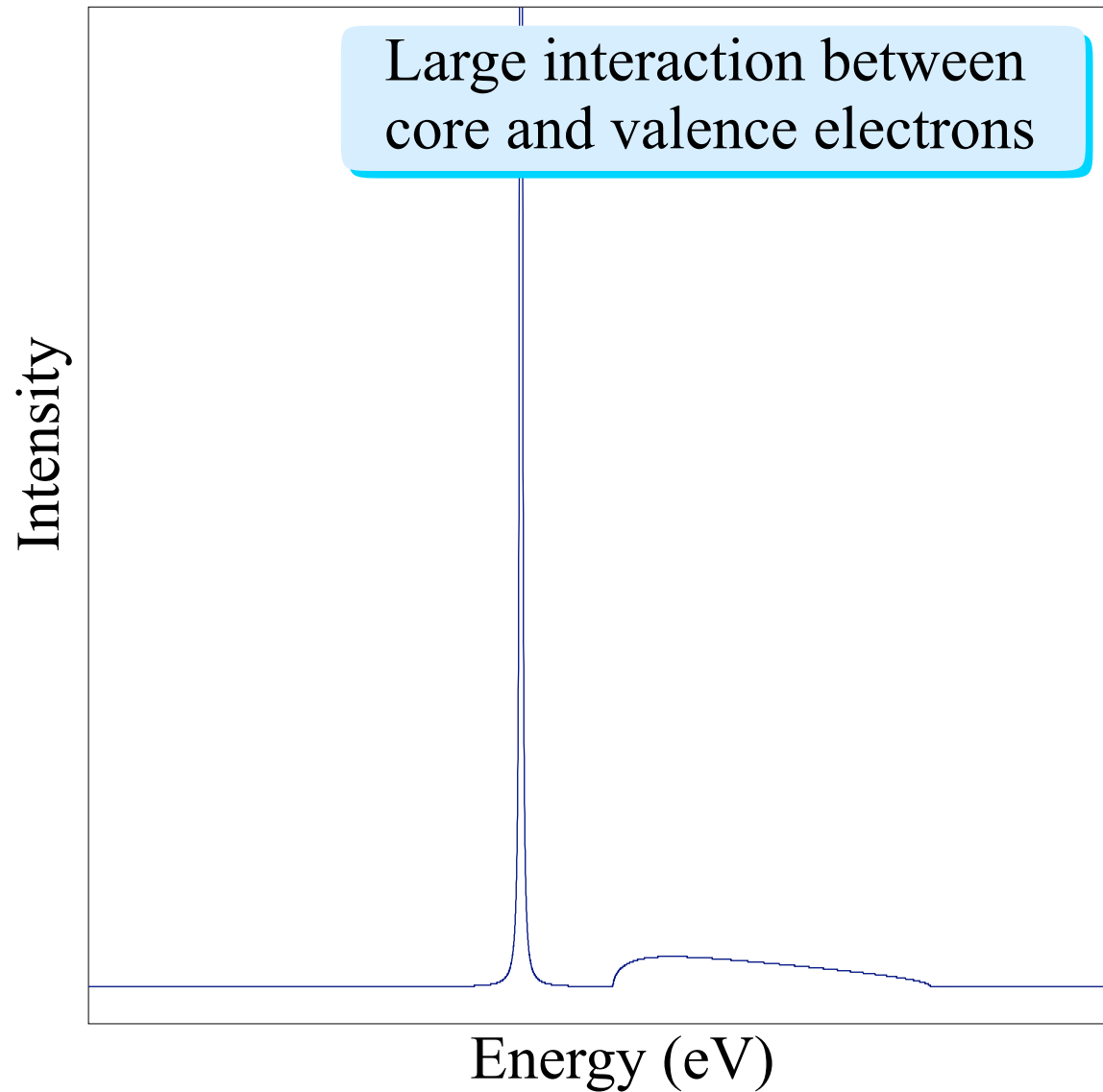
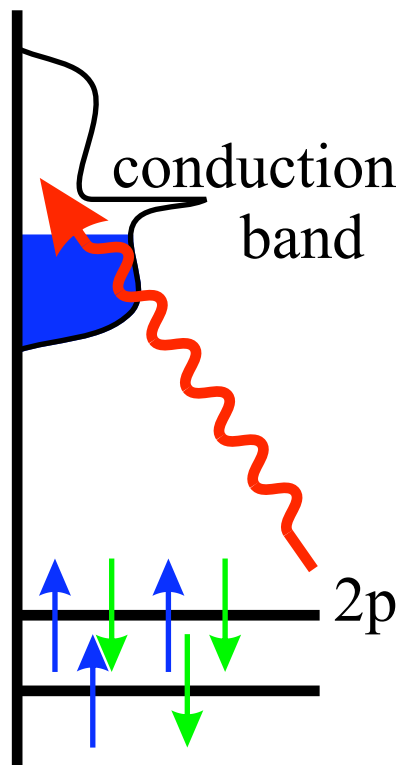
X-ray Absorption Spectroscopy: Excitons



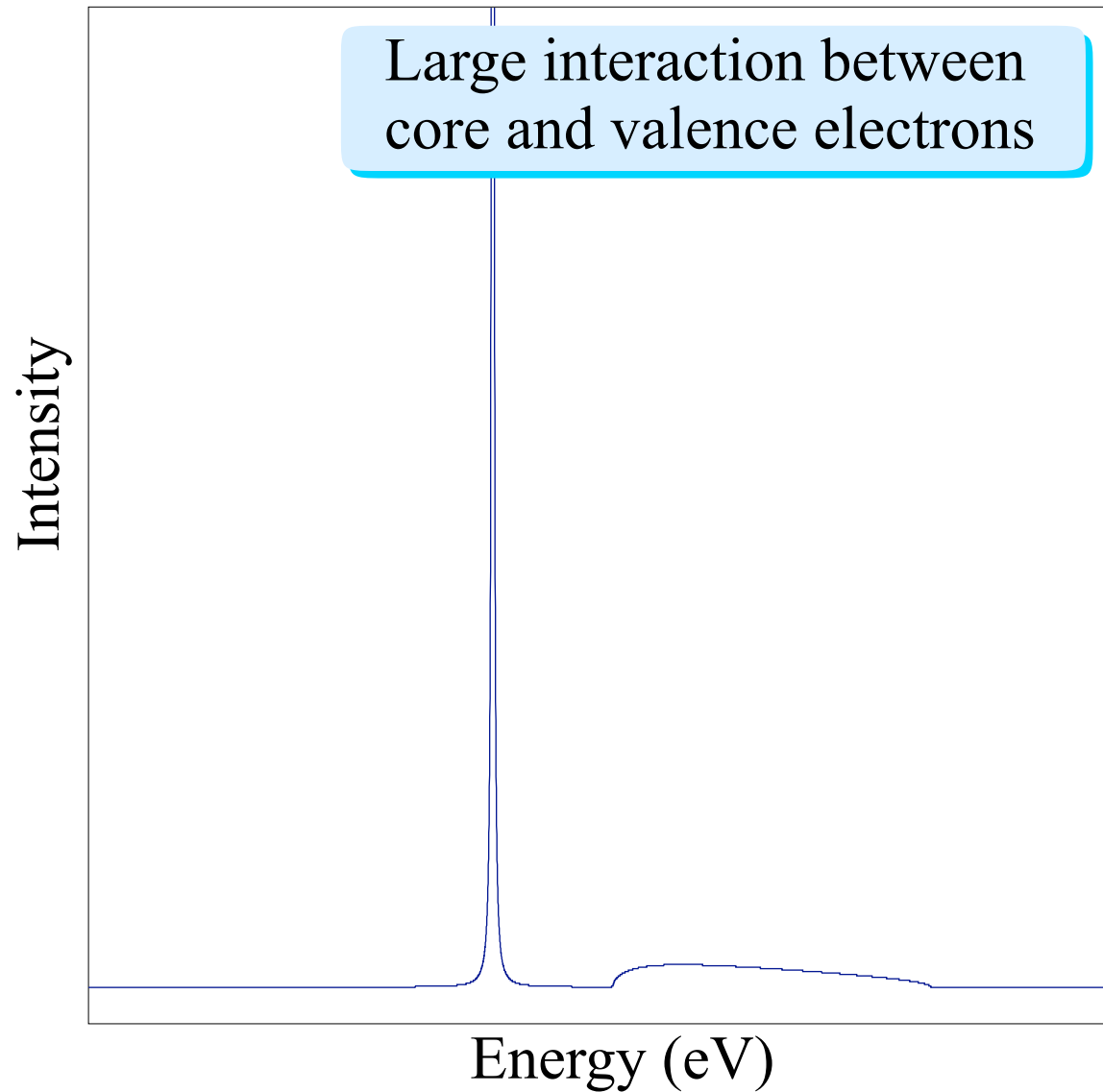
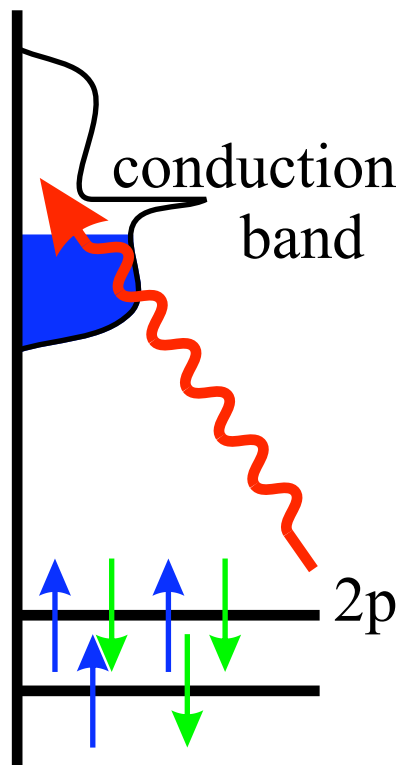
X-ray Absorption Spectroscopy: Excitons



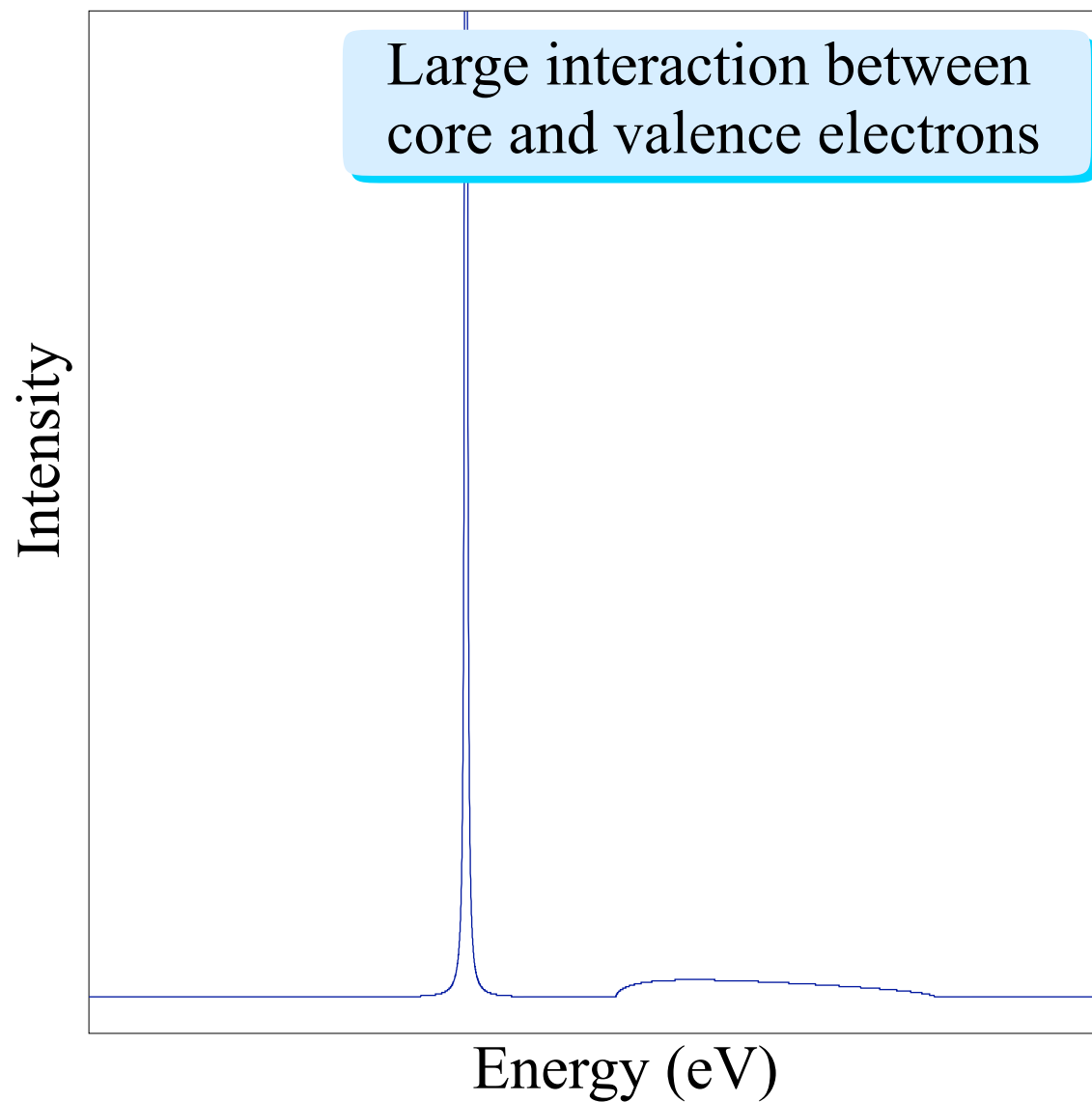
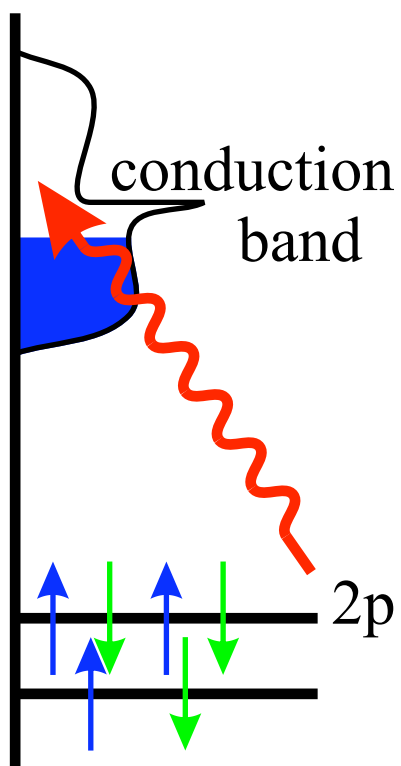
X-ray Absorption Spectroscopy: Excitons



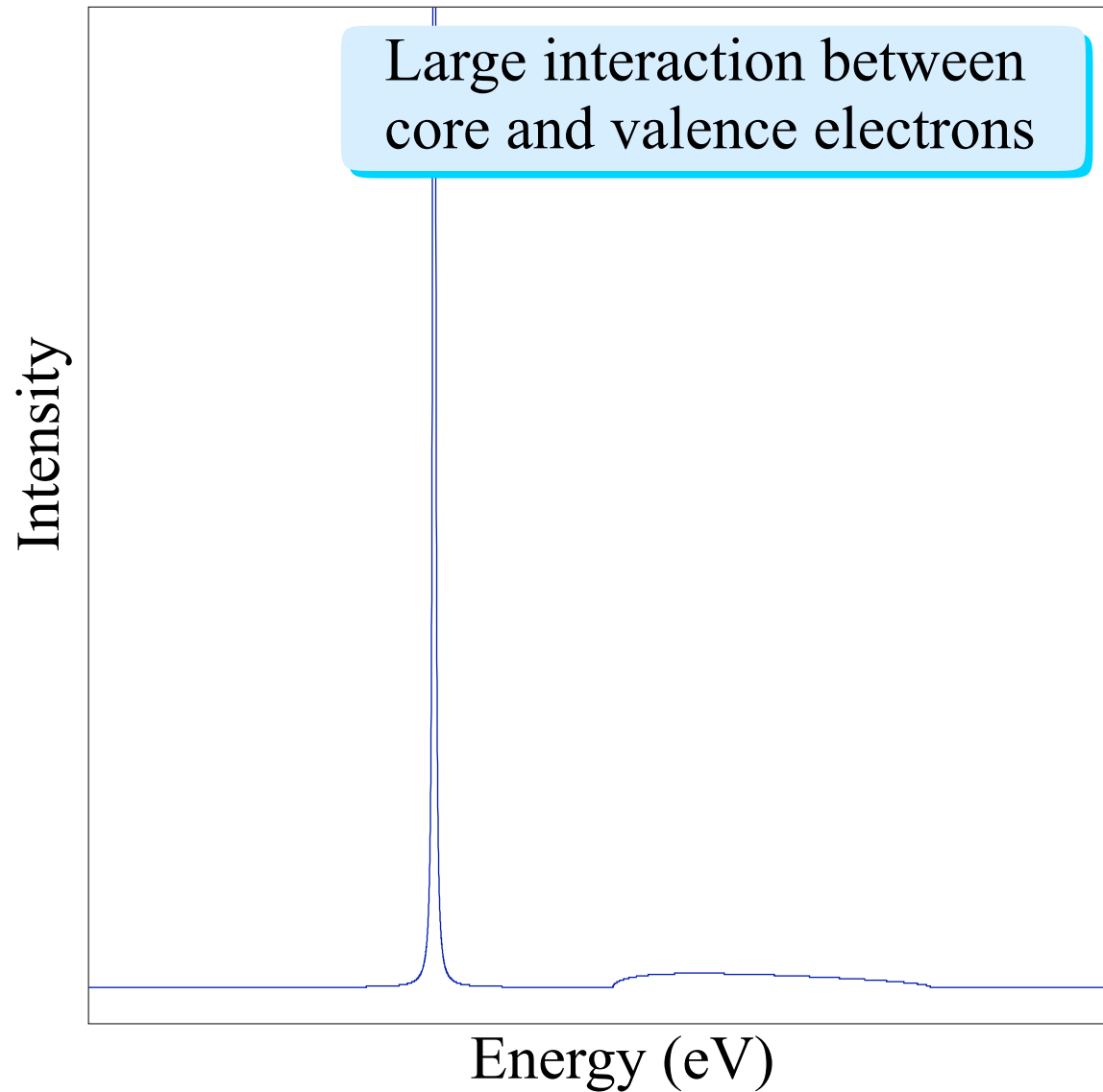
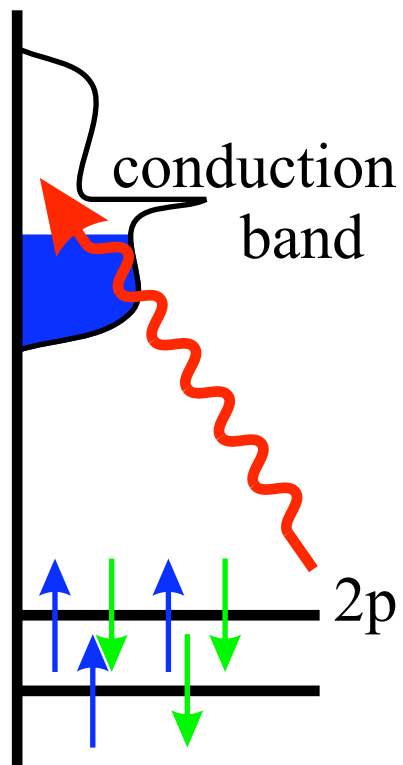
X-ray Absorption Spectroscopy: Excitons



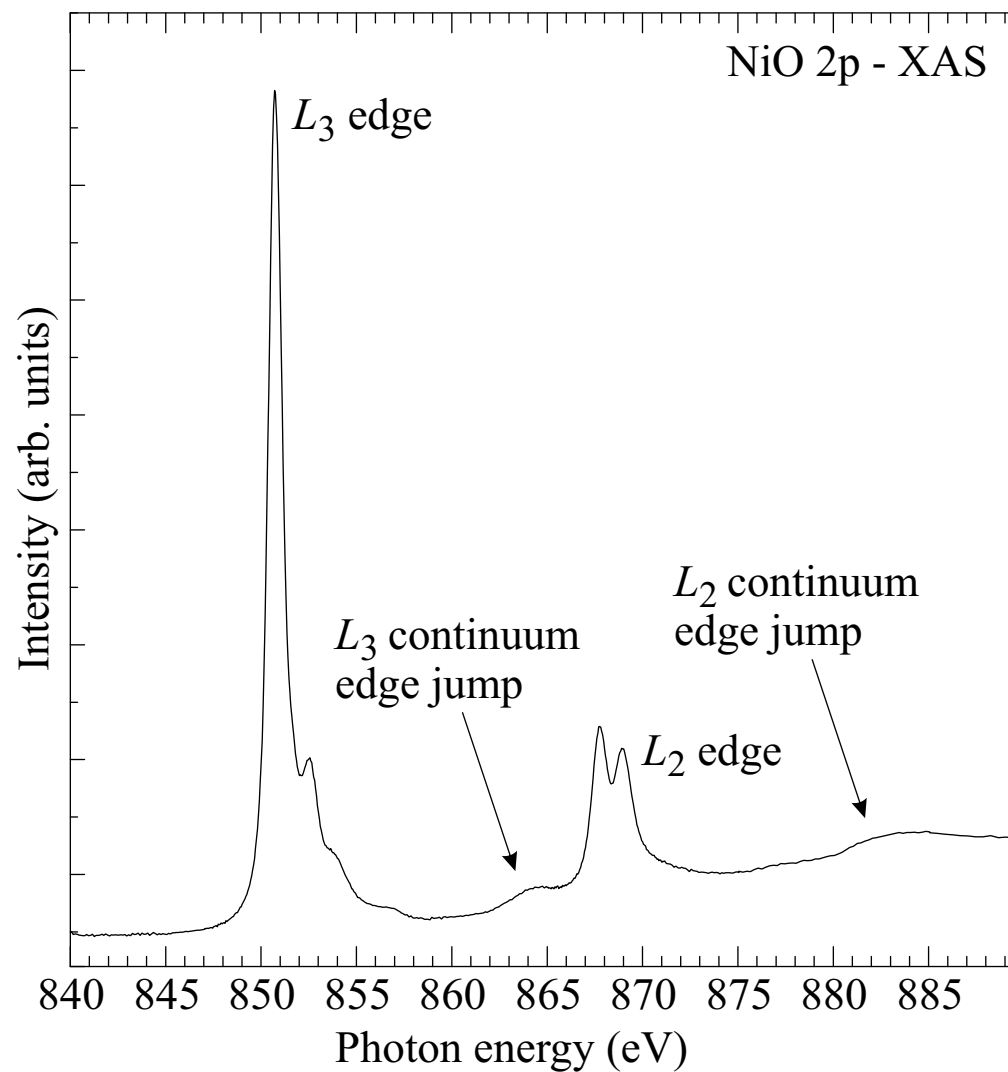
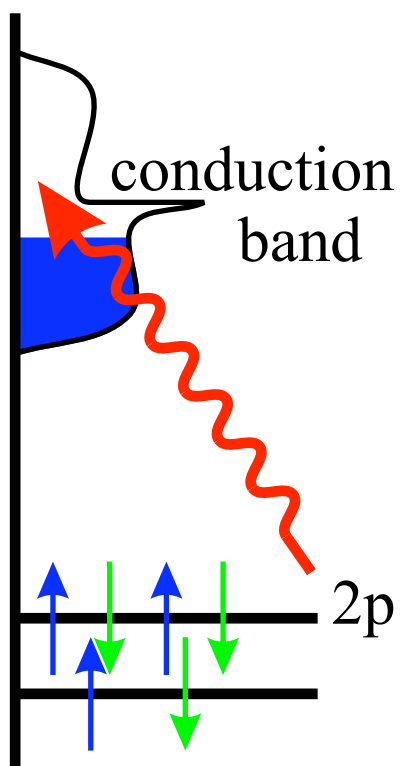
X-ray Absorption Spectroscopy: Excitons



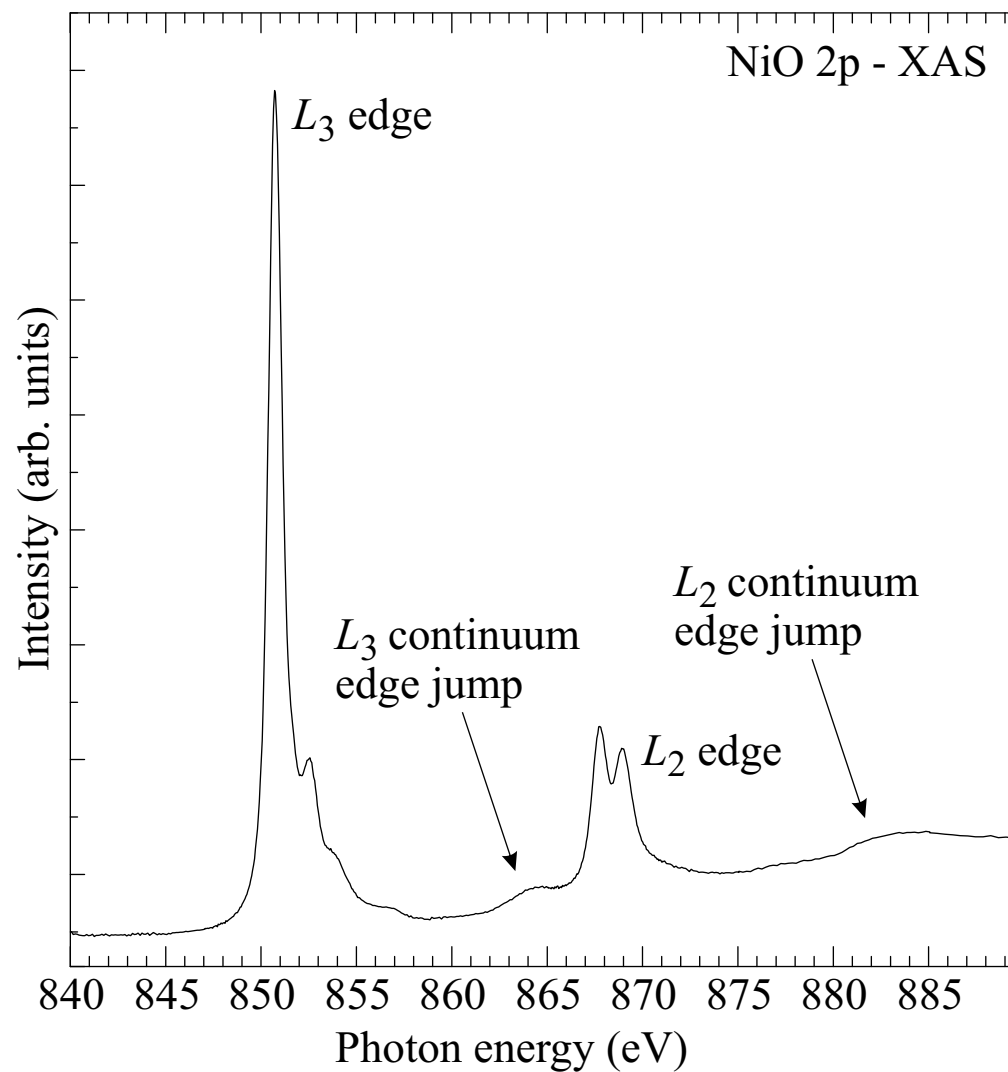
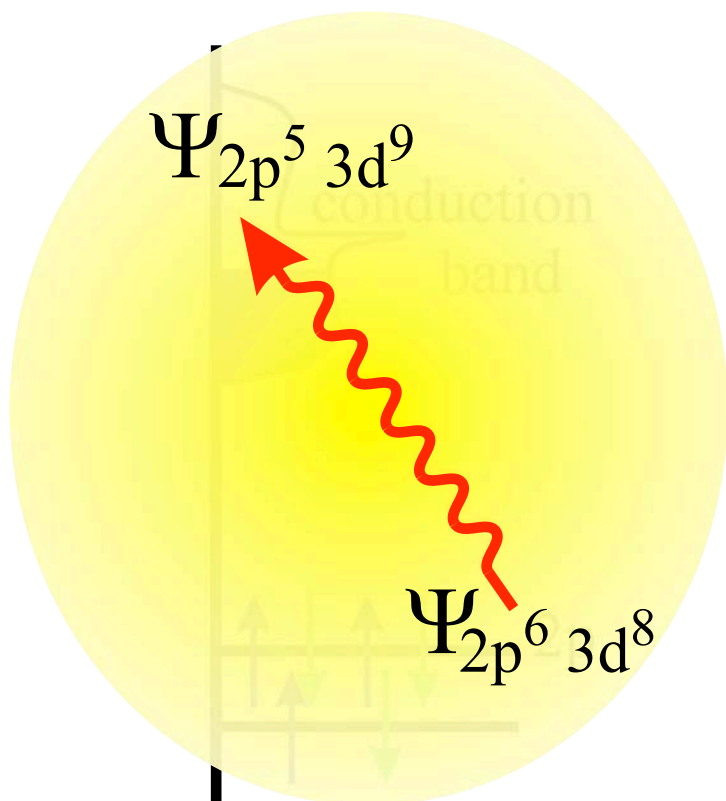
X-ray Absorption Spectroscopy: Excitons



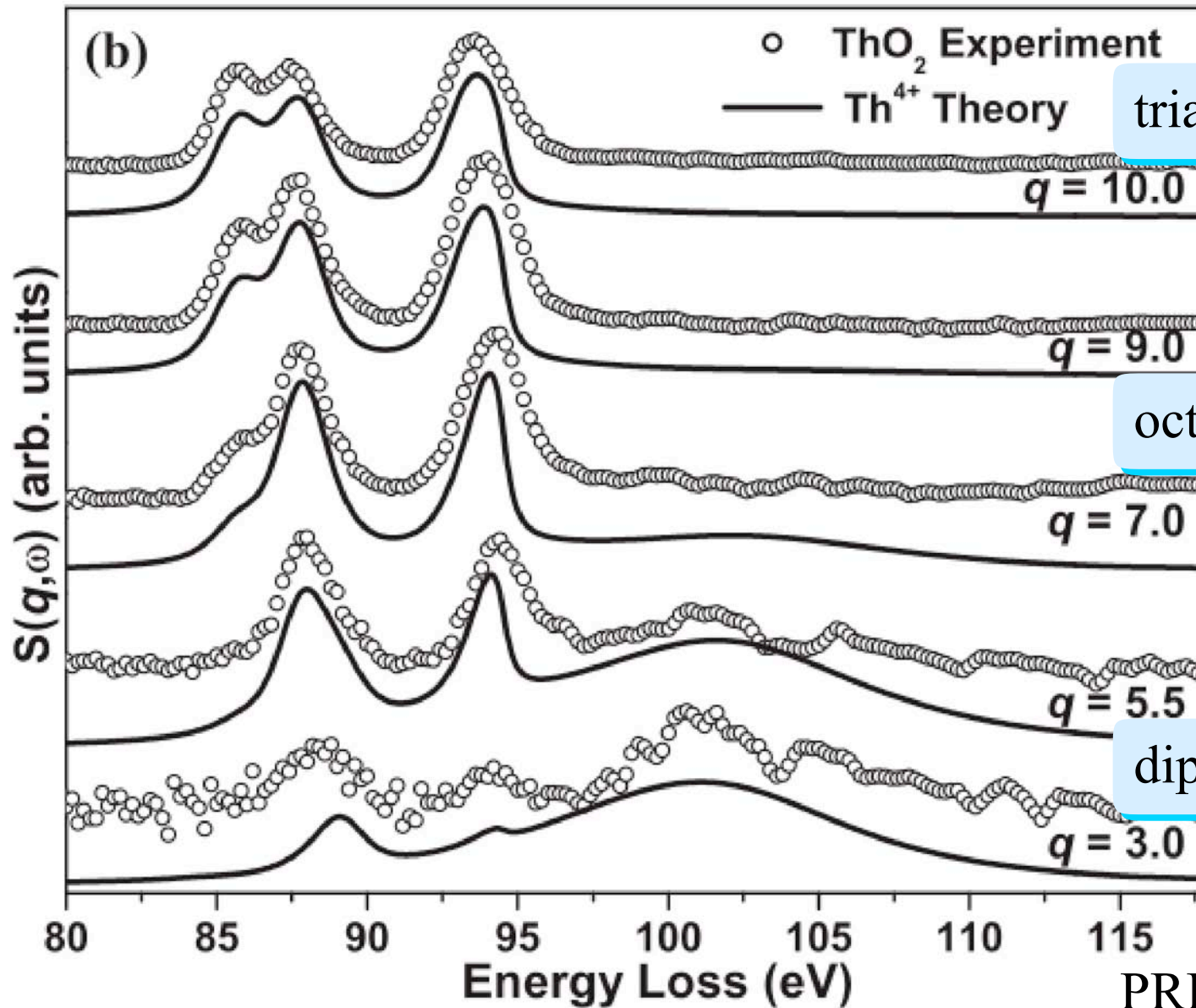
X-ray Absorption Spectroscopy: Excitons



X-ray Absorption Spectroscopy: Excitons



X-ray Absorption Spectroscopy: Excitons and continuum excitations

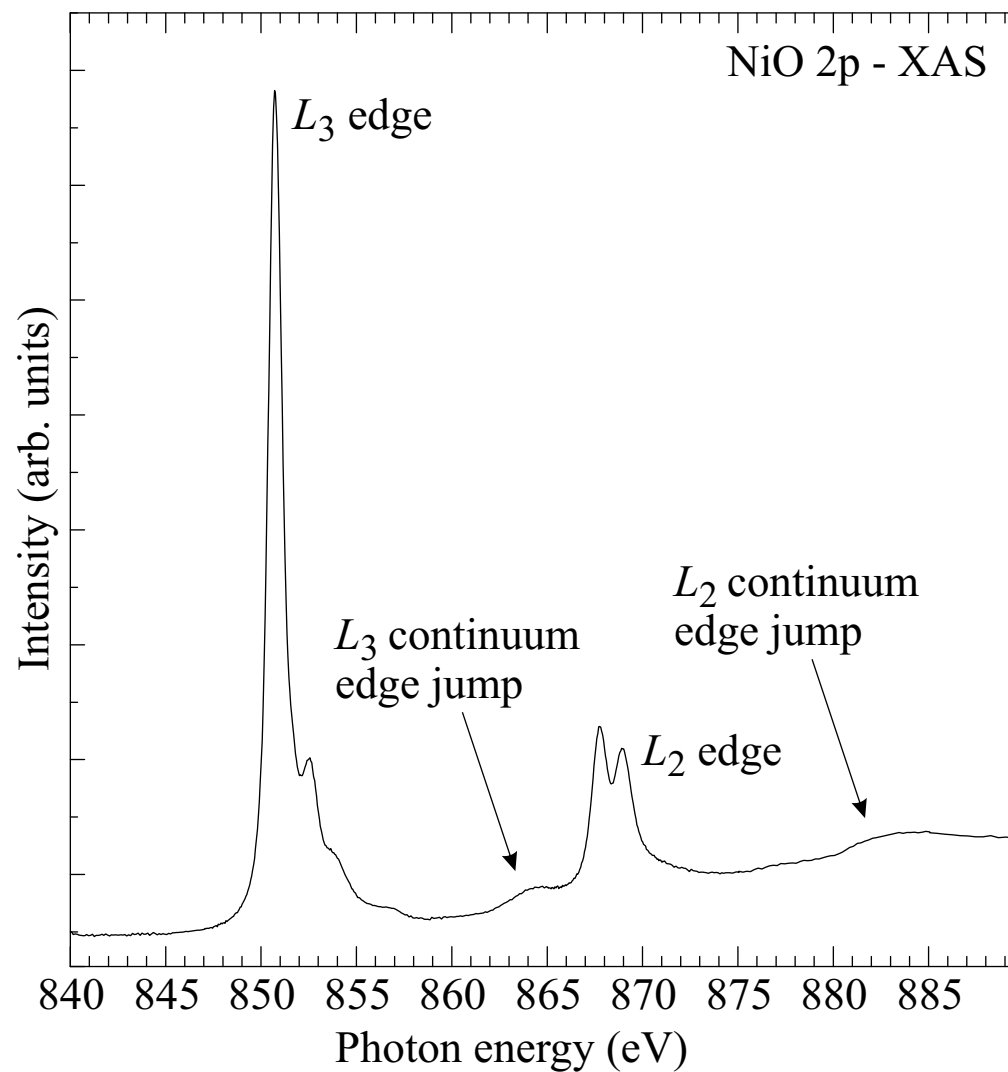
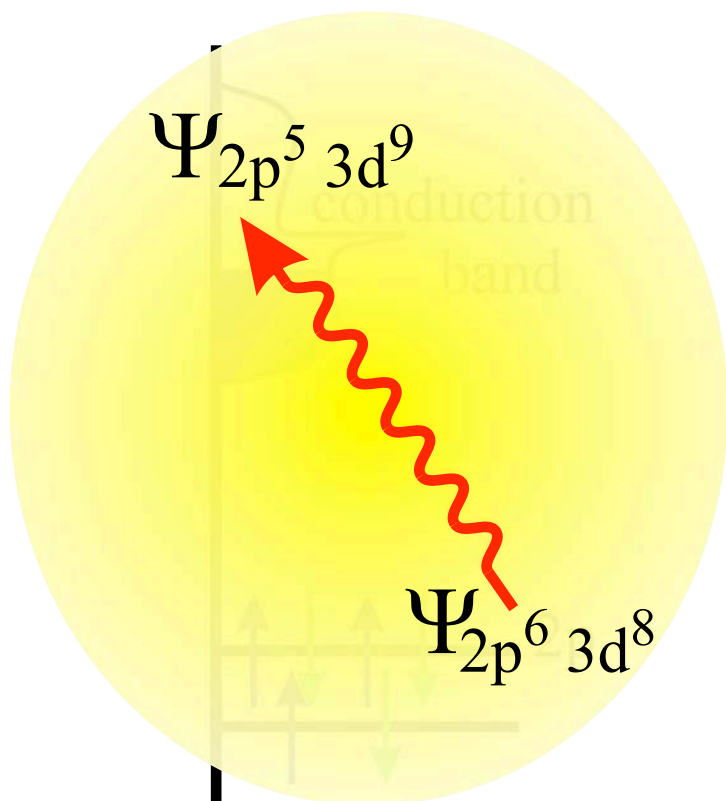


triacontadipole transitions

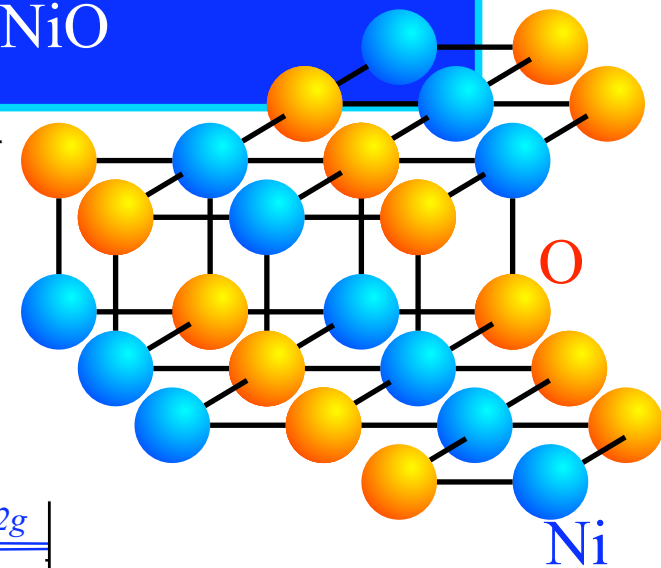
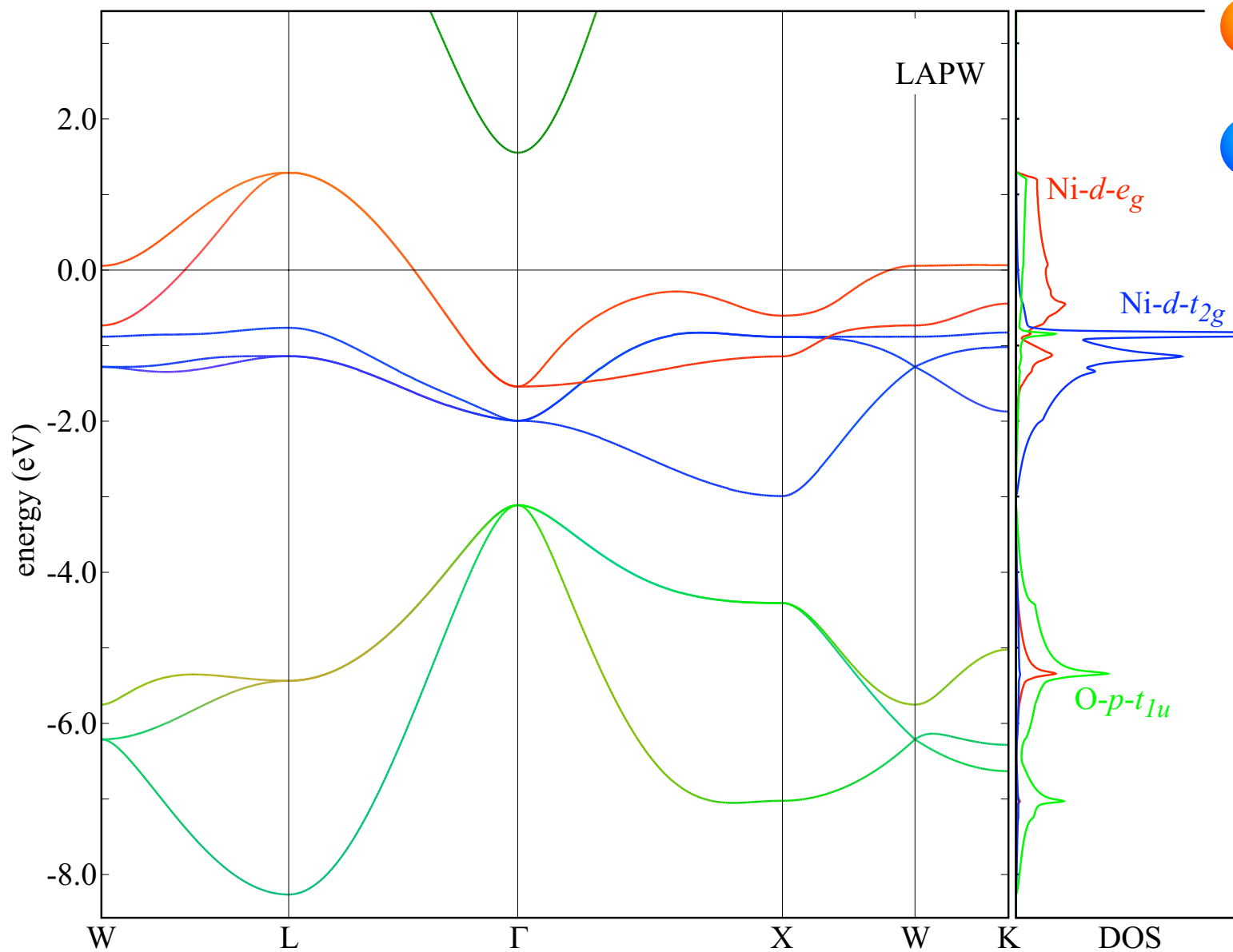
octupole transitions

dipole transitions

X-ray Absorption Spectroscopy: Excitons



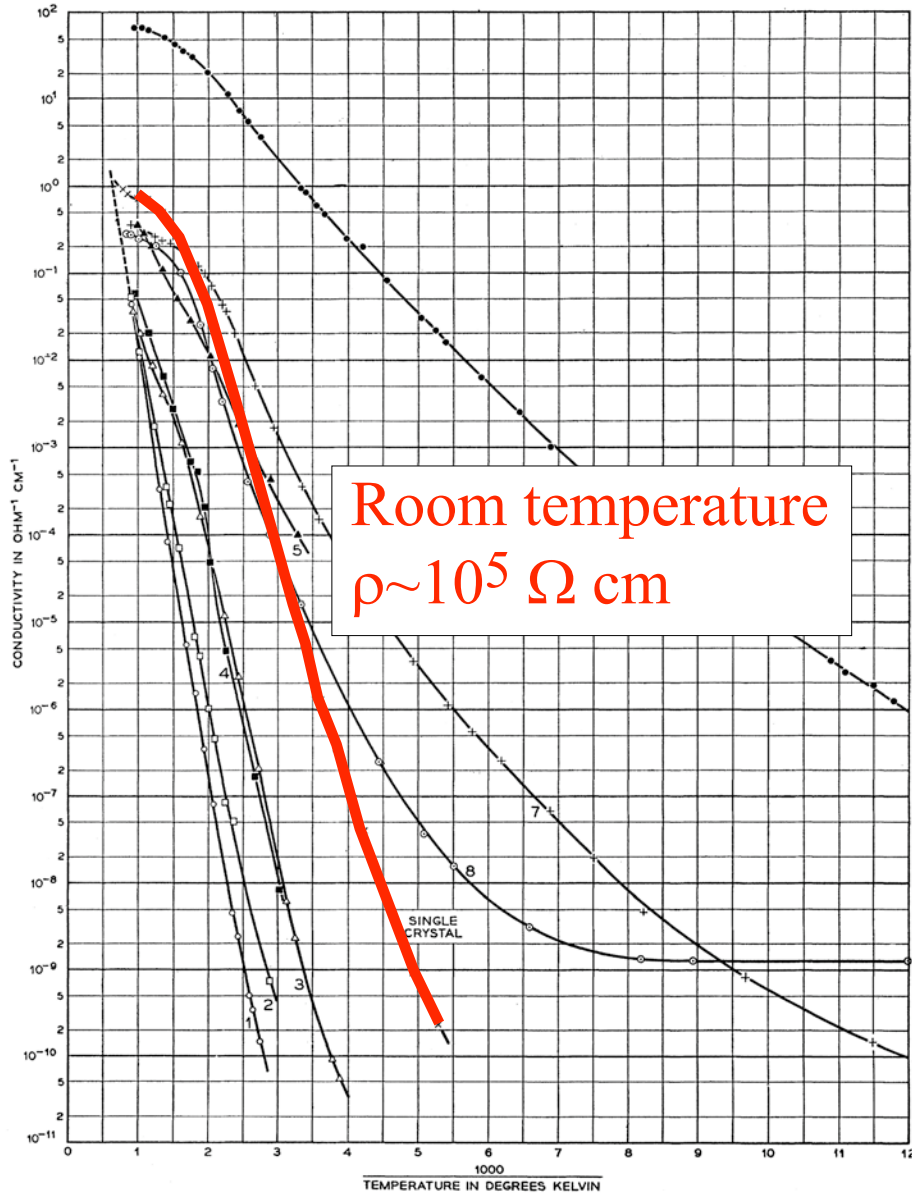
How to calculate XAS: an example on NiO



Within LDA NiO is a metal; experimentally it is a good insulator

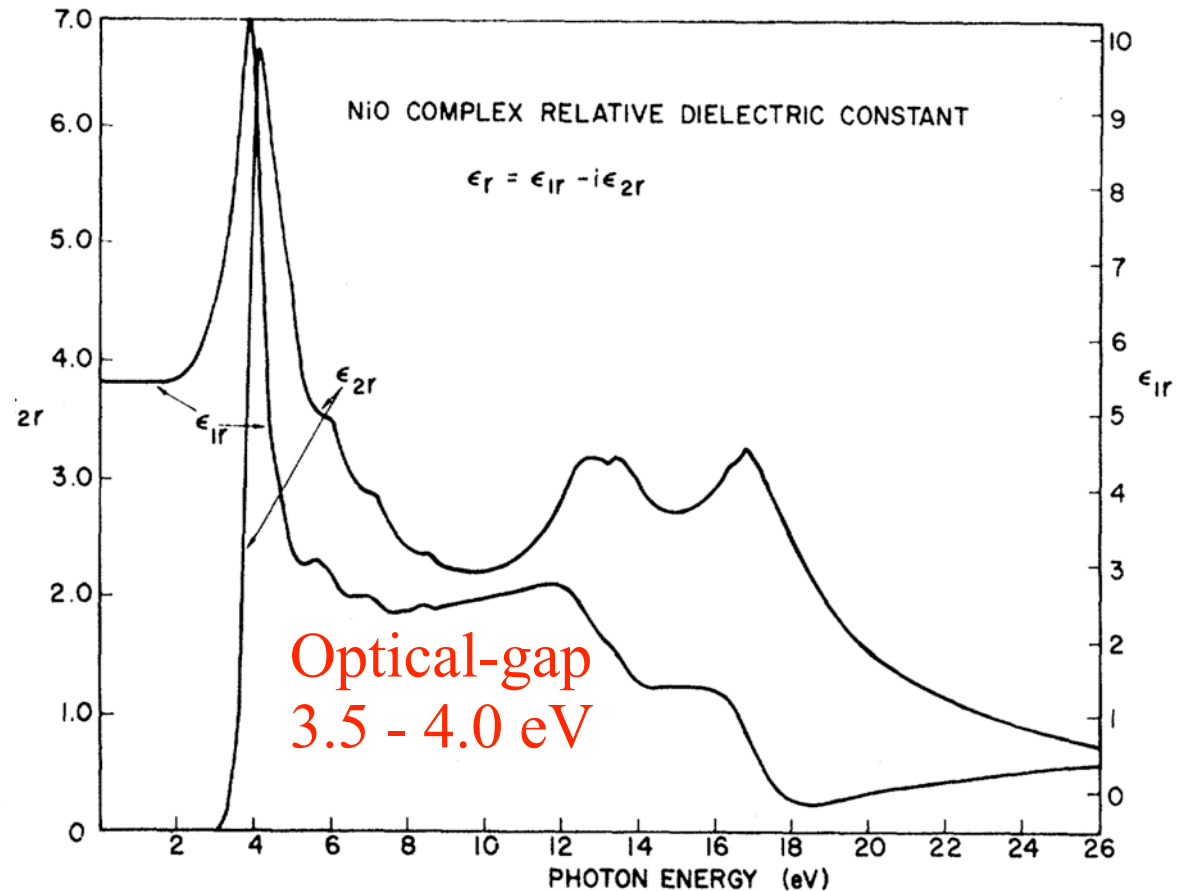
ELECTRICAL PROPERTIES OF NiO

1201

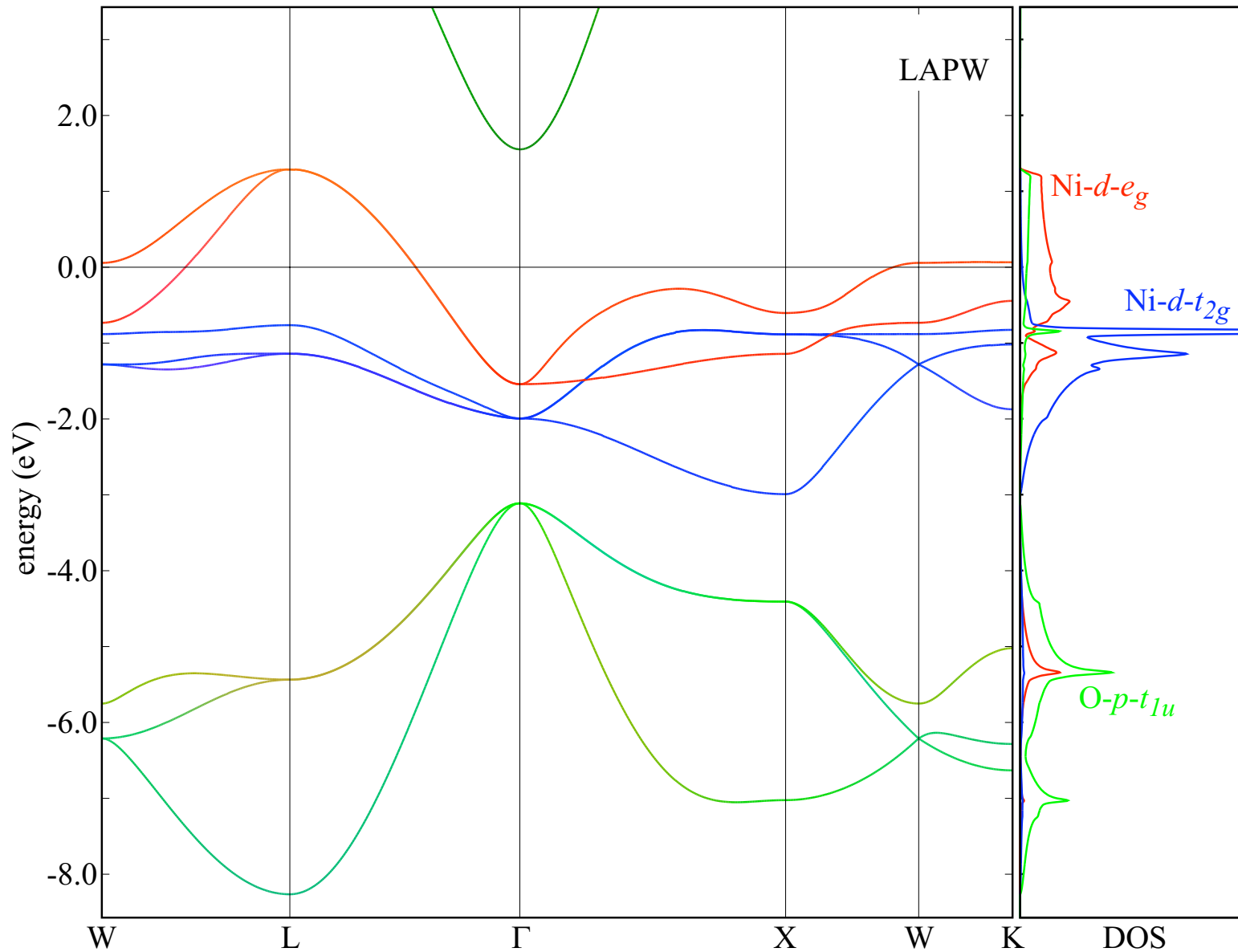


F. J. Morin Phys. Rev. **93**, 1199 (1954)

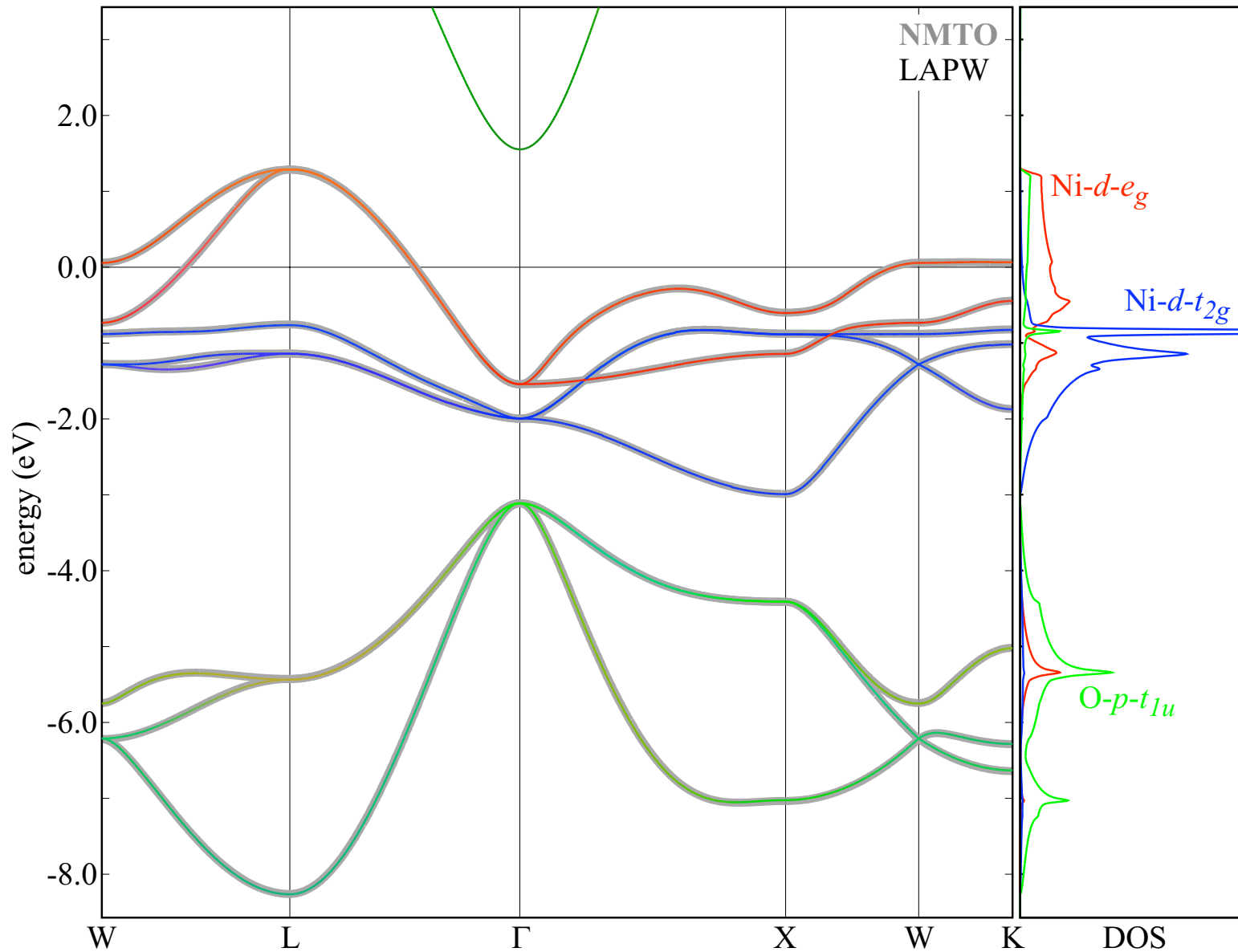
R. Newman and R. M. Chrenko
Phys. Rev. **114** 1507 (1959)



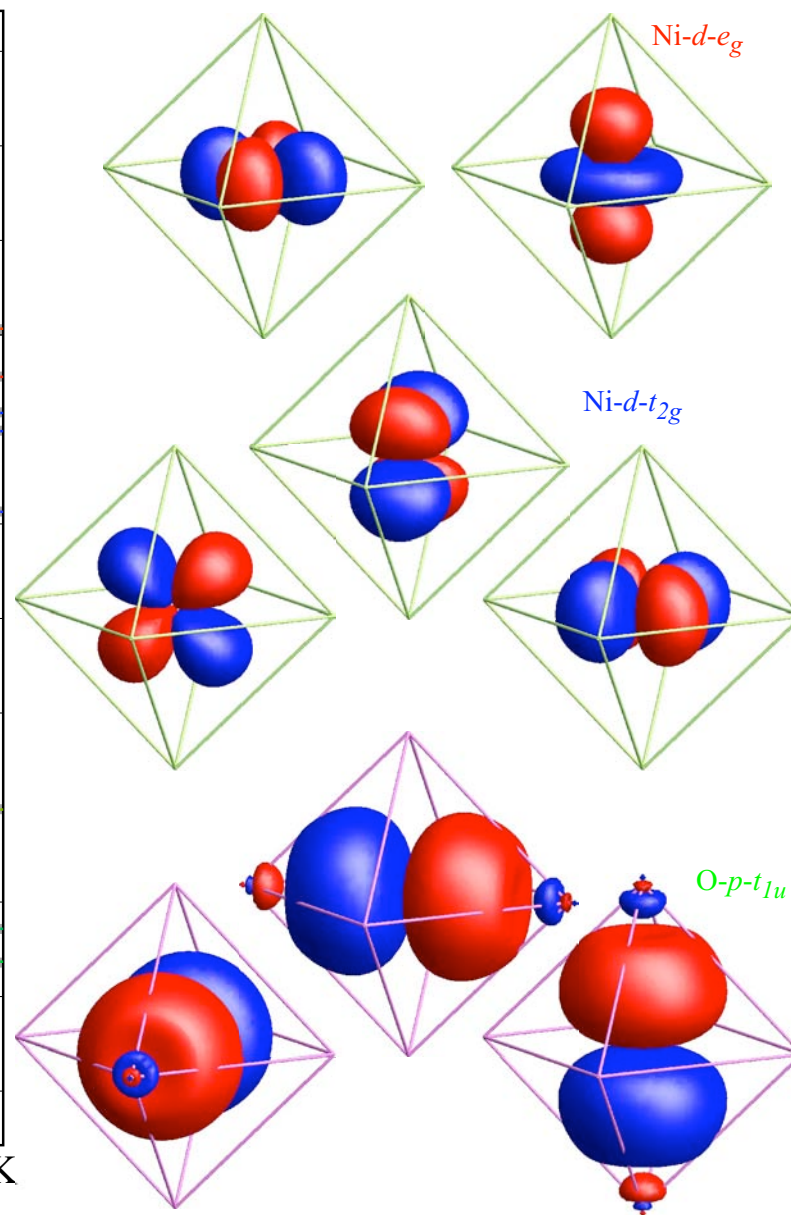
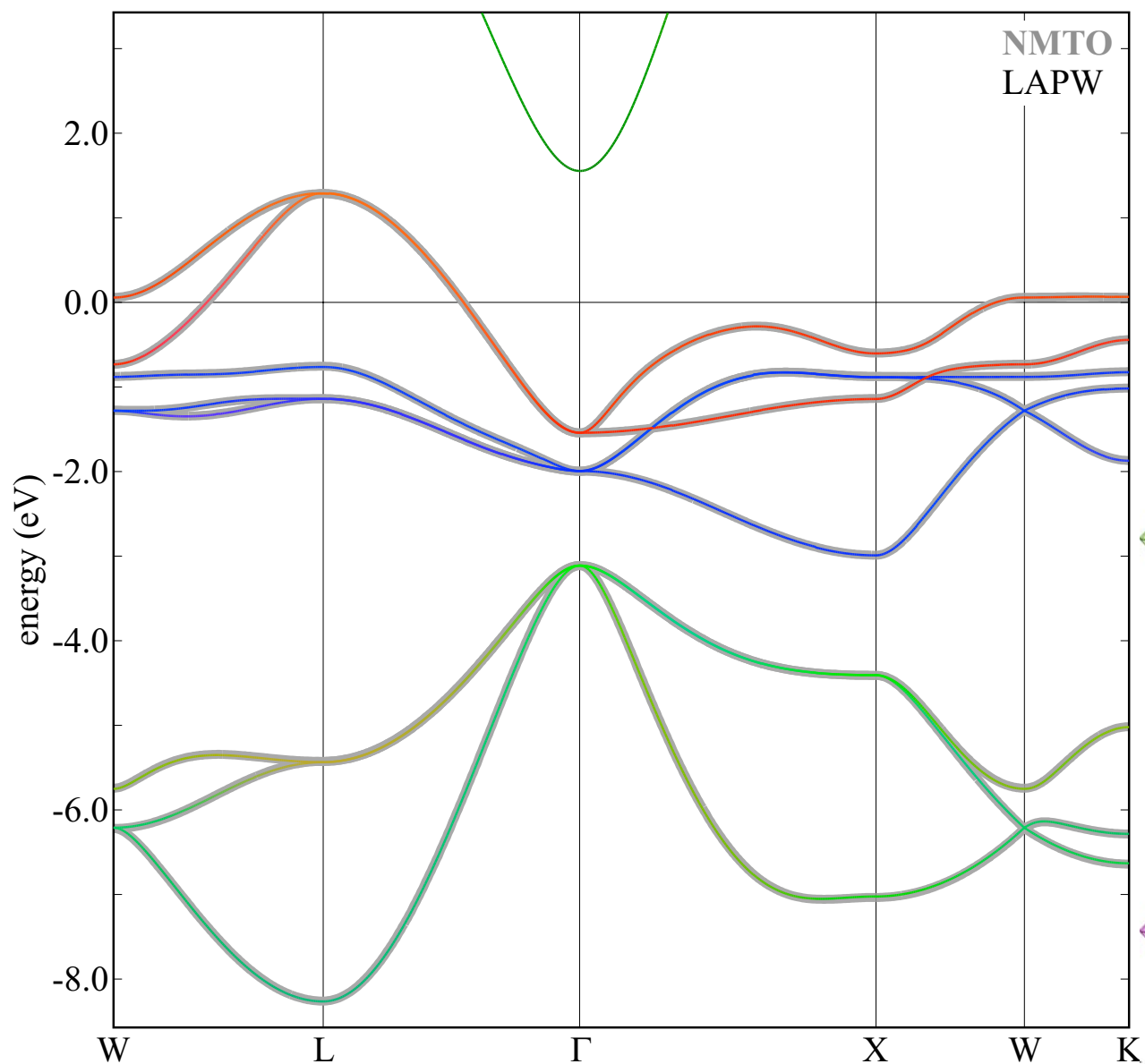
Although NiO is metal in LDA there are many things correct



Use the LAPW potential for an NMTO calculation

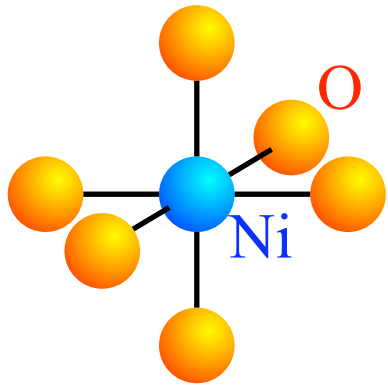
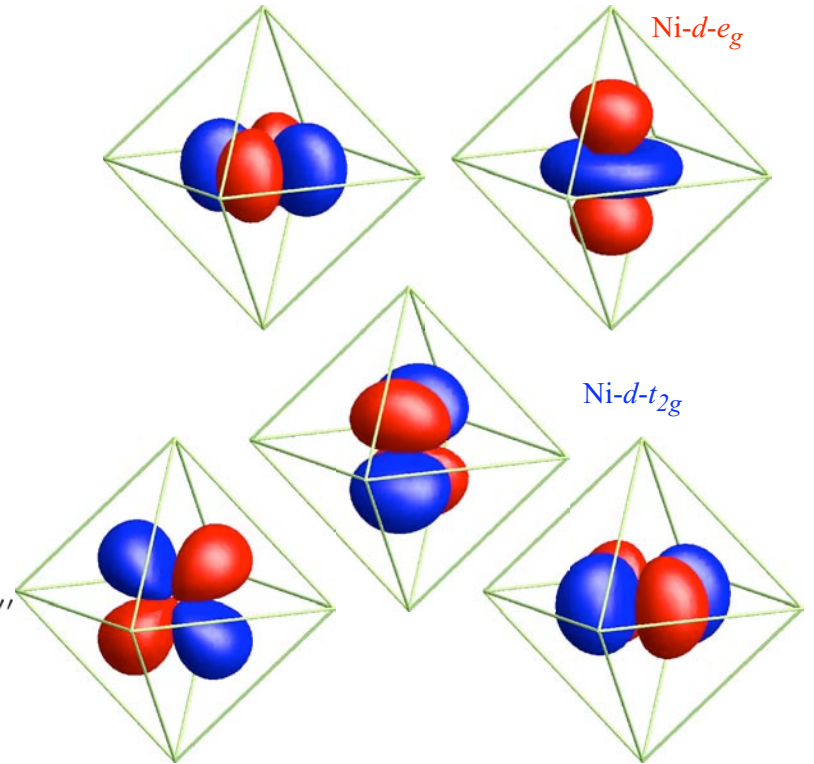


Obtain Wannier orbitals from the Muffin Tin Orbitals



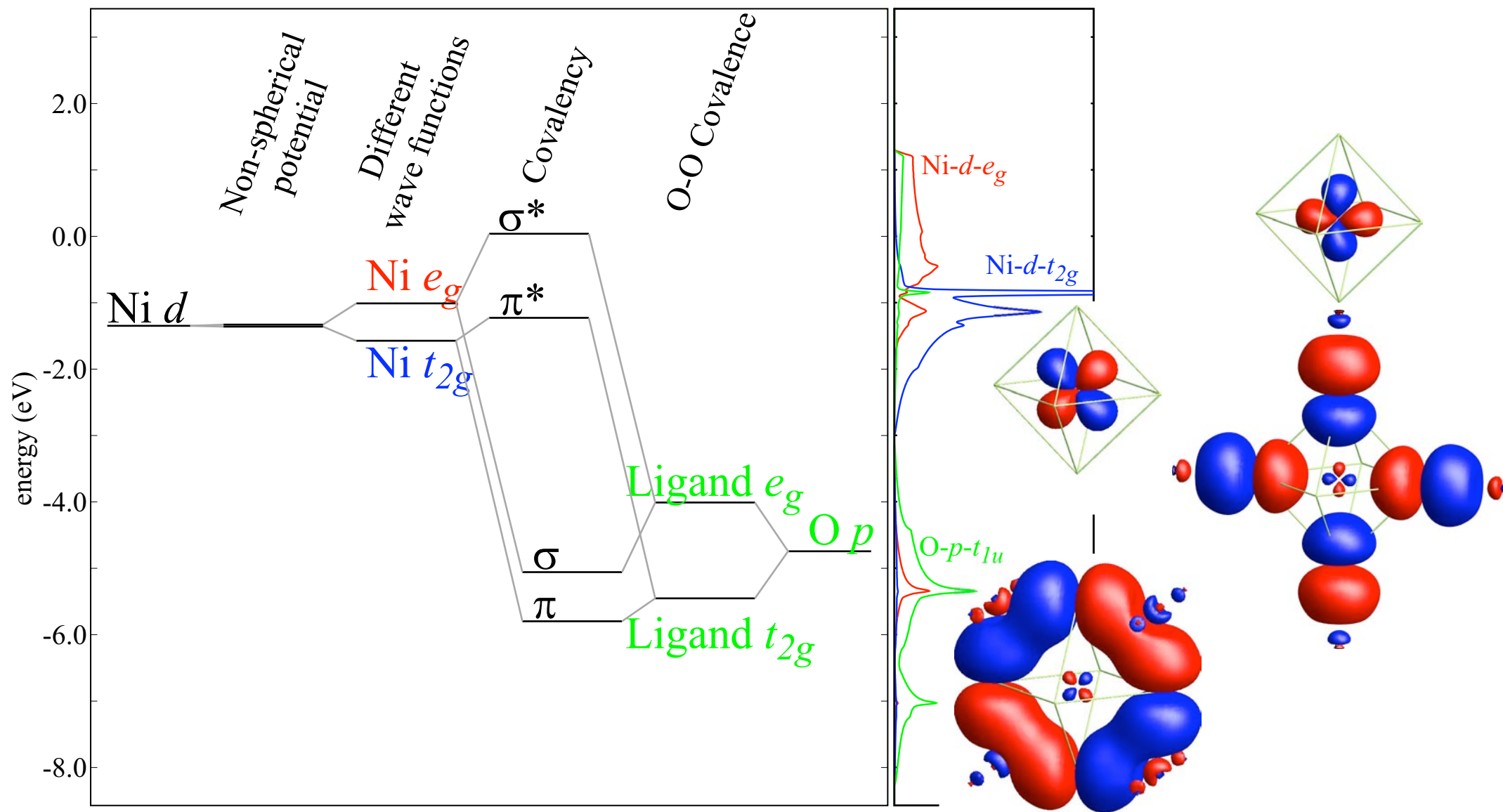
Full multiplet Ligand Field model based on LDA

$$\begin{aligned}
 H_i = & \sum_{\sigma, \tau, \tau'} \varepsilon_{\tau, \tau'} d_{\sigma, \tau}^\dagger d_{\sigma, \tau'} \\
 & + \sum_{\sigma, \tau, \tau'} t_{\tau, \tau'} (d_{\sigma, \tau}^\dagger L_{\sigma, \tau'} + L_{\sigma, \tau}^\dagger d_{\sigma, \tau'}) \\
 & + \sum_{\sigma, \tau, \tau'} t'_{\tau, \tau'} L_{\sigma, \tau}^\dagger L_{\sigma, \tau'} \\
 & + \sum_{\sigma, \tau} H_{EX} \sigma d_{\sigma, \tau}^\dagger d_{\sigma, \tau} \\
 & + \sum_{\sigma, \sigma', \tau, \tau'} \zeta_{\sigma, \sigma', \tau, \tau'} d_{\sigma, \tau}^\dagger d_{\sigma', \tau'} \\
 & + \sum_{\sigma, \sigma'', \sigma''', \tau, \tau', \tau'', \tau'''} U_{\sigma, \sigma'', \sigma''', \tau, \tau', \tau'', \tau'''} d_{\sigma, \tau}^\dagger d_{\sigma', \tau'}^\dagger d_{\sigma'', \tau''} d_{\sigma''', \tau'''}
 \end{aligned}$$



$$U_{\sigma, \sigma'', \sigma''', \tau, \tau', \tau'', \tau'''} = \frac{1}{2} \delta_{\sigma, \sigma''} \delta_{\sigma', \sigma'''} \int \int f_{\tau}^*(r_1) f_{\tau'}^*(r_2) \frac{e^2}{|r_1 - r_2|} f_{\tau''}(r_1) f_{\tau'''}(r_2) \delta r_1 \delta r_2$$

One particle energy level diagram obtained from LDA

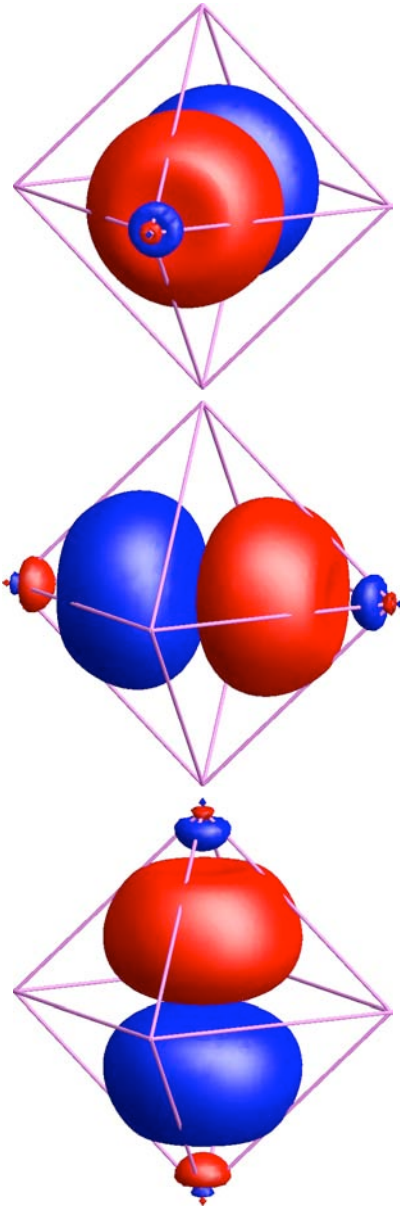


Two particle interactions on a level beyond LDA

O 2p orbitals

Orbitals are close to completely filled:
Chance that two holes meet on the same orbital is small.

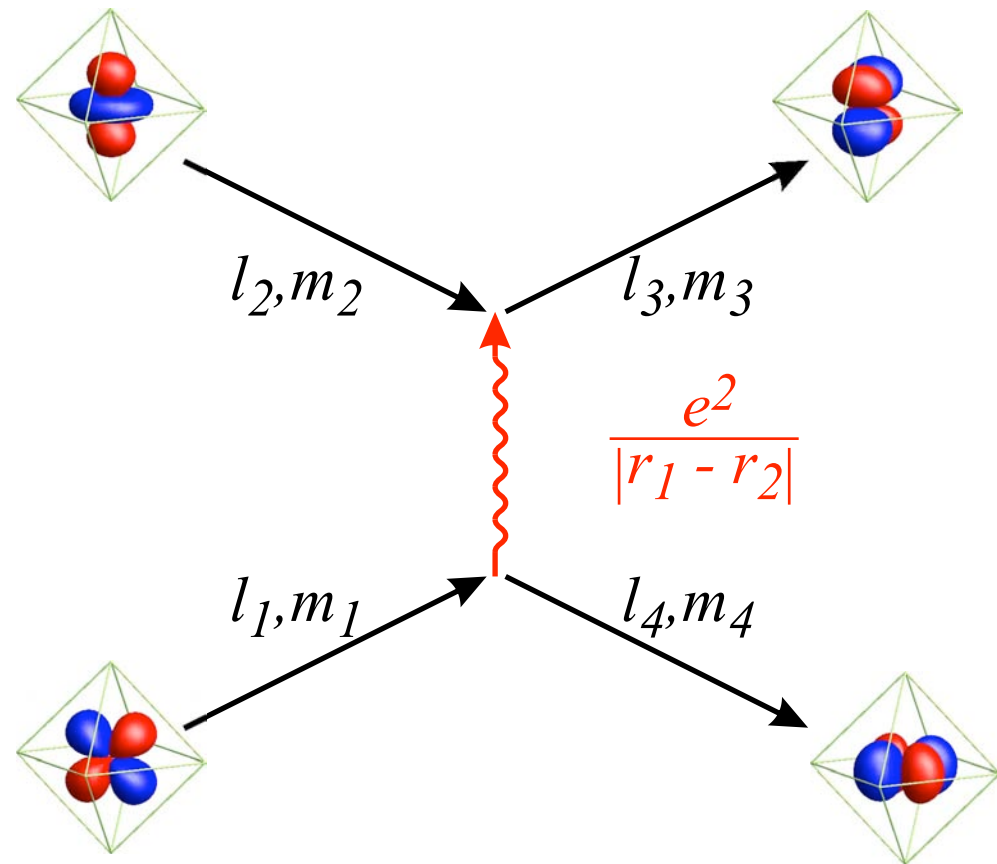
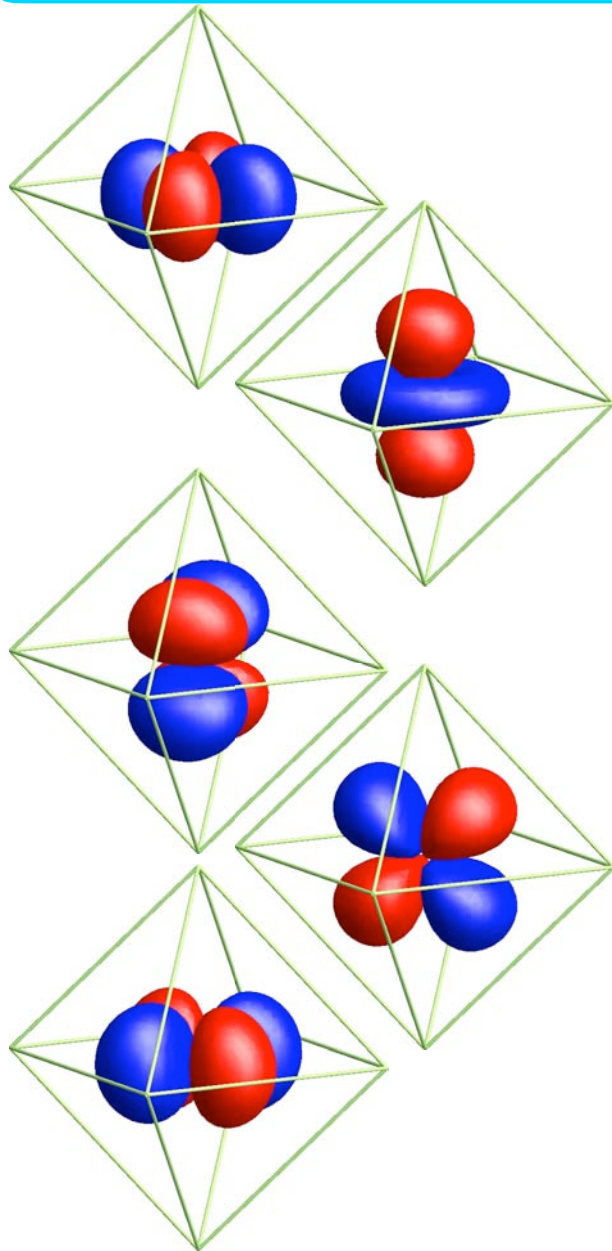
Assume interactions are well treated within the LDA



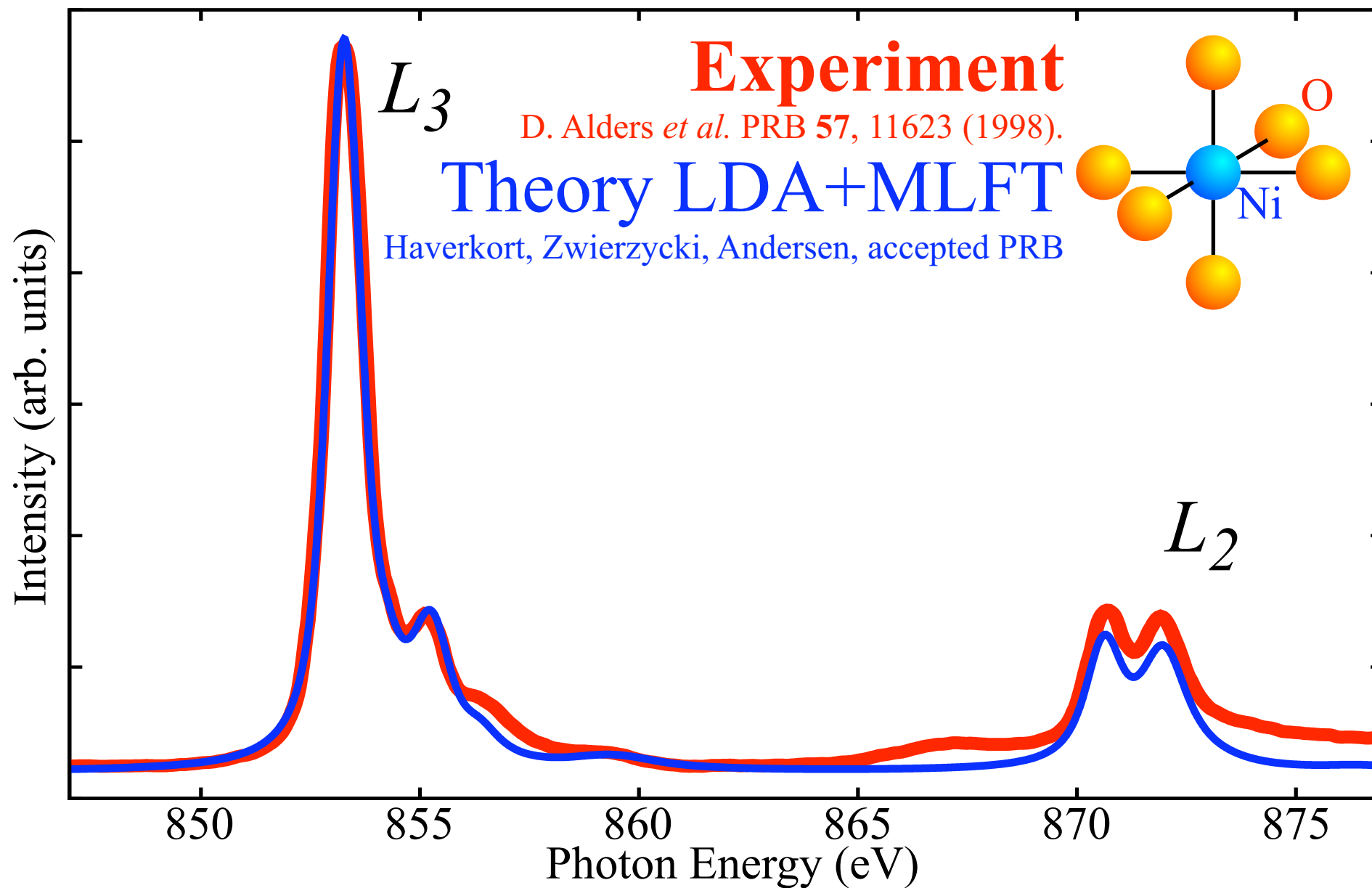
Two particle interactions on a level beyond LDA

TM 3d orbitals

Include the full Coulomb interaction



2p X-ray Absorption Spectroscopy *ab initio* theory



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

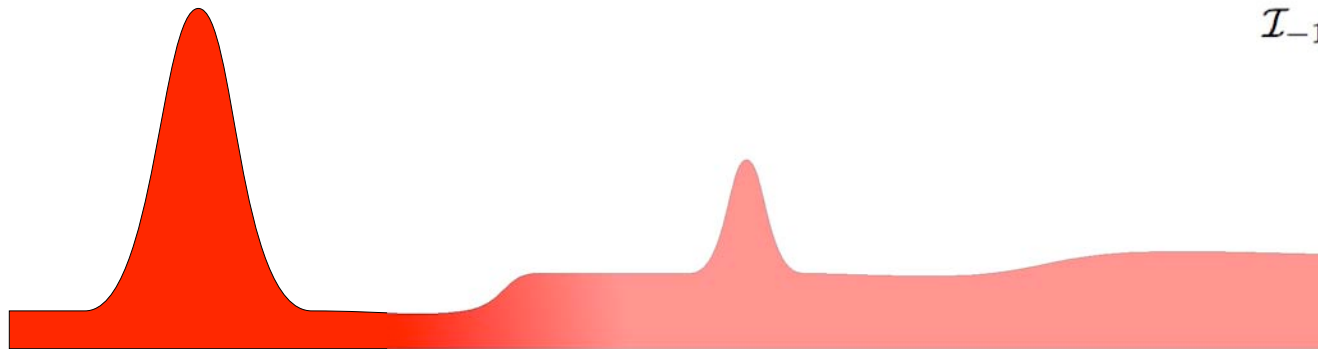
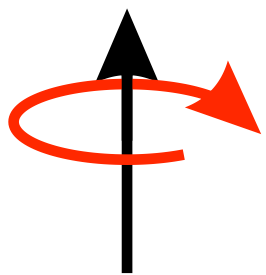
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_{1}^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_{1}^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_{1} = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

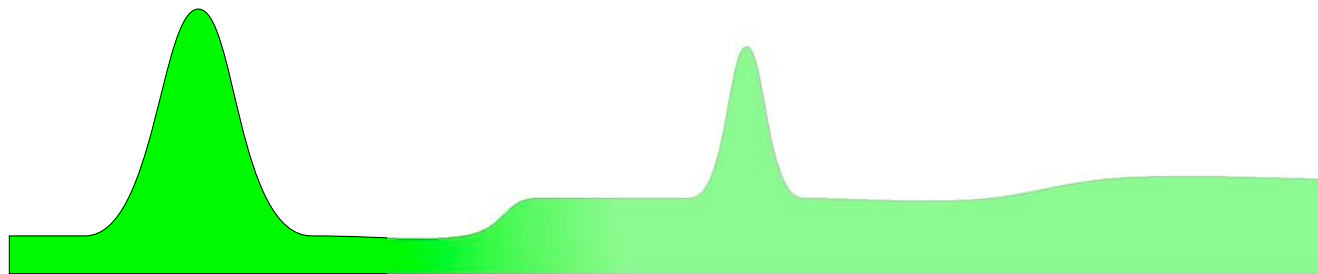
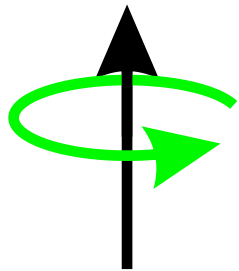
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_1^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_1^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_1 = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

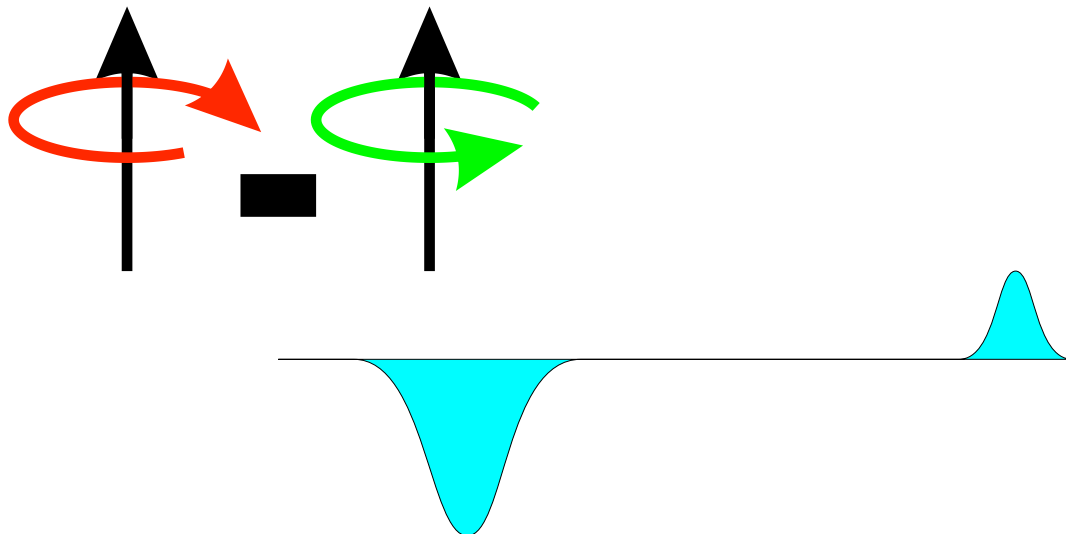
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_{1}^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_{1}^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_{1} = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

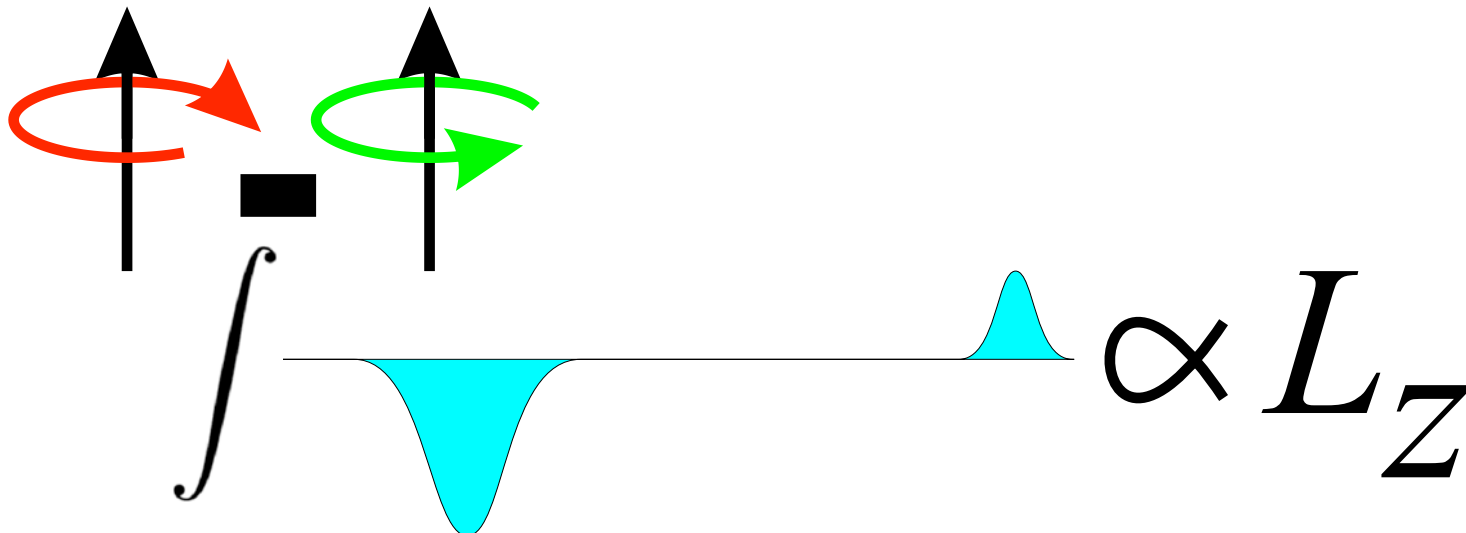
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_{1}^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_{1}^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_{1} = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

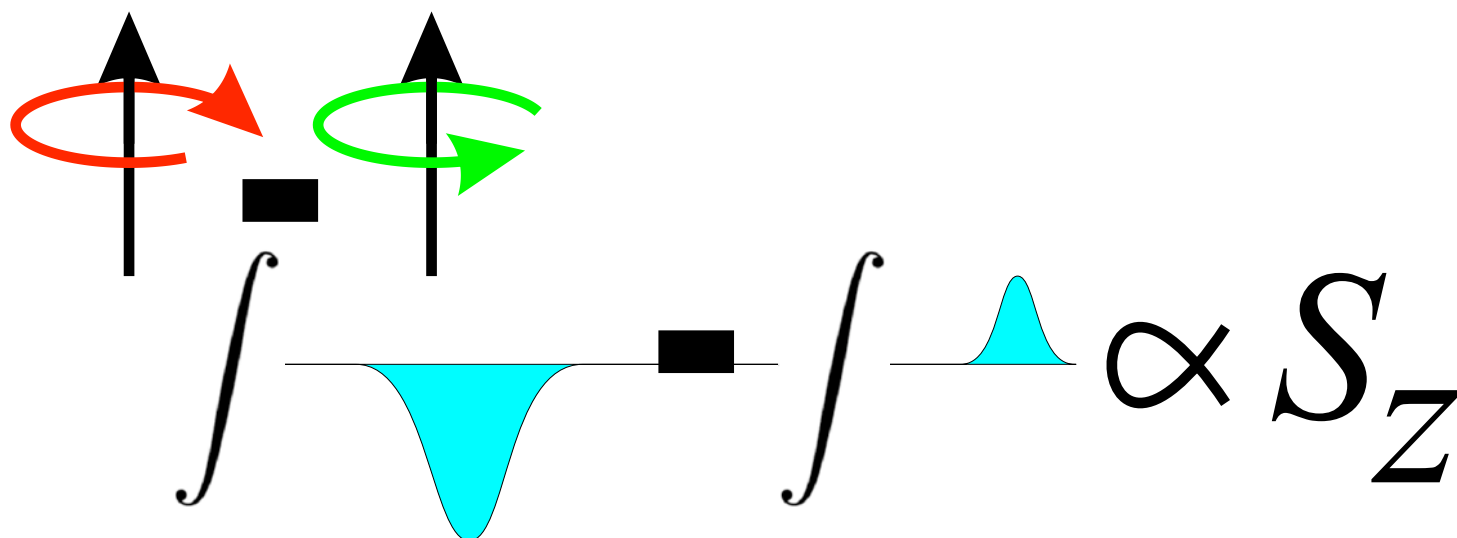
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_{1}^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_{1}^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_{1} = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

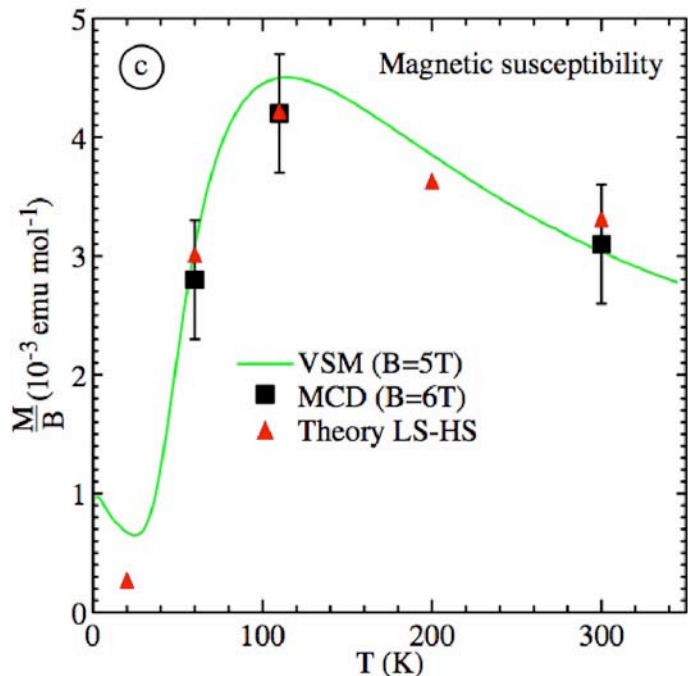
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

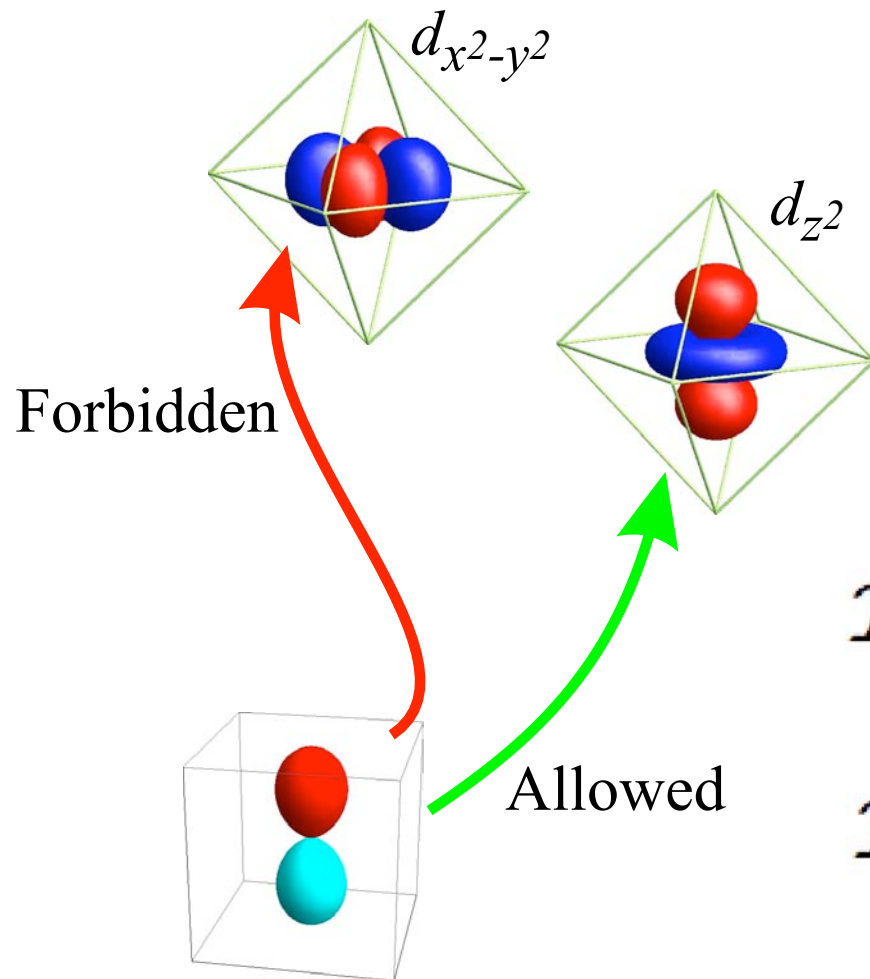
$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_{1}^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_{1}^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_{1} = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$



Orbital and spin resolved
magnetic susceptibility
(LaCoO_3)

XAS: Sum rules for linear dichroism orbital occupation

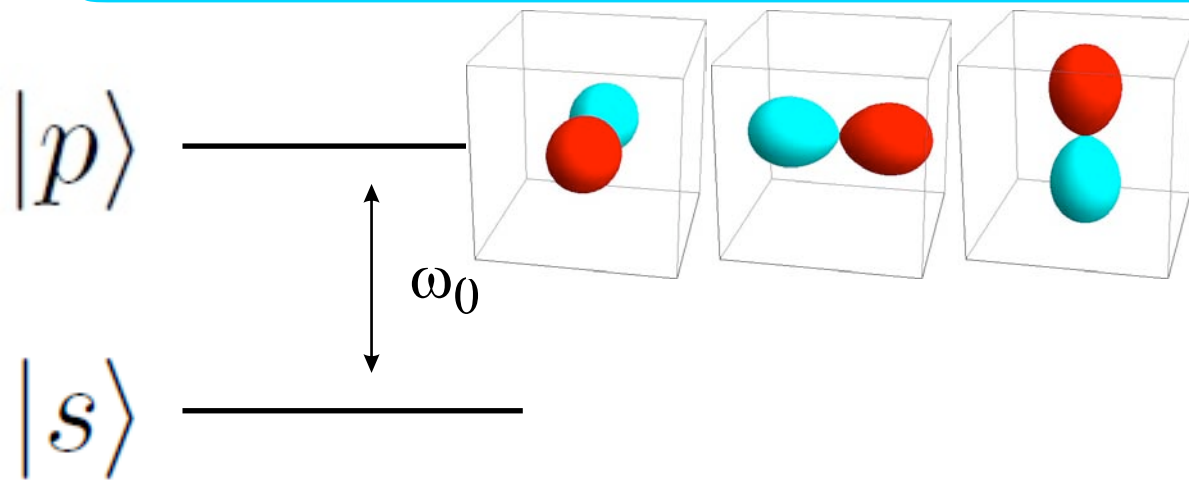


$$\mathcal{I}_x = \frac{1}{\underline{n}} \left(\frac{1}{2} \underline{n}_{xy} + \frac{1}{2} \underline{n}_{xz} + \frac{2}{3} \underline{n}_{x^2} \right)$$

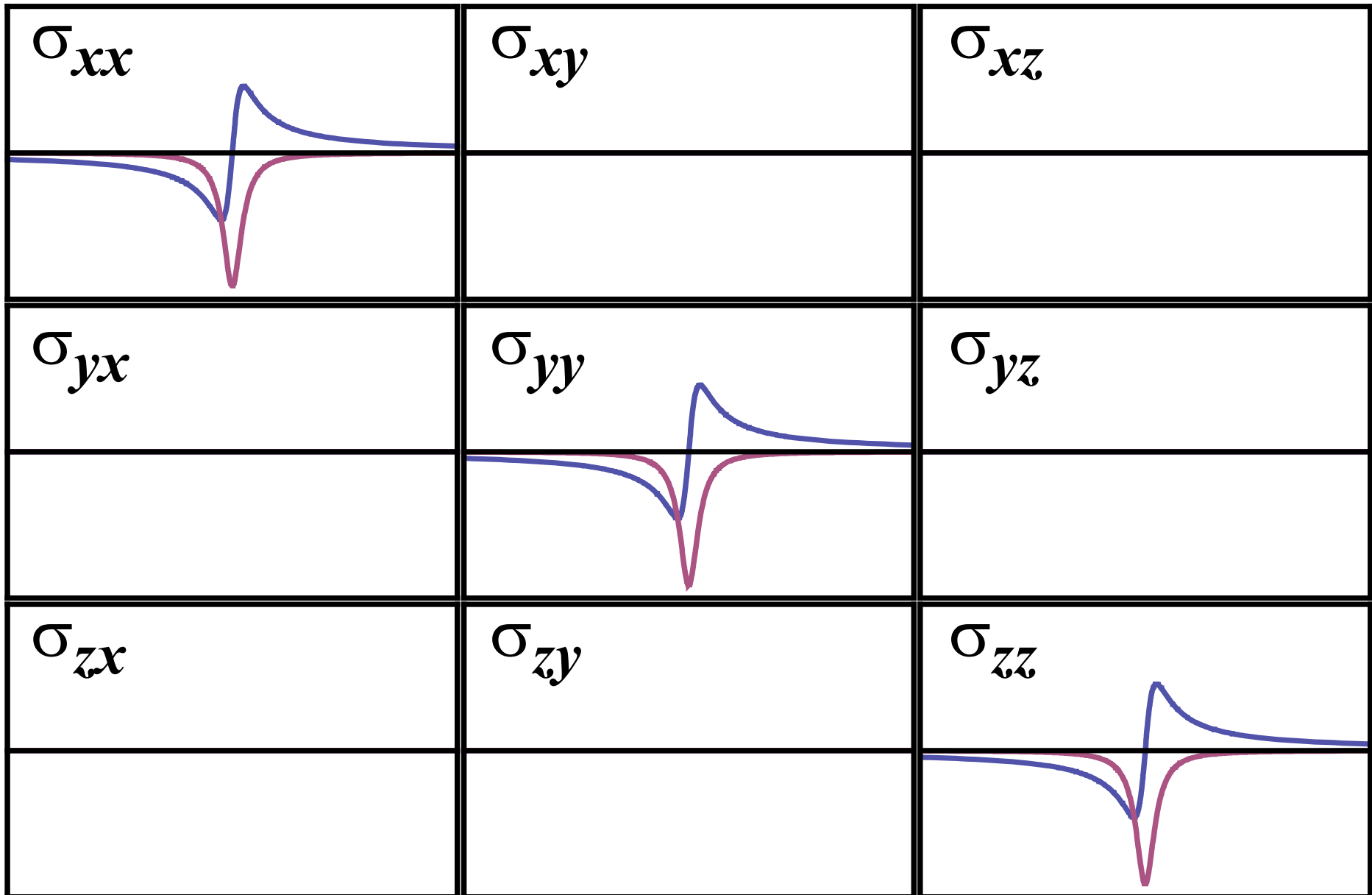
$$\mathcal{I}_y = \frac{1}{\underline{n}} \left(\frac{1}{2} \underline{n}_{xy} + \frac{1}{2} \underline{n}_{yz} + \frac{2}{3} \underline{n}_{y^2} \right)$$

$$\mathcal{I}_z = \frac{1}{\underline{n}} \left(\frac{1}{2} \underline{n}_{xz} + \frac{1}{2} \underline{n}_{yz} + \frac{2}{3} \underline{n}_{z^2} \right)$$

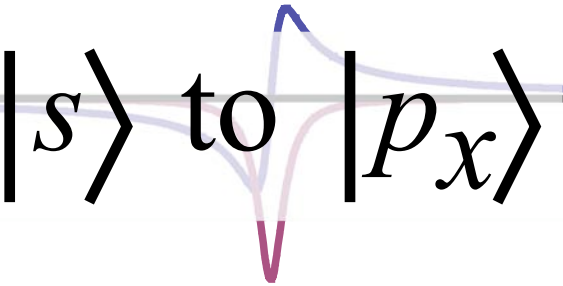
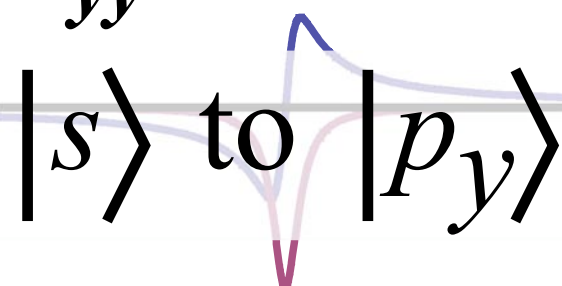
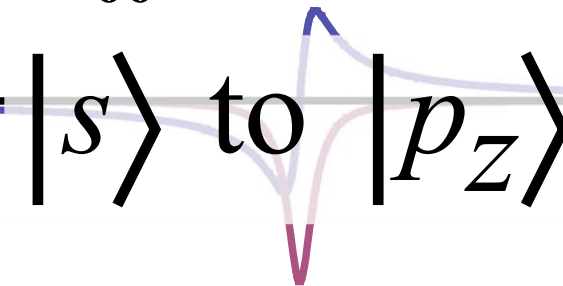
XAS and the conductivity tensor (σ)



But the conductivity tensor (σ) is a TENSOR

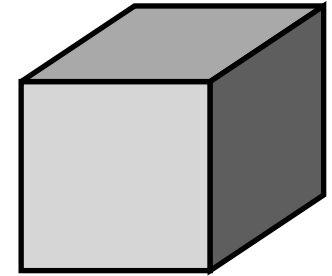
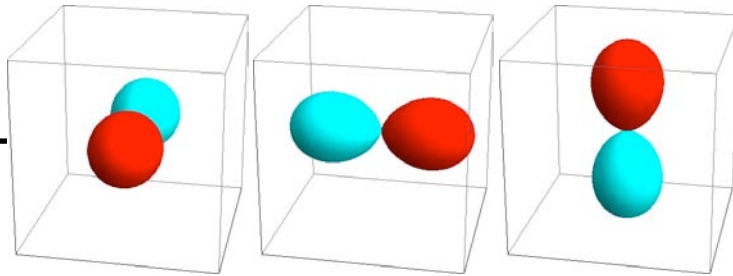


But the conductivity tensor (σ) is a TENSOR

σ_{xx} $ s\rangle$ to $ p_x\rangle$ 	σ_{xy}	σ_{xz}
σ_{yx}	σ_{yy} $ s\rangle$ to $ p_y\rangle$ 	σ_{yz}
σ_{zx}	σ_{zy}	σ_{zz} $ s\rangle$ to $ p_z\rangle$ 

the conductivity tensor (σ) in Cubic symmetry

$|p\rangle$

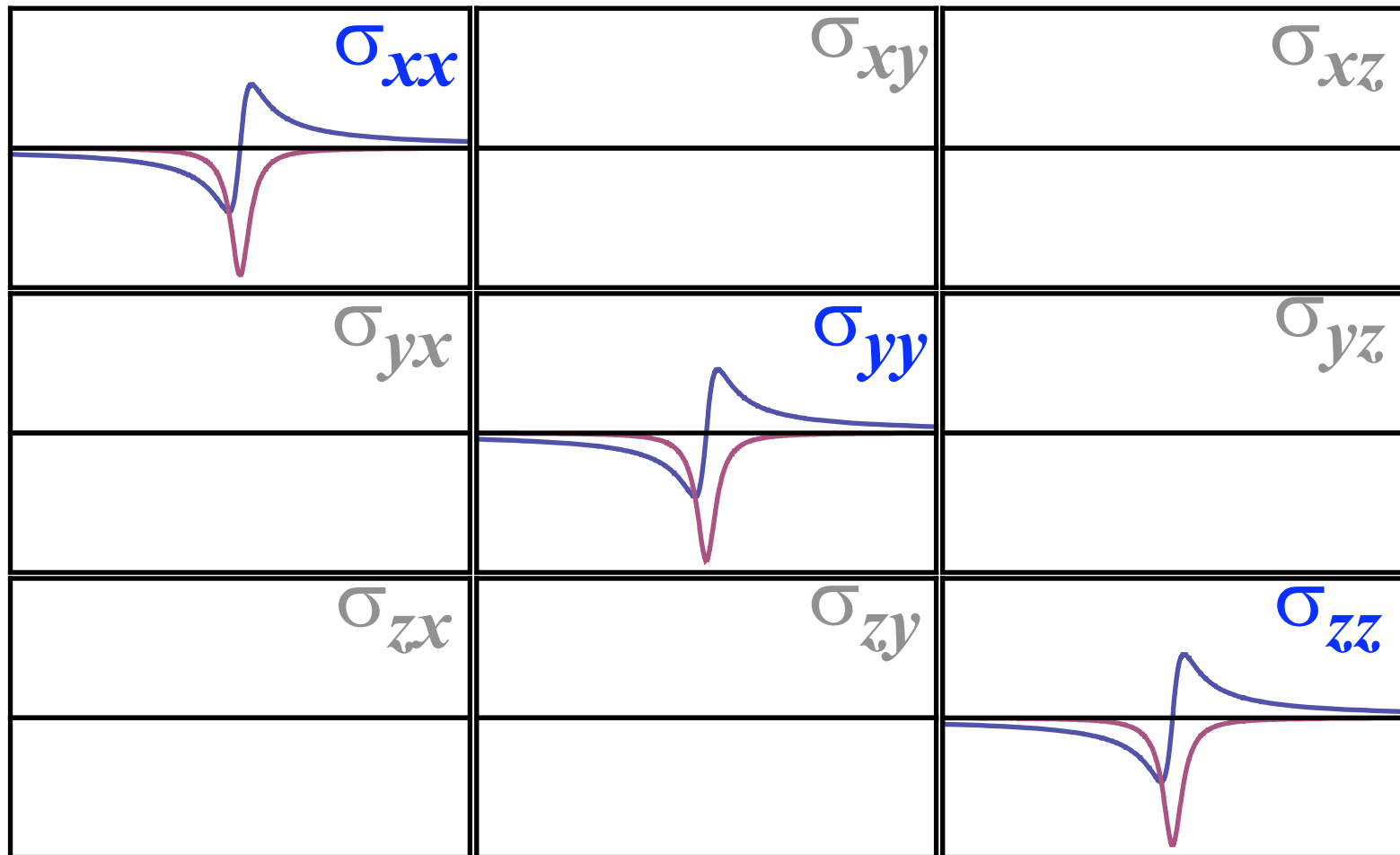


ω_0

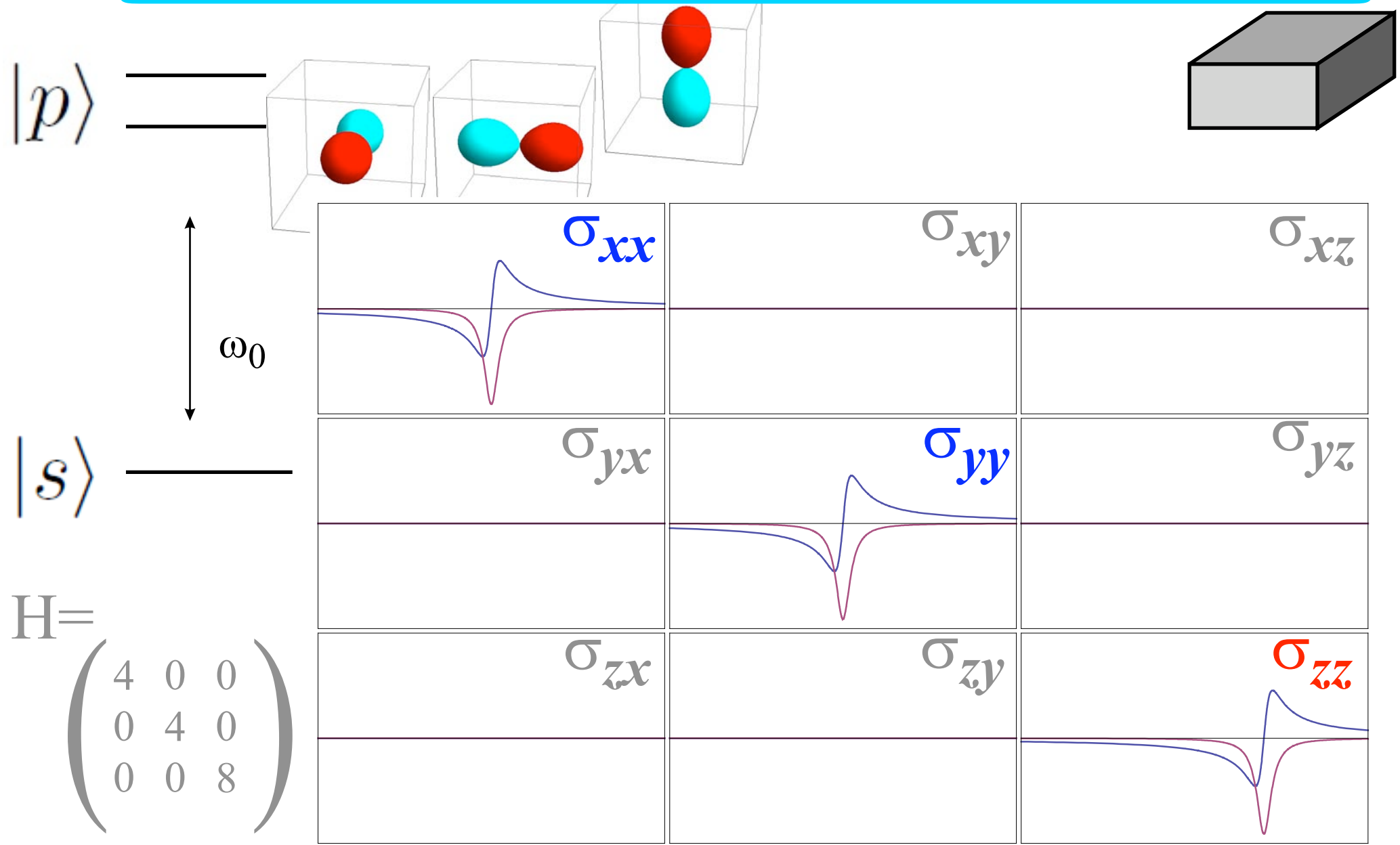
$|s\rangle$

H=

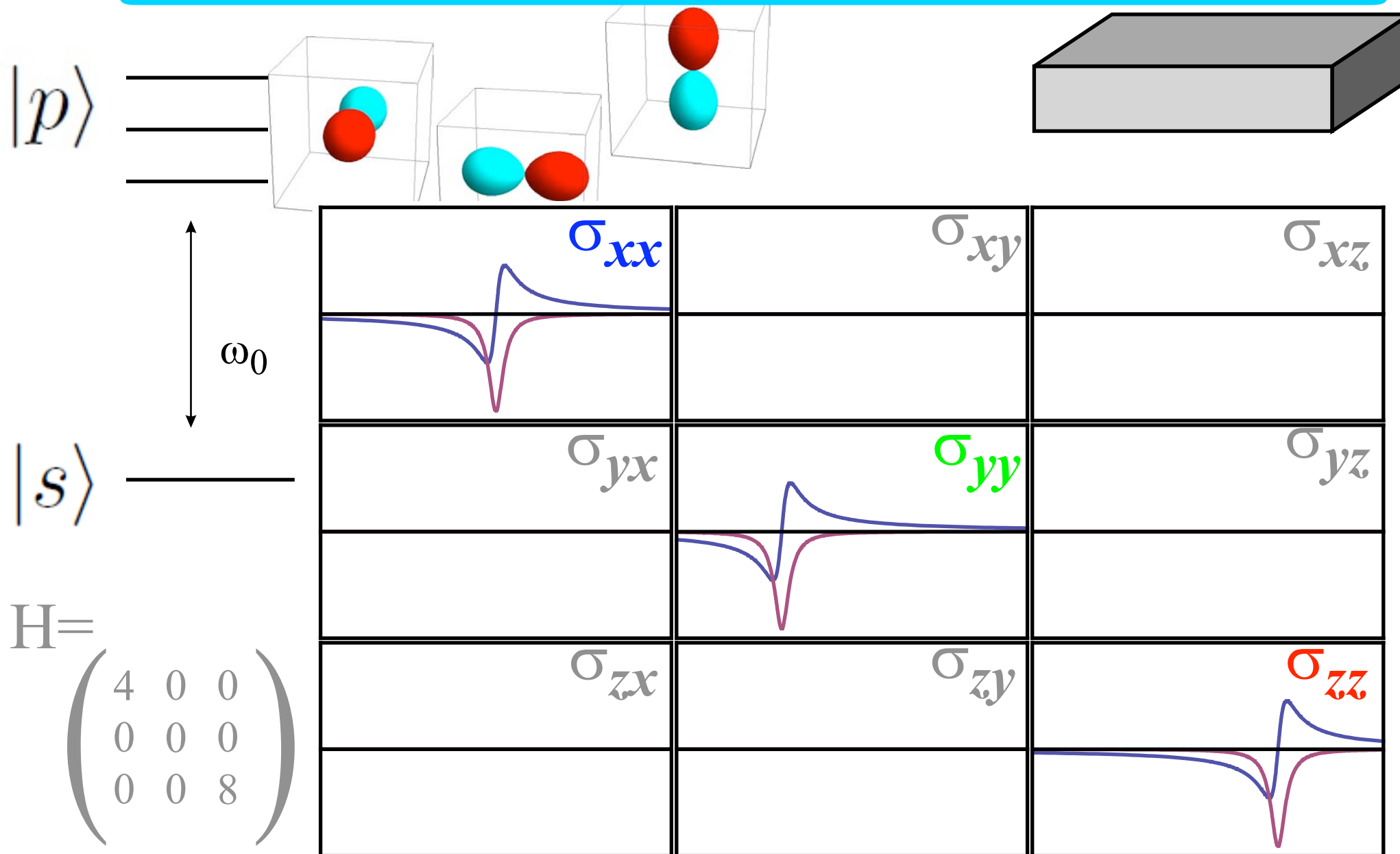
$$\begin{pmatrix} 4 & 0 & 0 \\ 0 & 4 & 0 \\ 0 & 0 & 4 \end{pmatrix}$$



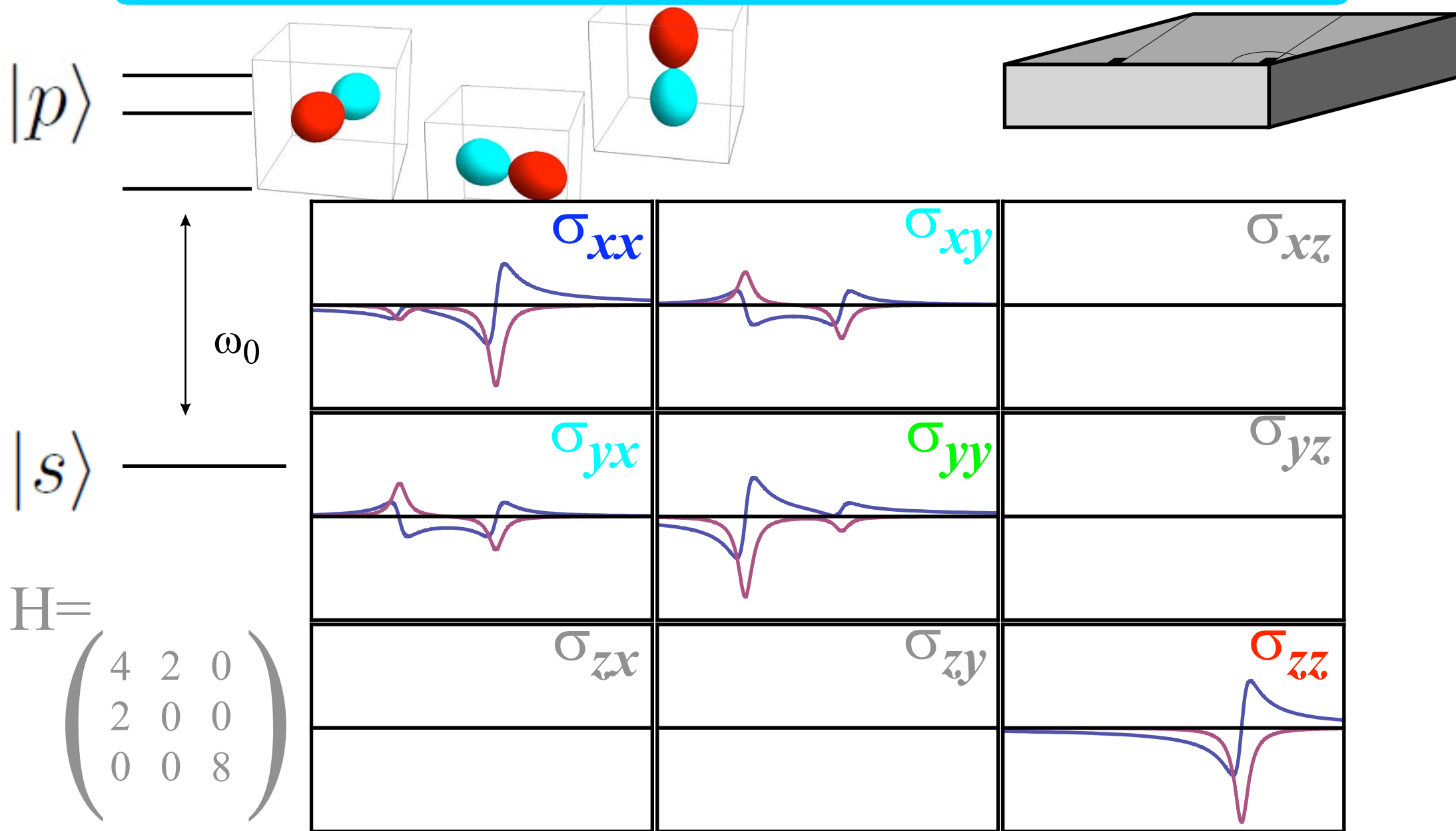
the conductivity tensor (σ) in **Tetragonal** symmetry



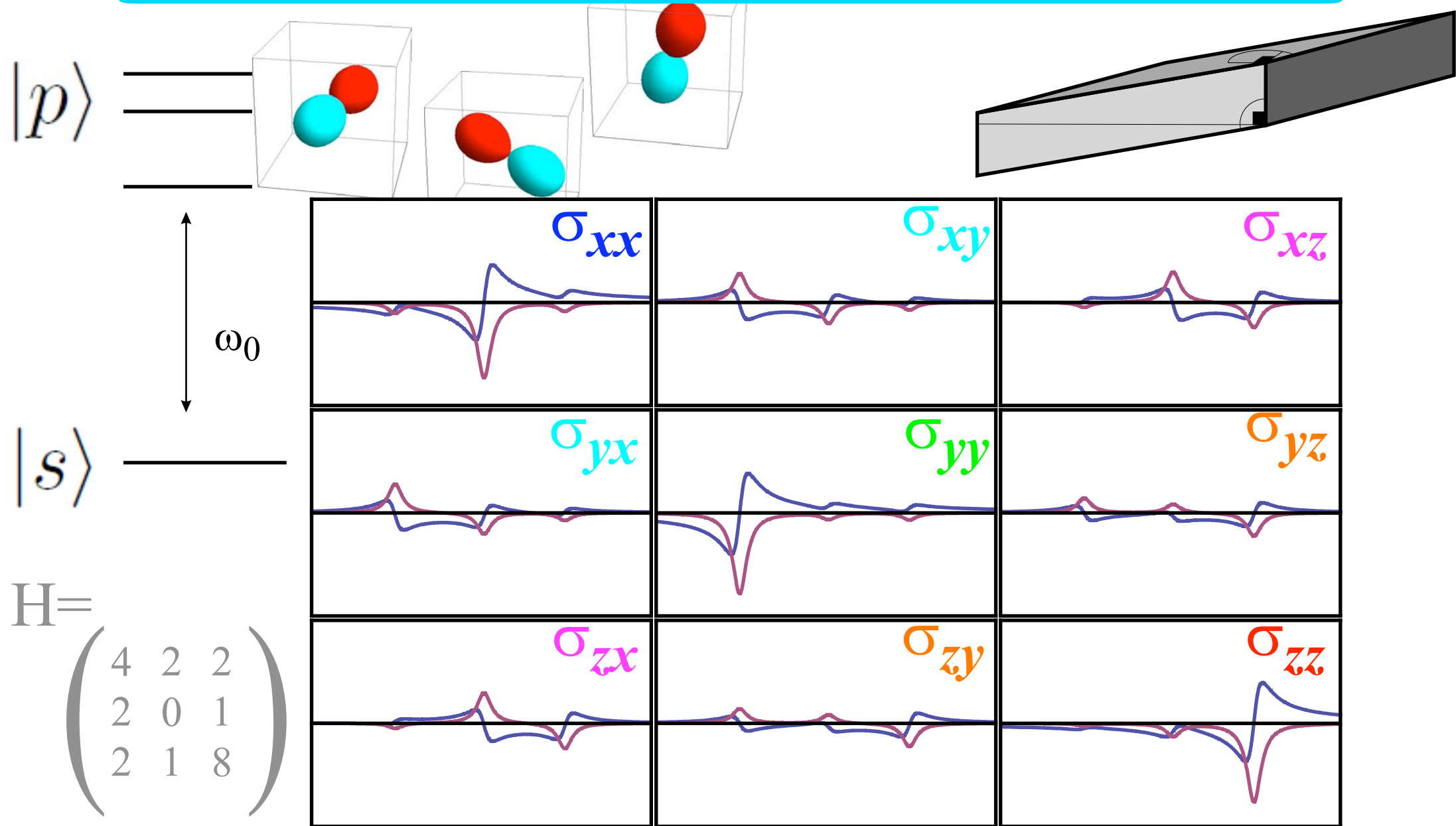
the conductivity tensor (σ) in Orthorhombic symmetry



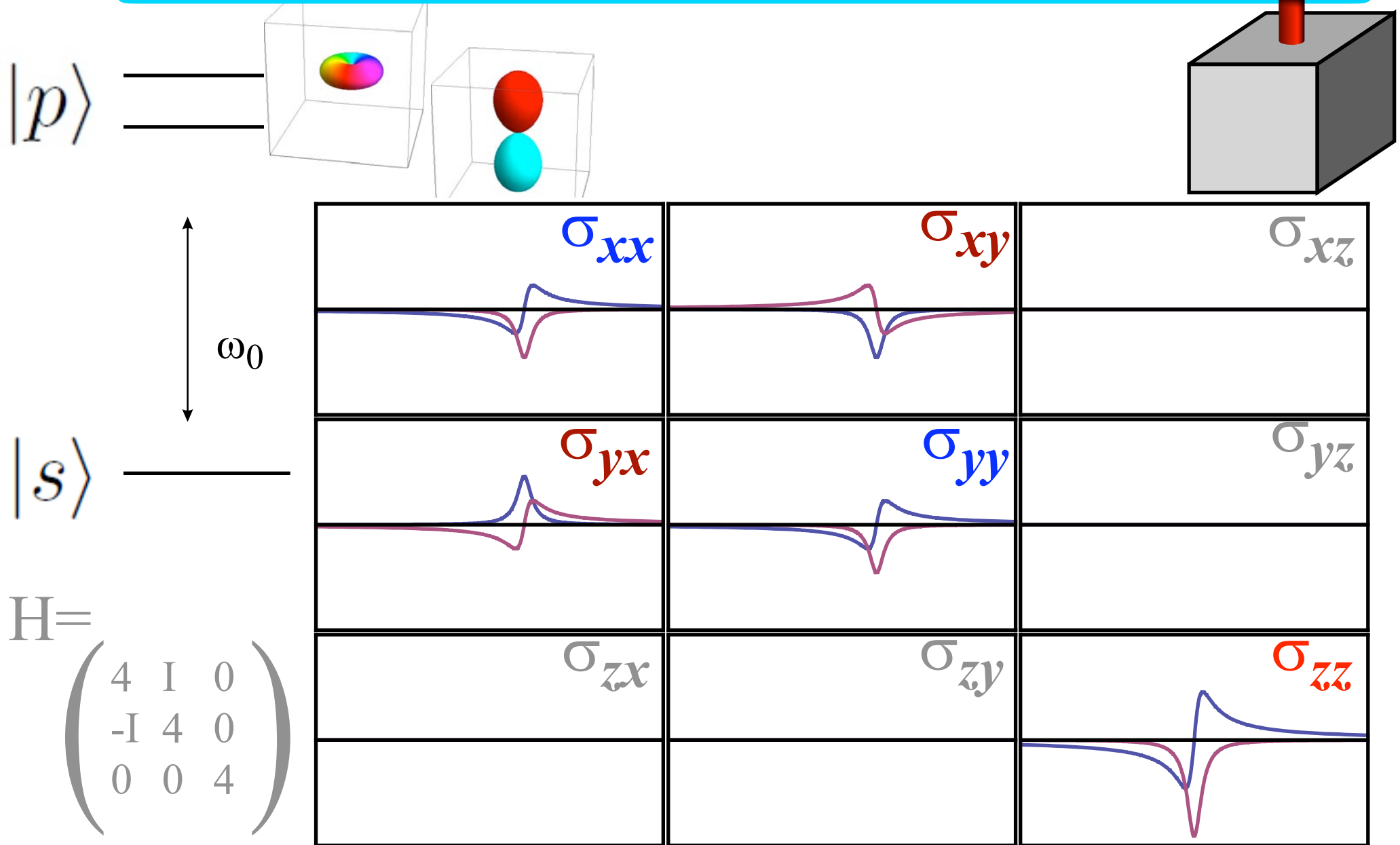
the conductivity tensor (σ) in **Monoclinic** symmetry



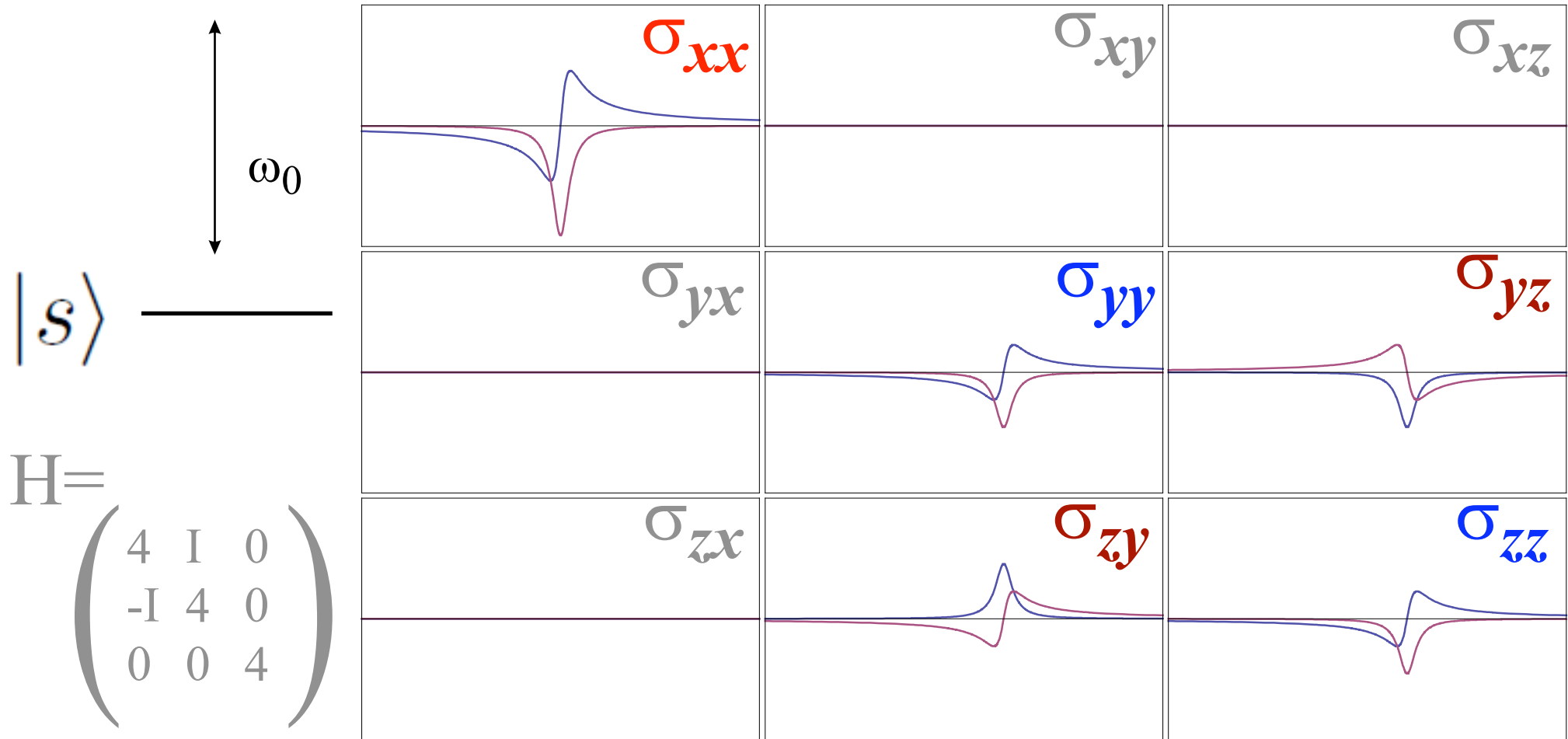
the conductivity tensor (σ) in *Triclinic* symmetry



the conductivity tensor (σ) in **Magnetic** materials



the conductivity tensor (σ) in **Magnetic** materials



XAS: Sum rules for circular dichroism orbital and spin moment

VOLUME 68, NUMBER 12

PHYSICAL REVIEW LETTERS

23 MARCH 1992

X-Ray Circular Dichroism as a Probe of Orbital Magnetization

B. T. Thole,⁽¹⁾ Paolo Carra,⁽²⁾ F. Sette,⁽²⁾ and G. van der Laan⁽³⁾

VOLUME 70, NUMBER 5

PHYSICAL REVIEW LETTERS

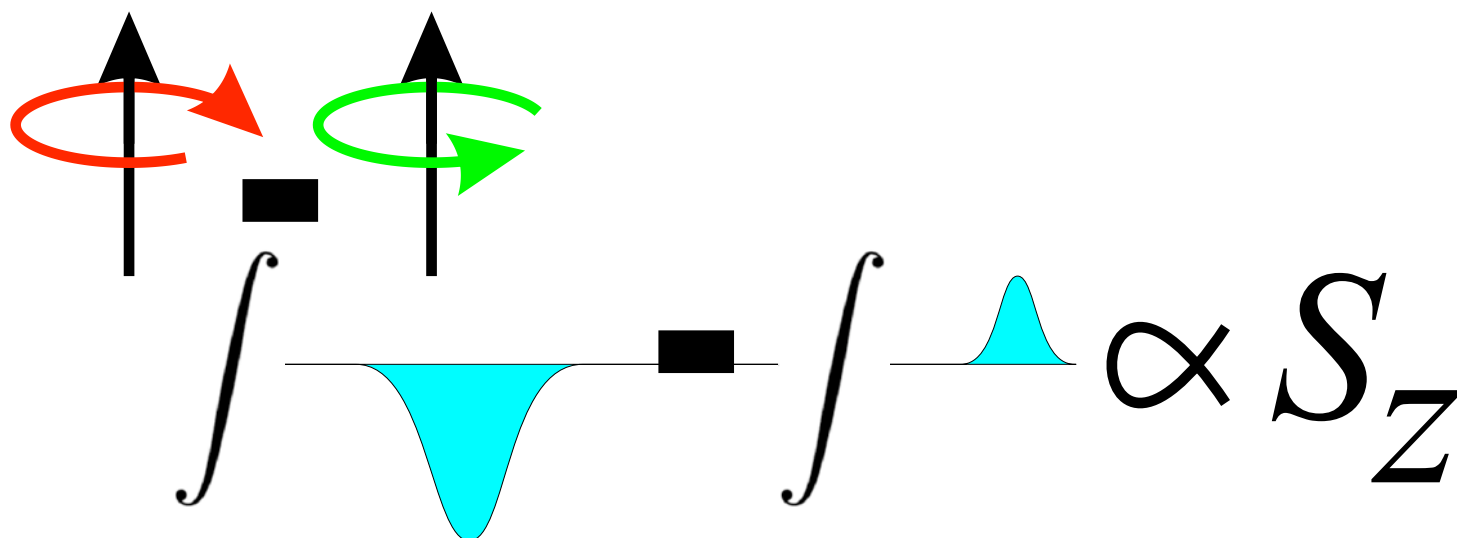
1 FEBRUARY 1993

X-Ray Circular Dichroism and Local Magnetic Fields

Paolo Carra,⁽¹⁾ B. T. Thole,^{(1),(2)} Massimo Altarelli,⁽¹⁾ and Xindong Wang⁽³⁾

$$(\mathcal{I}_{-1}^{c+\frac{1}{2}} - \mathcal{I}_{1}^{c+\frac{1}{2}}) - \frac{l}{l-1}(\mathcal{I}_{-1}^{c-\frac{1}{2}} - \mathcal{I}_{1}^{c-\frac{1}{2}}) = \frac{2}{3n} \mathbf{S}_z + \frac{2(2l+3)}{3ln} \mathbf{T}_z$$

$$\mathcal{I}_{-1} - \mathcal{I}_{1} = \frac{1}{n} \sum_{m,\sigma} n_{m\sigma} \frac{-m}{l} = \frac{\mathbf{L}_z}{ln}$$

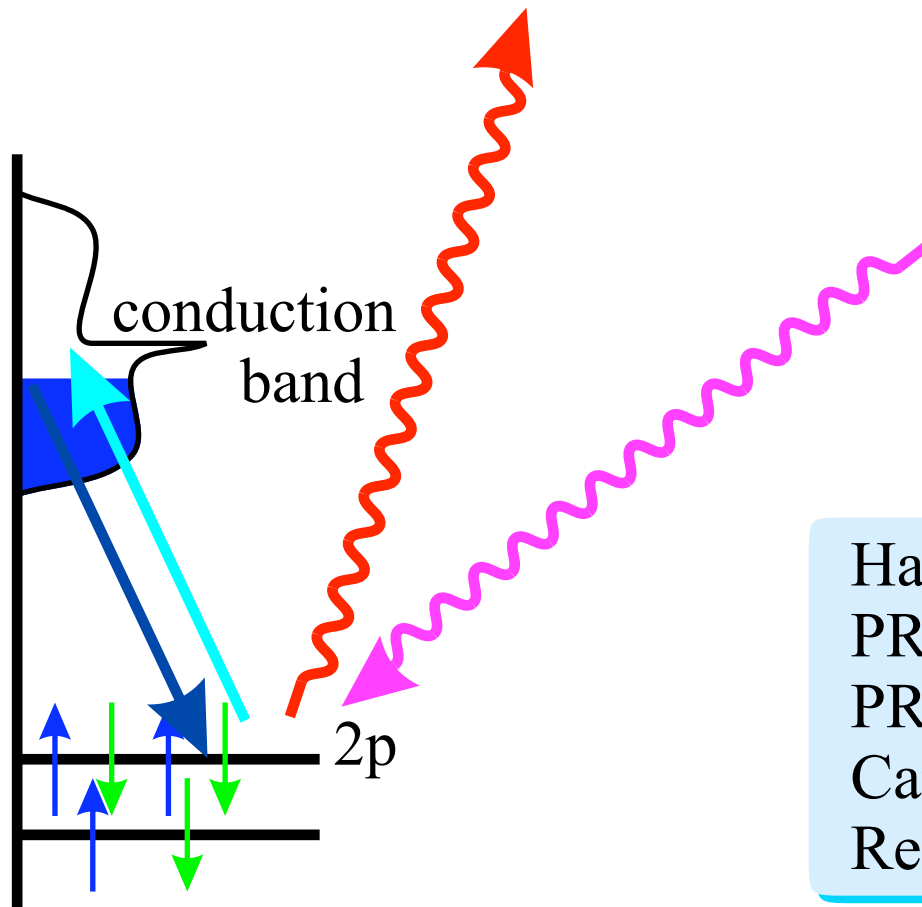


Thole *et al.*'s sum-rules in tensor form: s to p excitations

$$\int_{\mathbf{K}} \sigma = i \overline{\rho}_p$$

And now momentum, i.e. spatial resolved, from XAS to RXD

- ❑ X-ray Absorption Spectroscopy (XAS) at the TM $2p$ to $3d$ edge
- ❑ Resonant X-ray Diffraction (RXD)



Hannon, Trammel, Bloom, Gibbs
PRL **61**, 1245 (1988)
PRB **43**, 5663 (1991)
Carra, Thole
Rev. Mod. Phys. **66**, 1509 (1994).

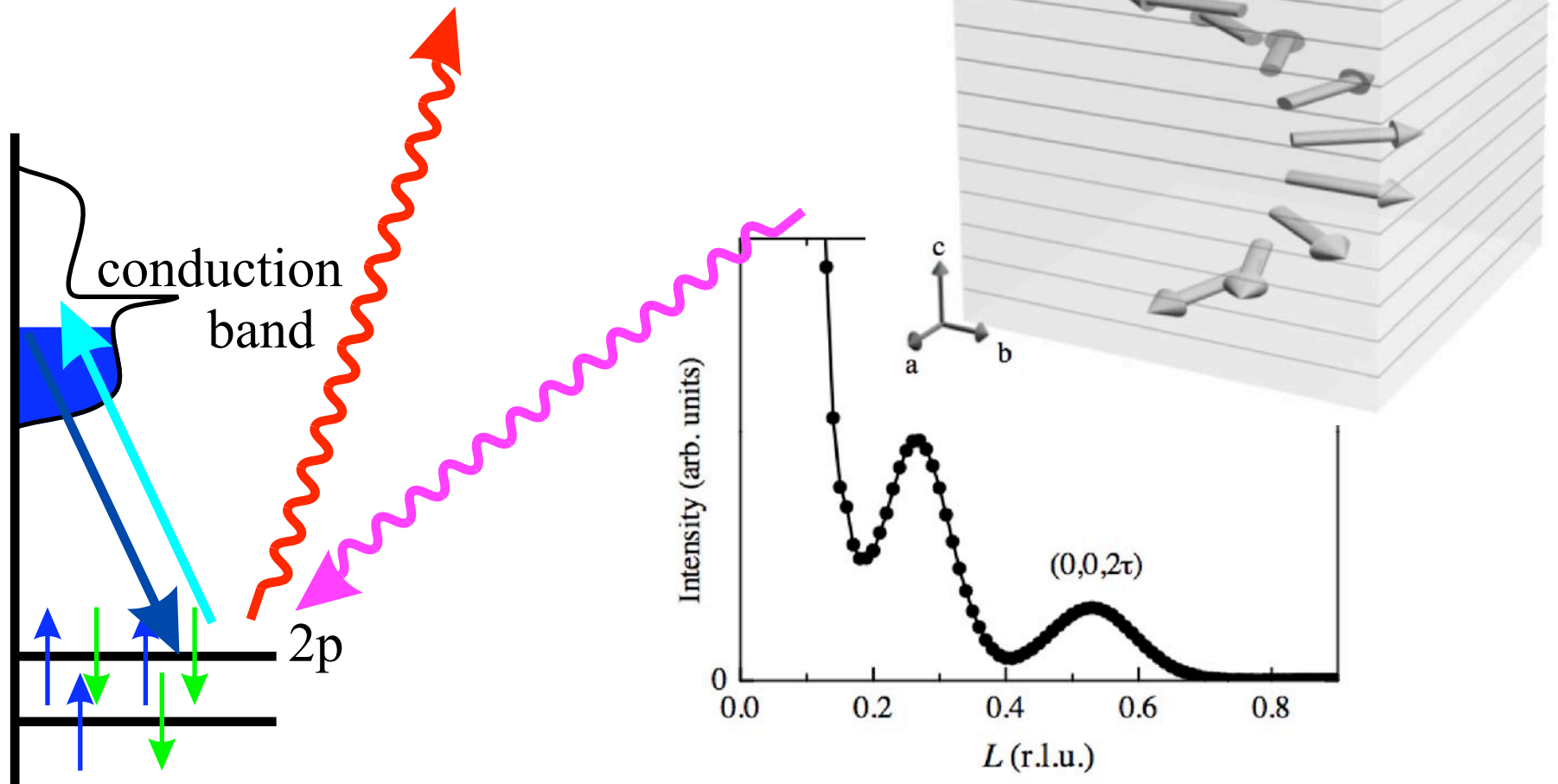
RXD: Magnetic Bragg reflection

Magnetic x-ray scattering at the M_5 absorption edge of Ho

Ho metal film

H. Ott,^{1,*} C. Schüßler-Langeheine,^{1,2} E. Schierle,¹ A. Yu. Grigoriev,^{1,†} V. Leiner,^{3,4,‡} H. Zabel,³
G. Kaindl,¹ and E. Weschke¹

PHYSICAL REVIEW B 74, 094412 (2006)



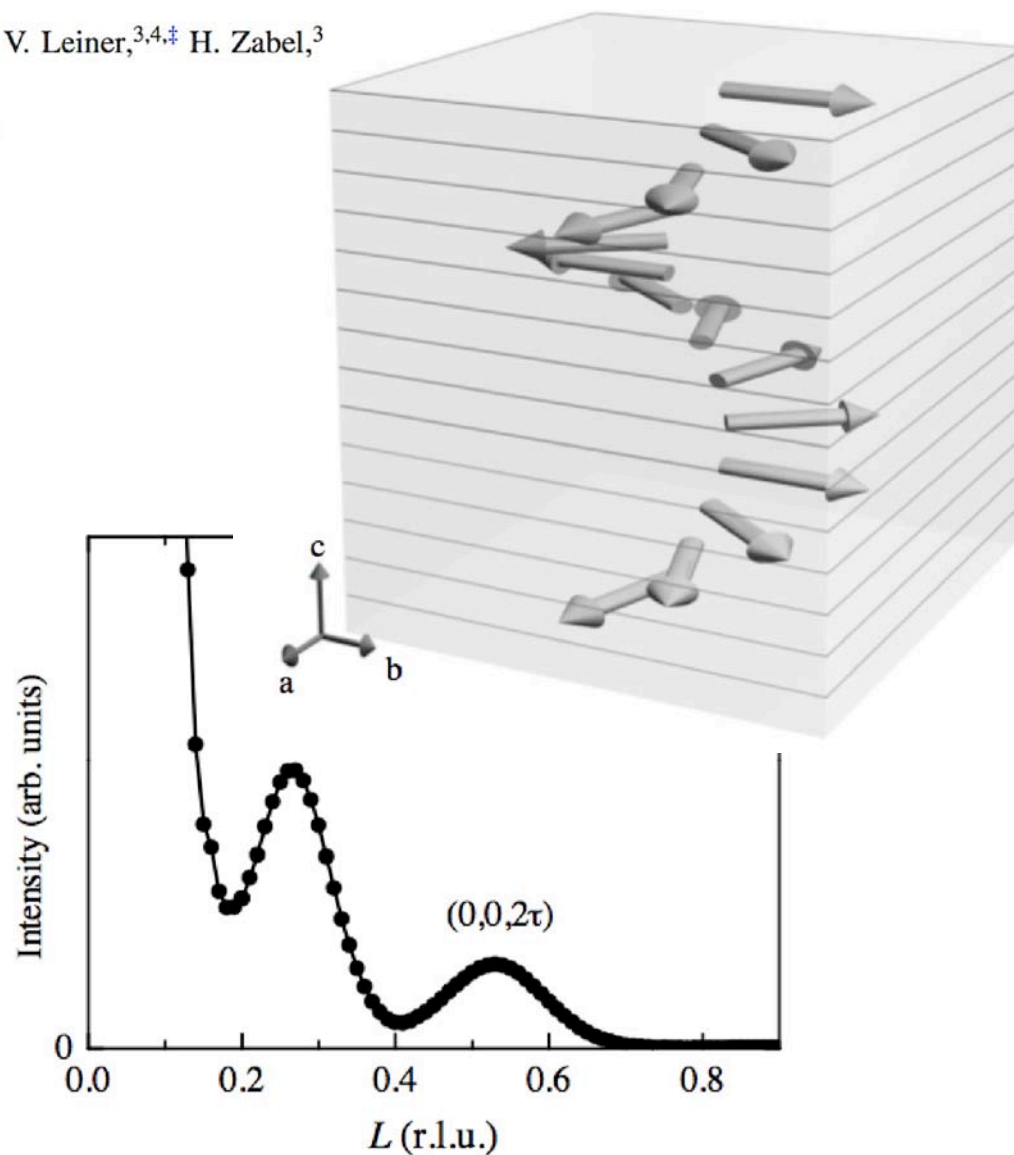
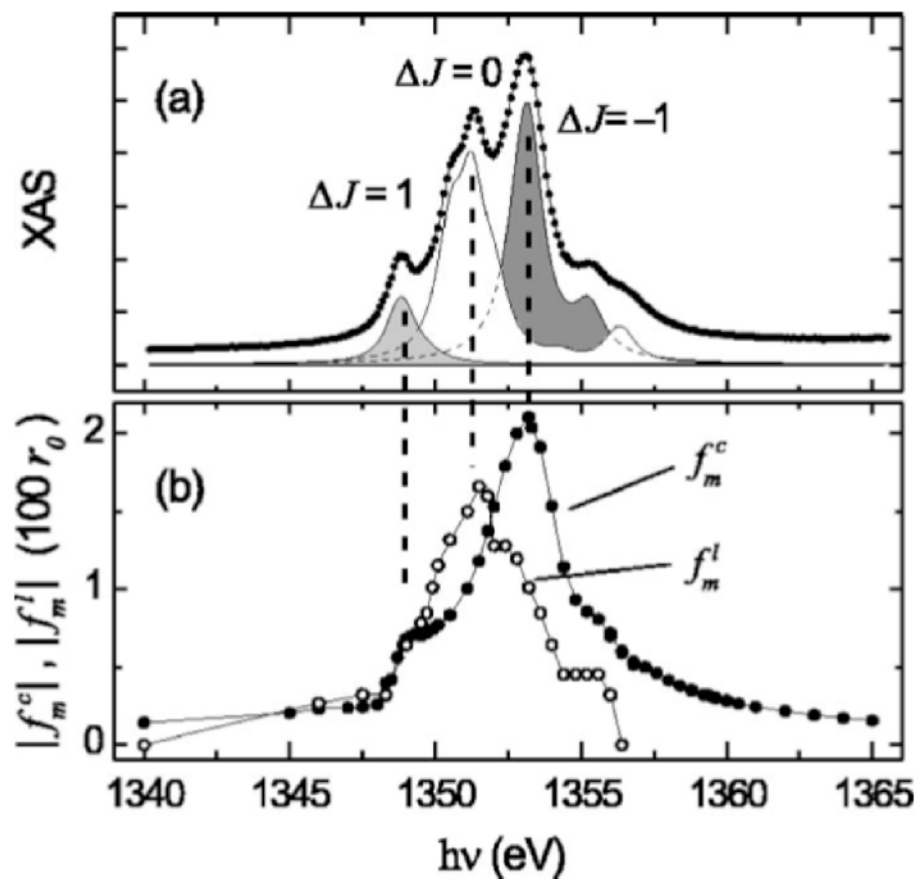
RXD: Magnetic Bragg reflection

Magnetic x-ray scattering at the M_5 absorption edge of Ho

Ho metal film

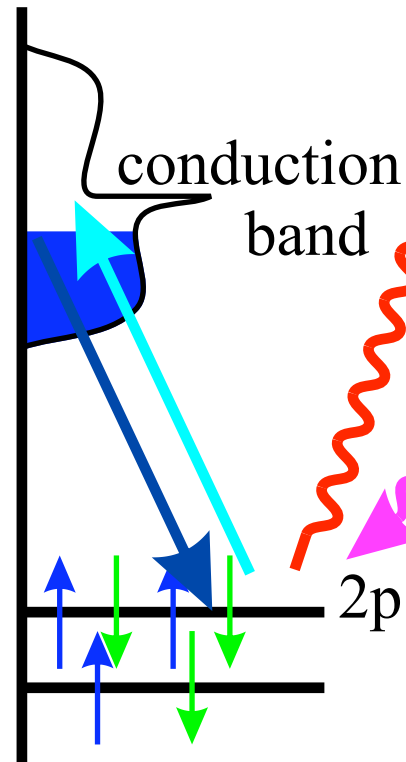
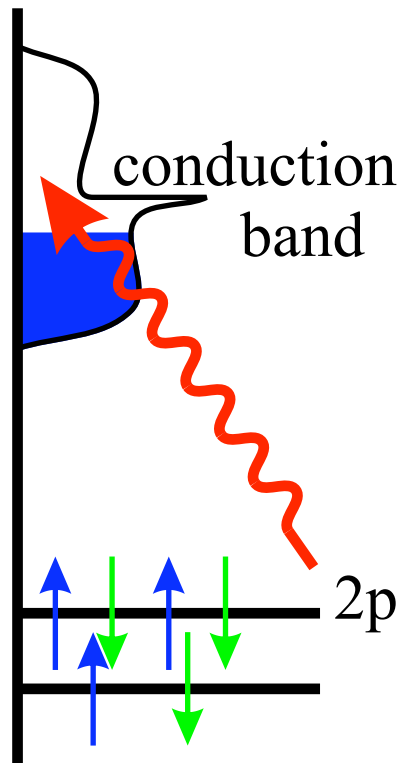
H. Ott,^{1,*} C. Schüßler-Langeheine,^{1,2} E. Schierle,¹ A. Yu. Grigoriev,^{1,†} V. Leiner,^{3,4,‡} H. Zabel,³
G. Kaindl,¹ and E. Weschke¹

PHYSICAL REVIEW B 74, 094412 (2006)

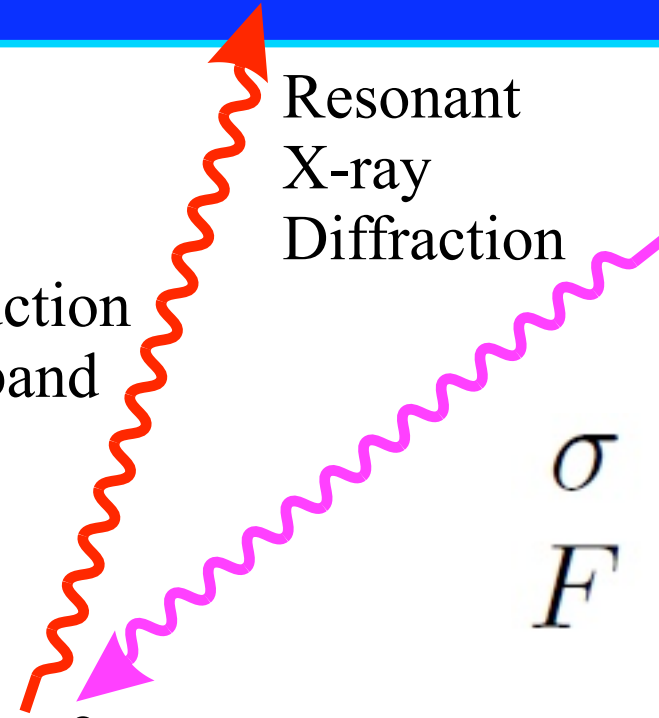


Relation between resonant x-ray diffraction and absorption

X-ray
Absorption
Spectroscopy



Resonant
X-ray
Diffraction



$$\sigma = \omega \chi$$

$$F = \omega \sigma$$

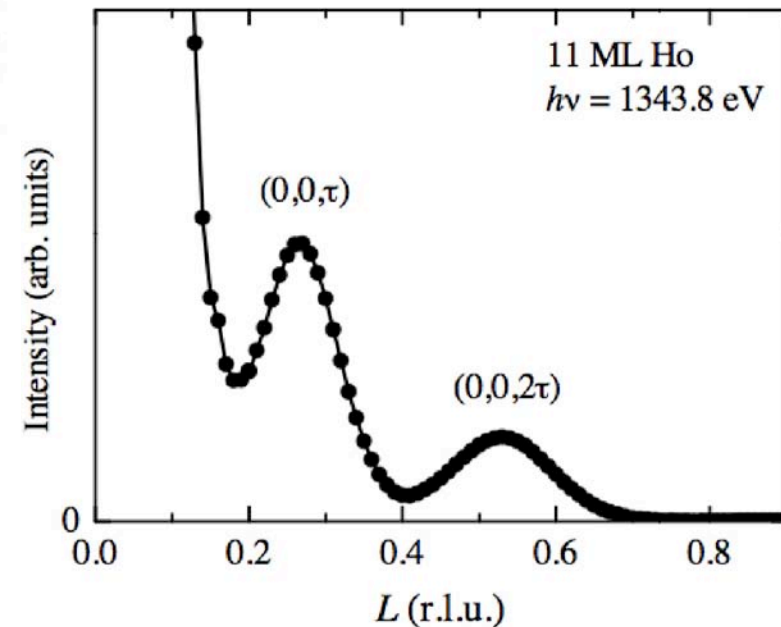
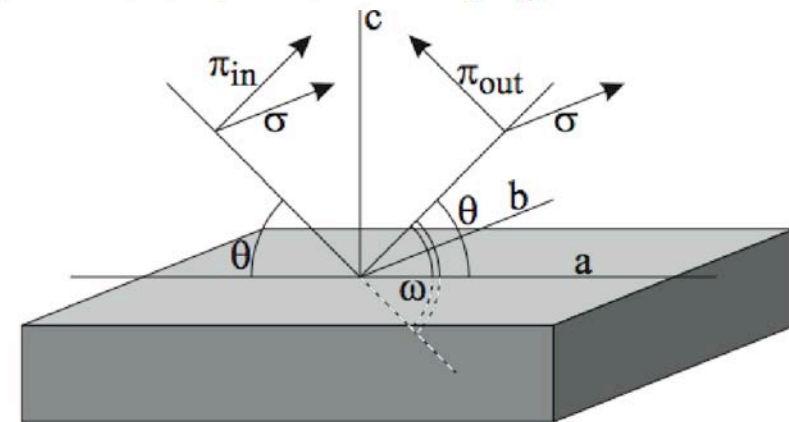
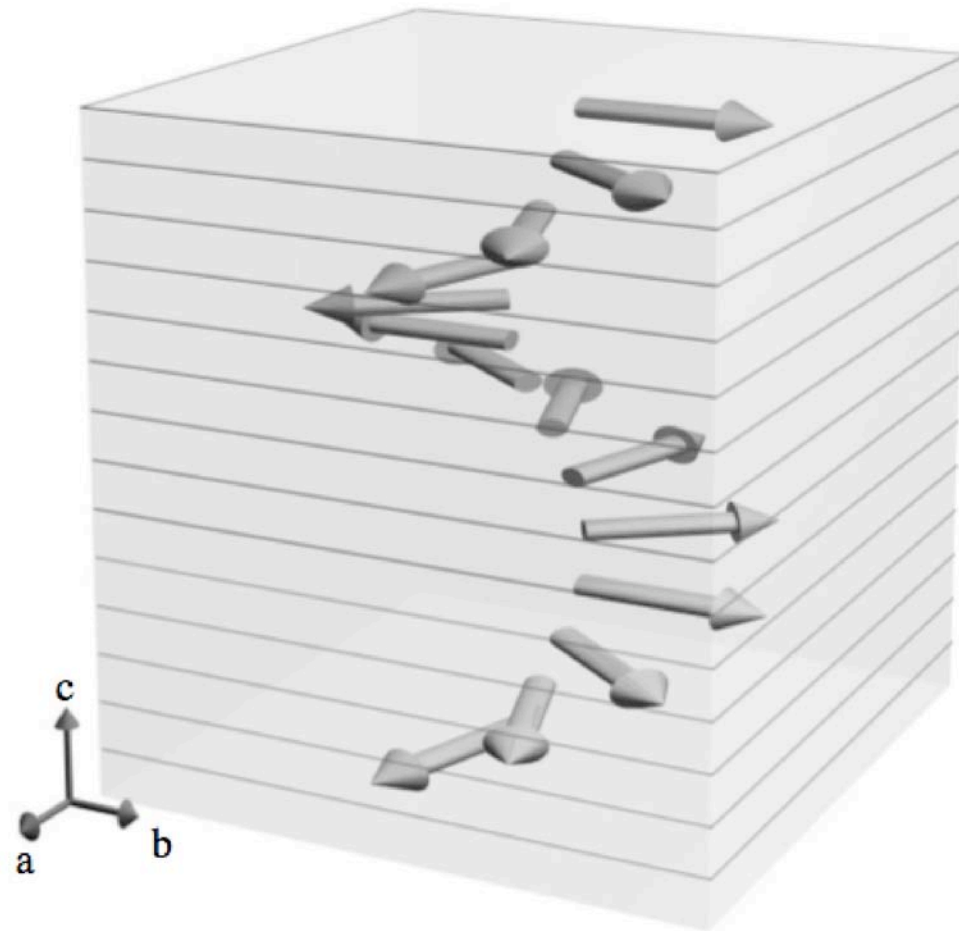
$$I_{XAS} = -\text{Im}[\hat{\epsilon}^* \cdot \sigma \cdot \hat{\epsilon}]$$

$$I_{RXD} = |\hat{\epsilon}_{out}^* \cdot F \cdot \hat{\epsilon}_{in}|^2$$

Example on Ho

Atomic multiplet calculation of $3d_{5/2} \rightarrow 4f$ resonant x-ray diffraction from Ho metal

M. W. Haverkort,^{1,2} C. Schüßler-Langeheine,¹ C. F. Chang,¹ M. Buchholz,¹
H.-H. Wu,¹ H. Ott,¹ E. Schierle,³ D. Schmitz,³ A. Tanaka,⁴ and L. H. Tjeng¹



Fundamental spectra from the literature

PHYSICAL REVIEW B

VOLUME 37, NUMBER 4

1 FEBRUARY 1988

Calculations of magnetic x-ray dichroism in the 3d absorption spectra of rare-earth compounds

J. B. Goedkoop

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

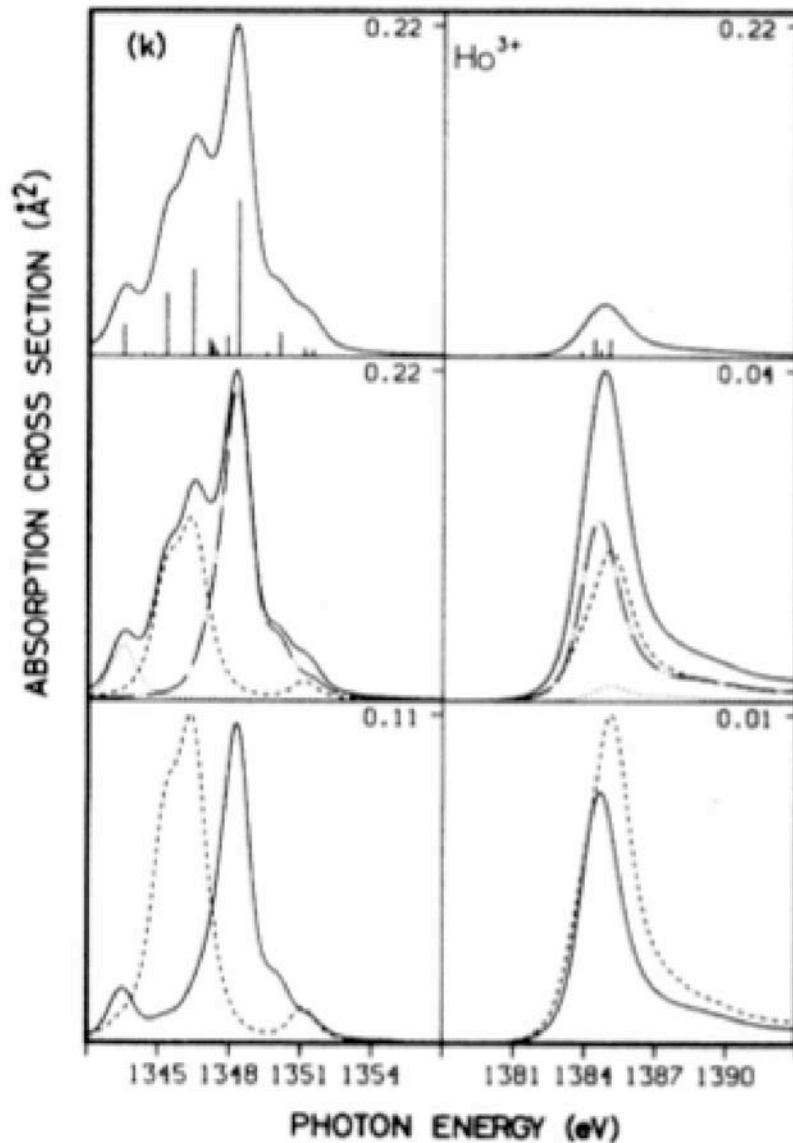
B. T. Thole, G. van der Laan, and G. A. Sawatzky

Department of Applied and Solid State Physics, Materials Science Center, University of Groningen, Nijenborgh 18, 9747 AG Groningen, The Netherlands

F. M. F. de Groot and J. C. Fuggle

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

(Received 13 April 1987)



Plotted is $-\text{Im}[\sigma]$

top: $\sigma^{(0)}$

middle $\sigma^{(1)}$ (difference between the two spectra)

bottom $\sigma^{(2)}$ (difference between the two spectra)

Fundamental spectra from the literature

PHYSICAL REVIEW B

VOLUME 37, NUMBER 4

1 FEBRUARY 1988

Calculations of magnetic x-ray dichroism in the 3d absorption spectra of rare-earth compounds

J. B. Goedkoop

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

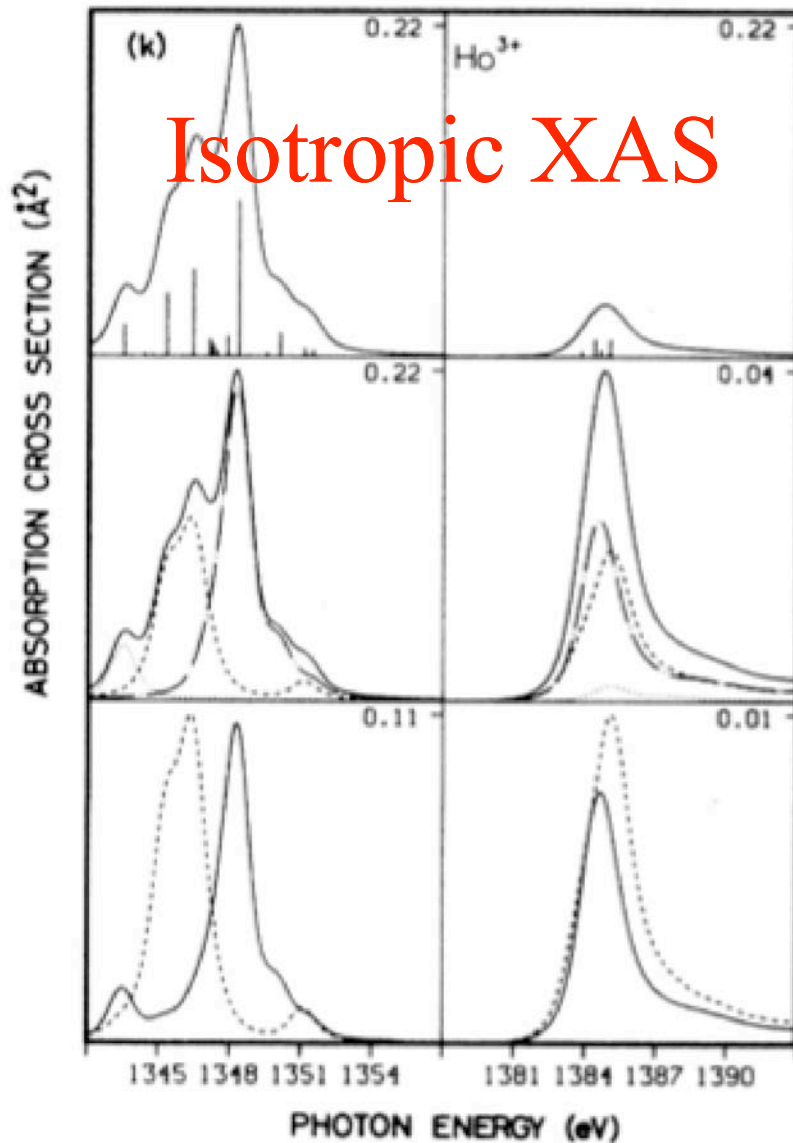
B. T. Thole, G. van der Laan, and G. A. Sawatzky

Department of Applied and Solid State Physics, Materials Science Center, University of Groningen, Nijenborgh 18, 9747 AG Groningen, The Netherlands

F. M. F. de Groot and J. C. Fuggle

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

(Received 13 April 1987)



Plotted is $-\text{Im}[\sigma]$

top: $\sigma^{(0)}$

middle $\sigma^{(1)}$ (difference between the two spectra)

bottom $\sigma^{(2)}$ (difference between the two spectra)

Fundamental spectra from the literature

PHYSICAL REVIEW B

VOLUME 37, NUMBER 4

1 FEBRUARY 1988

Calculations of magnetic x-ray dichroism in the 3d absorption spectra of rare-earth compounds

J. B. Goedkoop

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

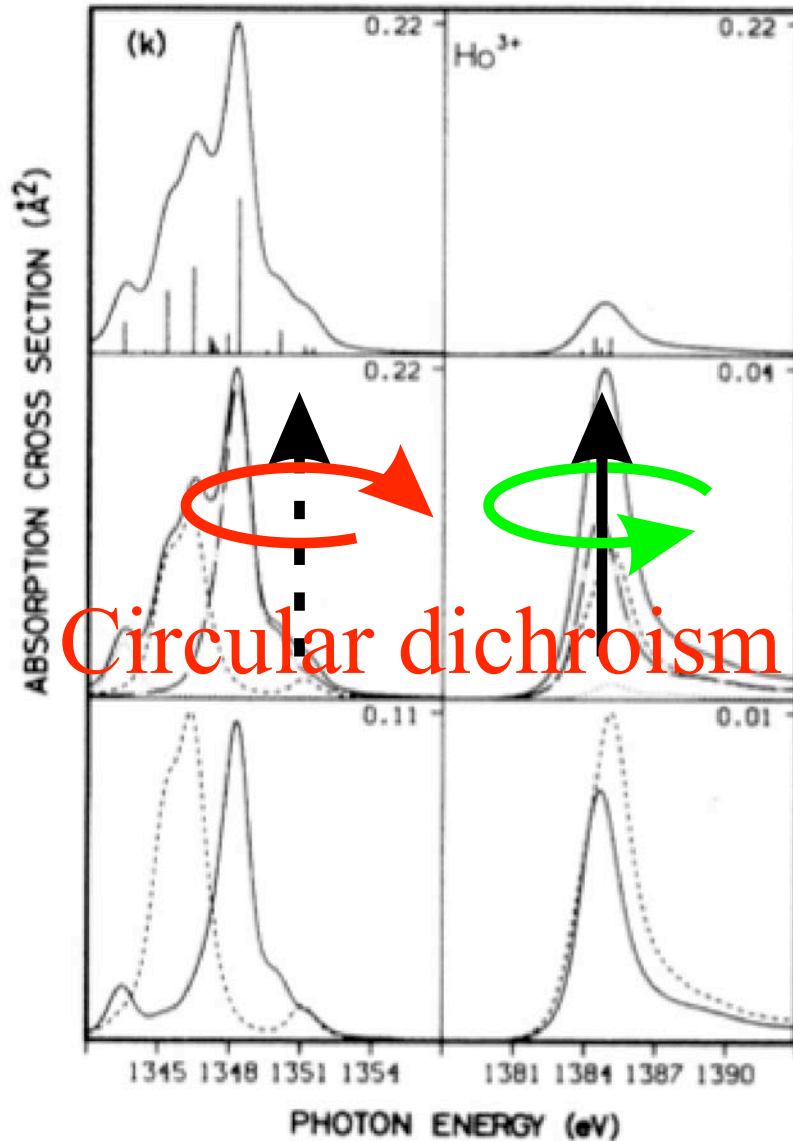
B. T. Thole, G. van der Laan, and G. A. Sawatzky

Department of Applied and Solid State Physics, Materials Science Center, University of Groningen, Nijenborgh 18, 9747 AG Groningen, The Netherlands

F. M. F. de Groot and J. C. Fuggle

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

(Received 13 April 1987)



Plotted is $-\text{Im}[\sigma]$

top: $\sigma^{(0)}$

middle $\sigma^{(1)}$ (difference between the two spectra)

bottom $\sigma^{(2)}$ (difference between the two spectra)

Circular dichroism

Fundamental spectra from the literature

PHYSICAL REVIEW B

VOLUME 37, NUMBER 4

1 FEBRUARY 1988

Calculations of magnetic x-ray dichroism in the 3d absorption spectra of rare-earth compounds

J. B. Goedkoop

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

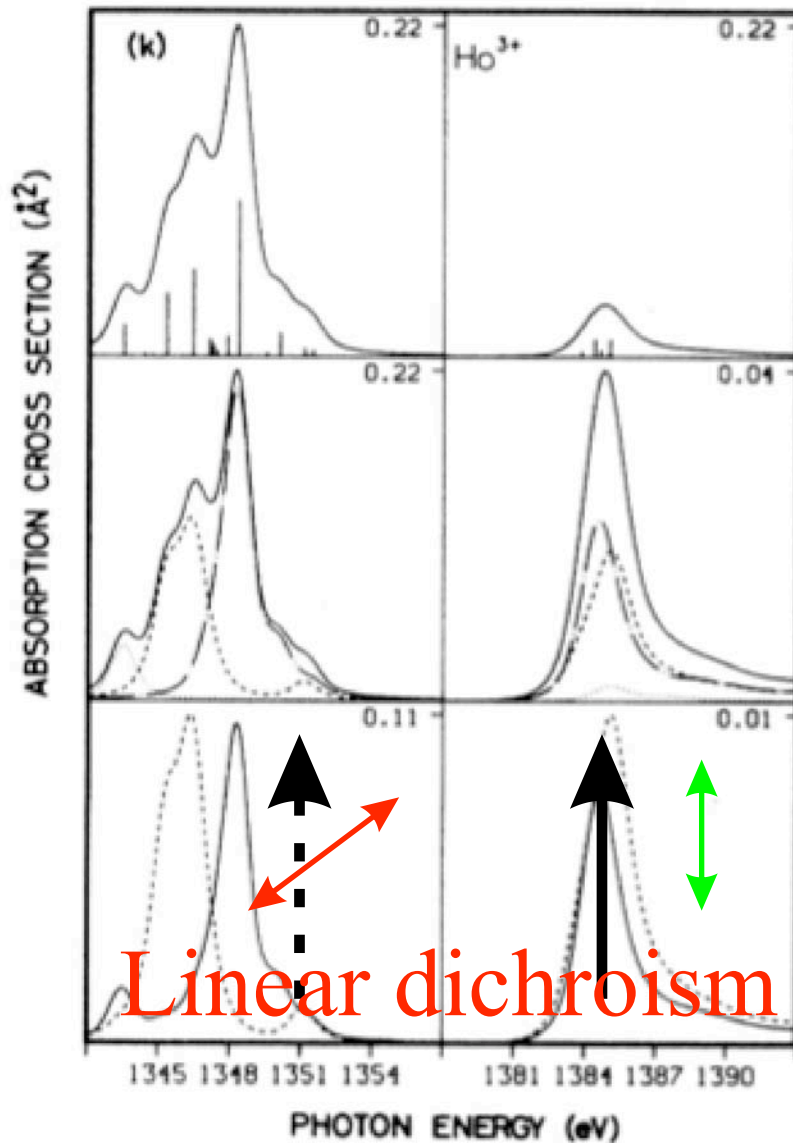
B. T. Thole, G. van der Laan, and G. A. Sawatzky

Department of Applied and Solid State Physics, Materials Science Center, University of Groningen, Nijenborgh 18, 9747 AG Groningen, The Netherlands

F. M. F. de Groot and J. C. Fuggle

Research Institute of Materials, Faculty of Science, University of Nijmegen, Toernooiveld, 6525 ED Nijmegen, The Netherlands

(Received 13 April 1987)



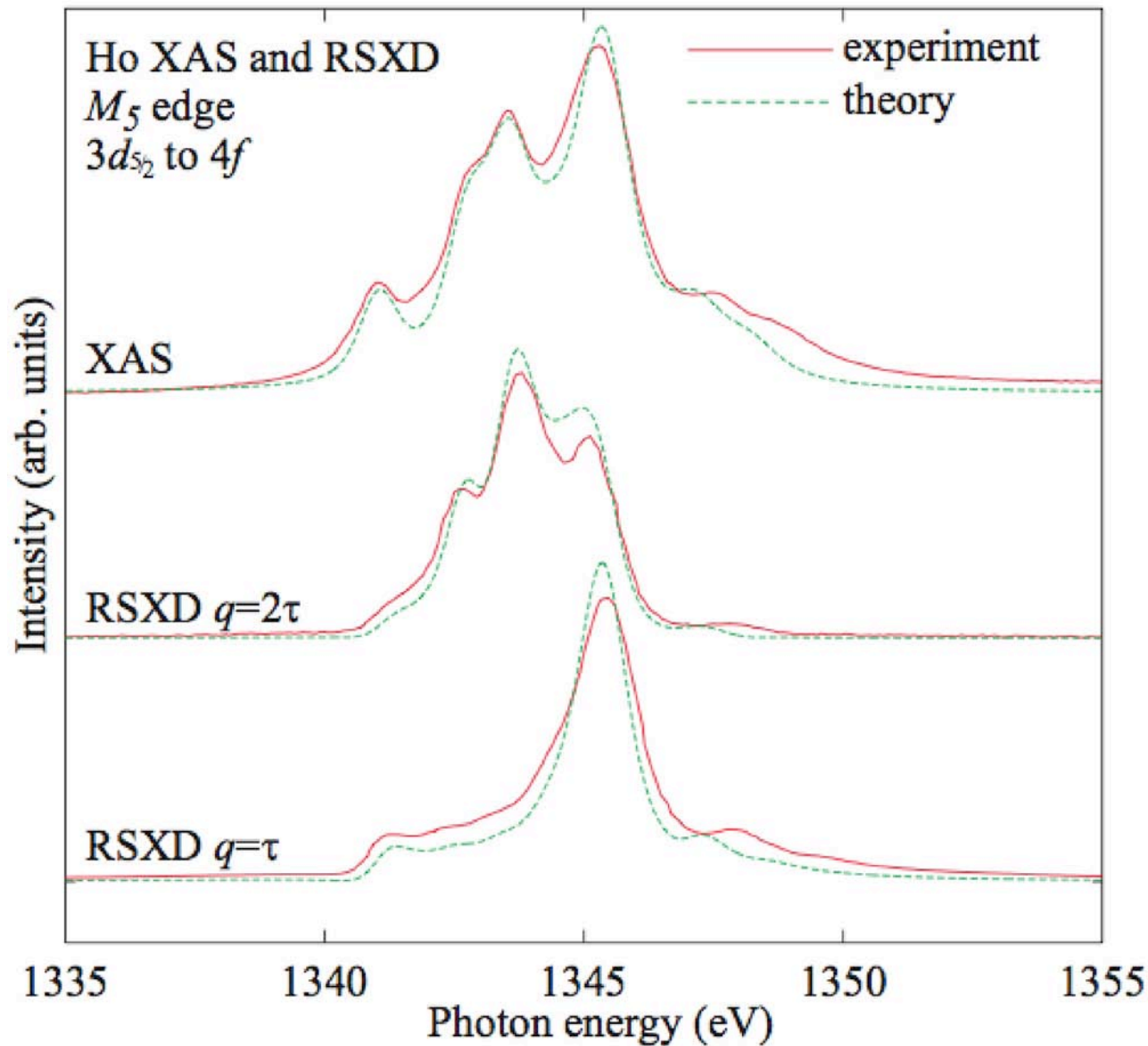
Plotted is $-\text{Im}[\sigma]$

top: $\sigma^{(0)}$

middle $\sigma^{(1)}$ (difference between the two spectra)

bottom $\sigma^{(2)}$ (difference between the two spectra)

Example on Ho: Elastic diffraction



top: $-\text{Im}[\sigma^{(0)}]$

middle $|\omega \sigma^{(2)}|^2$

bottom $|\omega \sigma^{(1)}|^2$

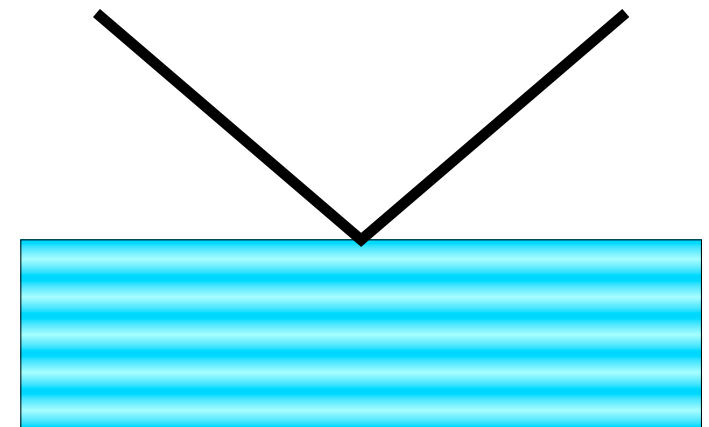
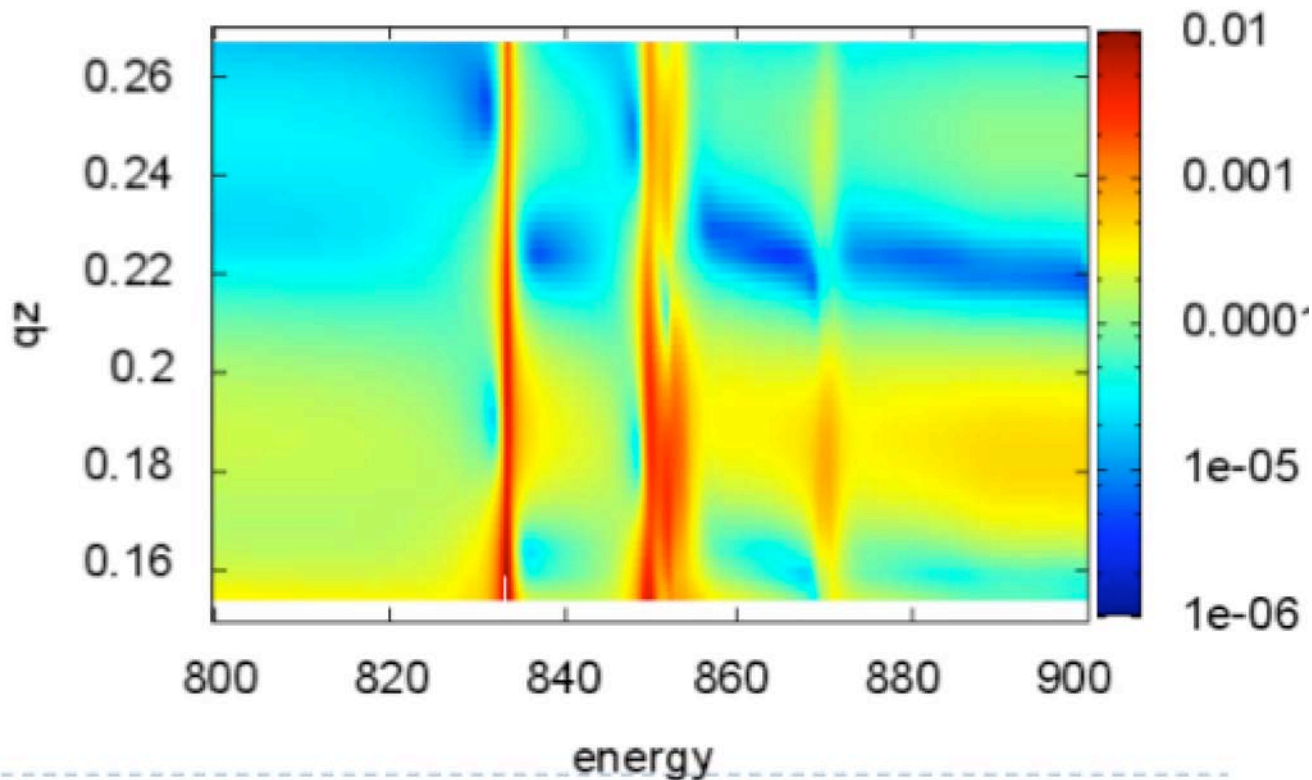
Can we use sum-rules to get quantitative information?

- ❑ Measure intensity for different q and polarizations
- ❑ Find F_q from the measured values of $|F_q|^2$
- ❑ Use that $F_i = \omega \sigma_i$

$$\int_{\mathbf{K}} \sigma = i \overline{\rho}_p$$

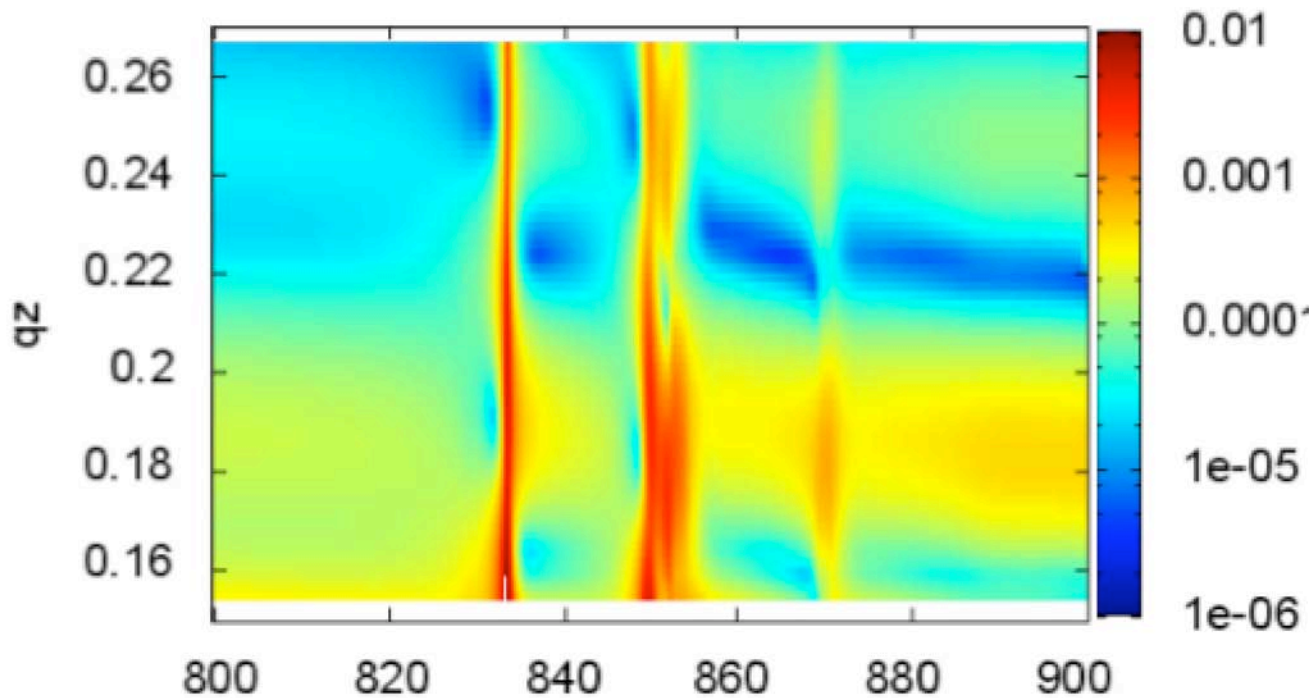
Can we use sum-rules to get quantitative information?

- Yes, but one needs to include dynamical scattering effects



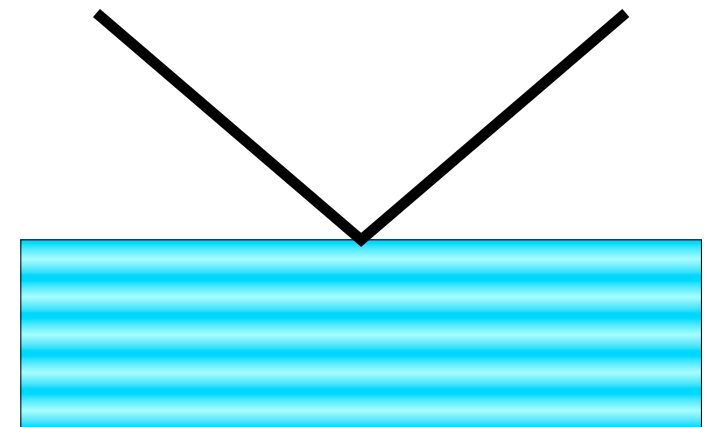
Can we use sum-rules to get quantitative information?

- Yes, but one needs to include dynamical scattering effects



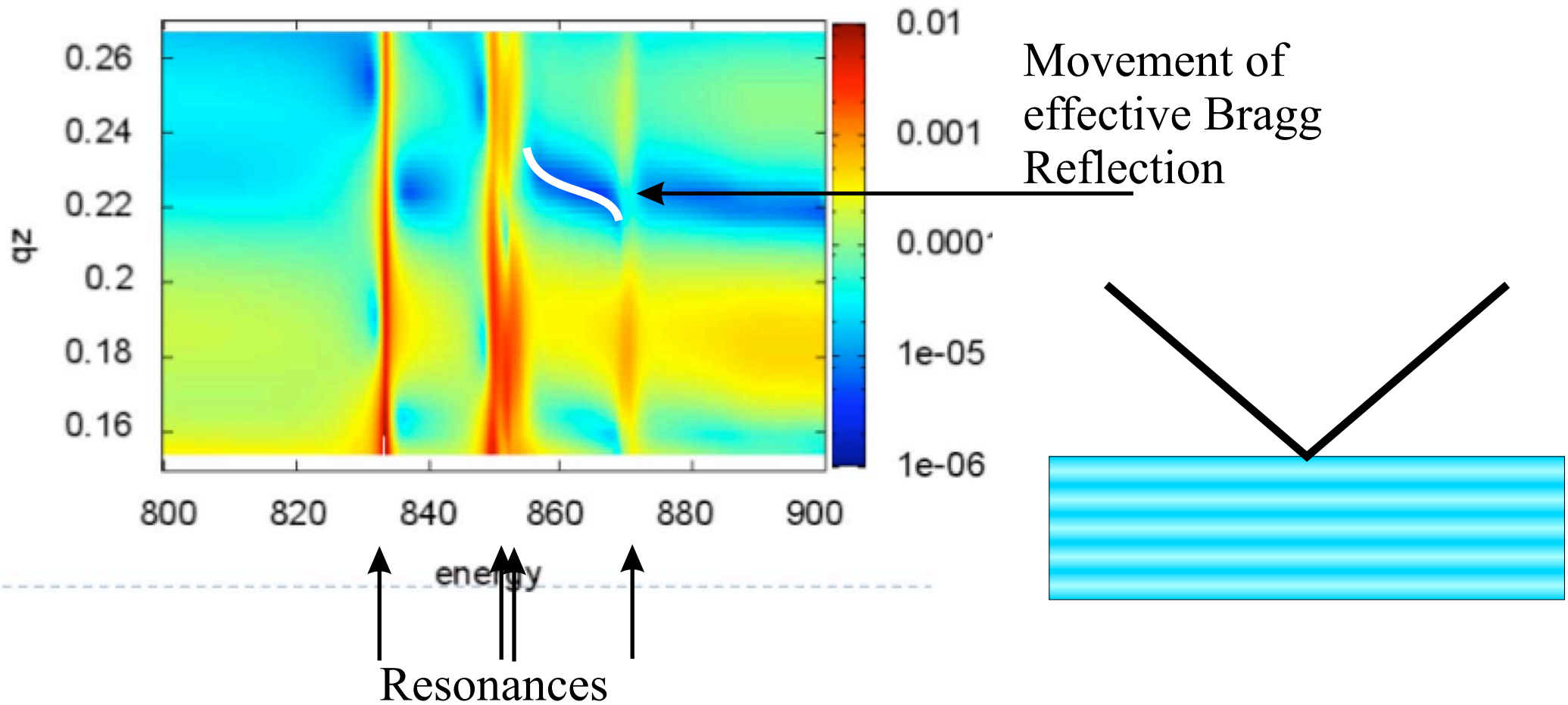
energy

Resonances



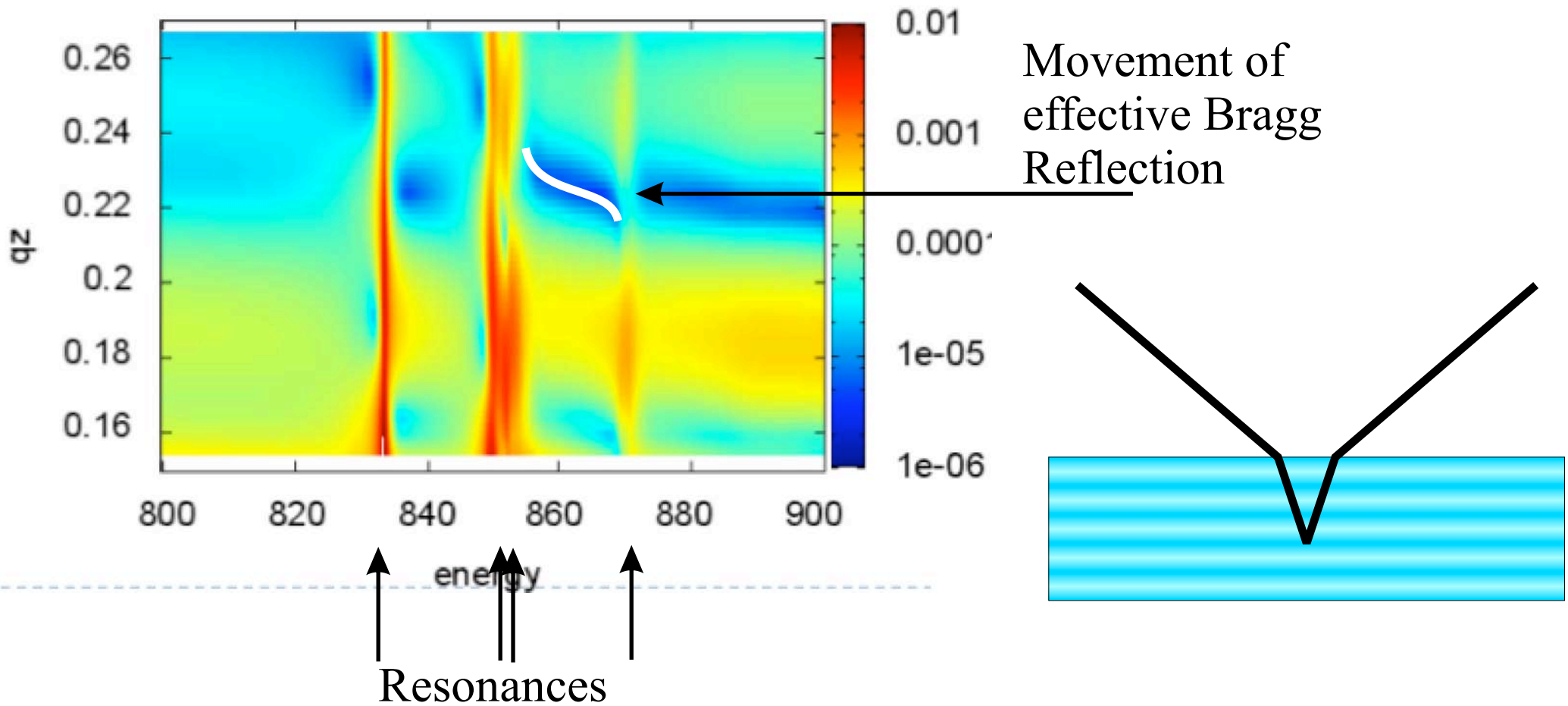
Can we use sum-rules to get quantitative information?

- Yes, but one needs to include dynamical scattering effects



Can we use sum-rules to get quantitative information?

- Yes, but one needs to include dynamical scattering effects



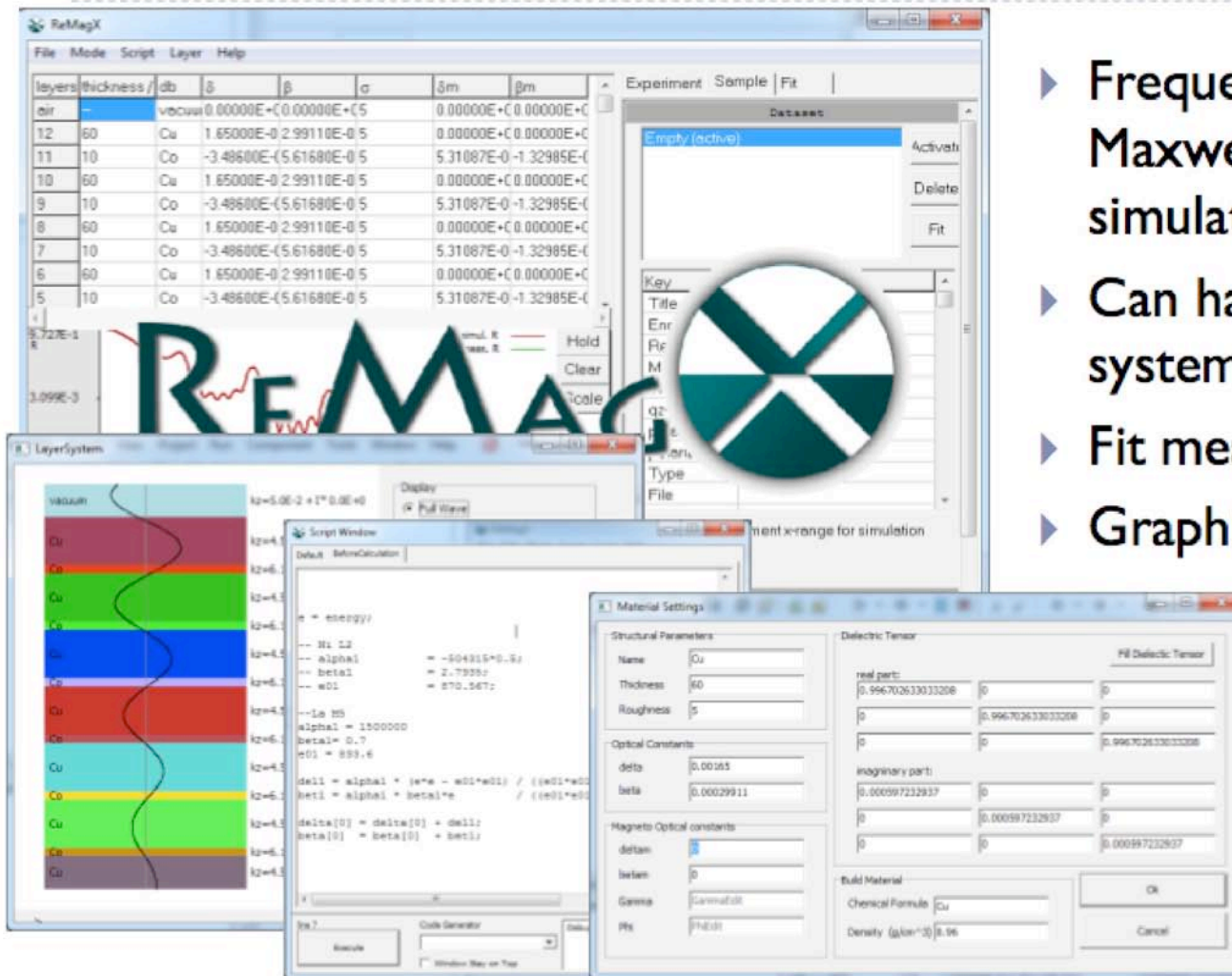
Can we use sum-rules to get quantitative information?

- ❑ Measure intensity for different q and polarizations
- ❑ Starting from $|F_q|^2$ reversely solve Maxwell equations to obtain F_i
- ❑ Use that $F_i = \omega \sigma_i$

$$\int_{\mathbf{K}} \sigma = i \overline{\rho}_p$$

Solve Maxwell equations (with side conditions)

Analysis tool **REMag** 



layers	thickness	db	δ	β	σ	δ_m	β_m
air			vacuum	$0.00000E+0$	$0.00000E+0$	$0.00000E+0$	$0.00000E+0$
12	60		Cu	$1.65000E-0$	$2.99110E-0$	5.0	$0.00000E+0$
11	10		Co	$-3.48600E-0$	$5.61680E-0$	5	$5.31087E-0$
10	60		Cu	$1.65000E-0$	$2.99110E-0$	5.0	$0.00000E+0$
9	10		Co	$-3.48600E-0$	$5.61680E-0$	5	$5.31087E-0$
8	60		Cu	$1.65000E-0$	$2.99110E-0$	5.0	$0.00000E+0$
7	10		Co	$-3.48600E-0$	$5.61680E-0$	5	$5.31087E-0$
6	60		Cu	$1.65000E-0$	$2.99110E-0$	5.0	$0.00000E+0$
5	10		Co	$-3.48600E-0$	$5.61680E-0$	5	$5.31087E-0$

```
--- beta[0] = -0.4815+0.0j
--- alpha[1] = 2.7993j
--- m01 = 0.70.547j
---lo m0
alpha[1] = 1500000
beta[1] = 0.7
m01 = 899.6
dell = alpha[1] + (m0 - m01*m01) / ((m01+m0)
bet1 = alpha[1] + beta[1]*e
delta[0] = delta[0] + dell;
beta[0] = beta[0] + bet1;
```

- ▶ Frequency-domain Maxwell-equation simulator
- ▶ Can handle dichroic systems
- ▶ Fit measurements
- ▶ Graphical user interface

Developed at MPI Stuttgart
and now at UBC
by S. Macke

Creating a High-Tc Cuprate like fermi-surface in Nickelates

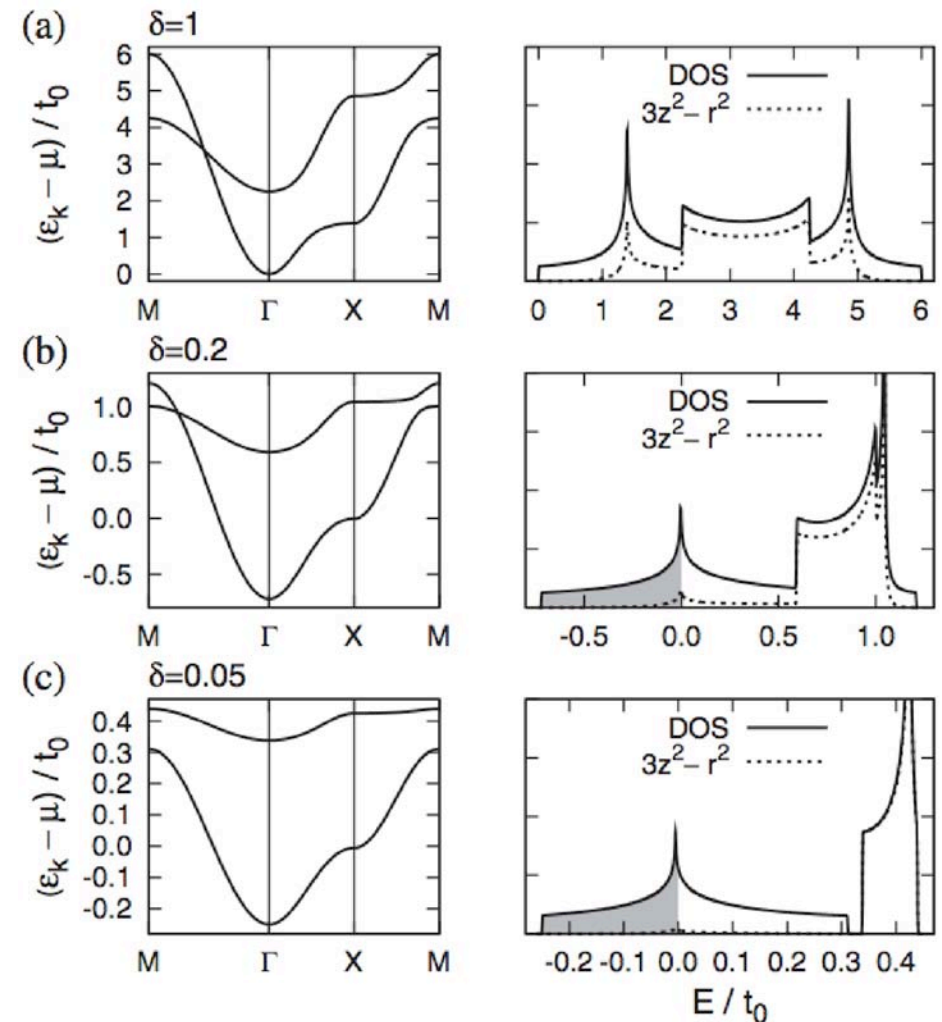
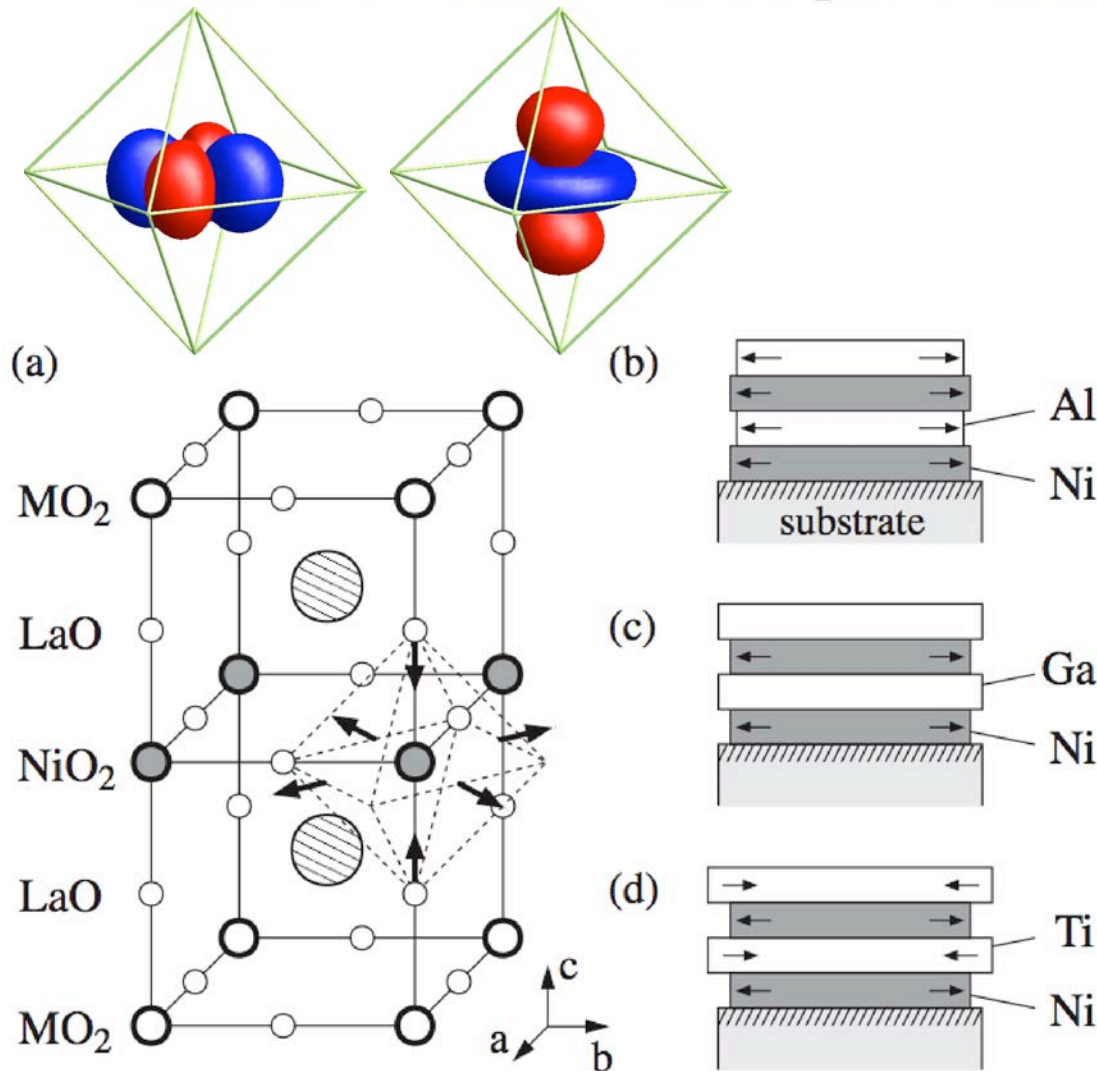
PRL **100**, 016404 (2008)

PHYSICAL REVIEW LETTERS

week ending
11 JANUARY 2008

Orbital Order and Possible Superconductivity in $\text{LaNiO}_3/\text{LaMO}_3$ Superlattices

Jiří Chaloupka^{1,2} and Giniyat Khaliullin¹



Creating a High-Tc Cuprate like fermi-surface in Nickelates

PRL 103, 016401 (2009)

PHYSICAL REVIEW LETTERS

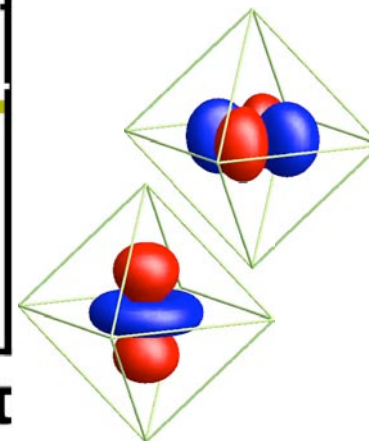
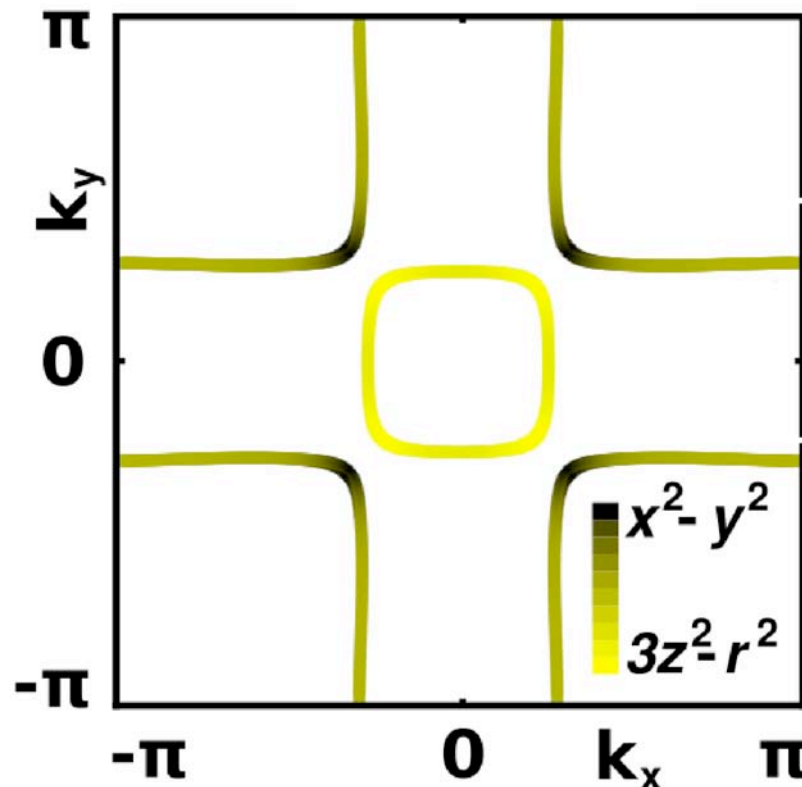
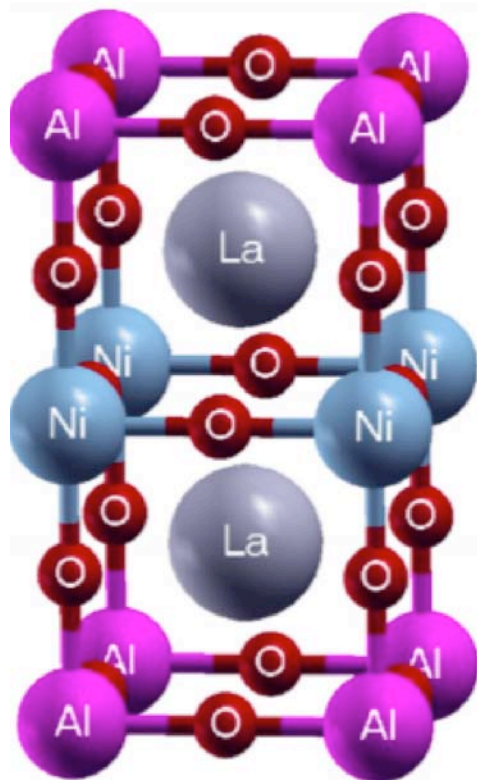
week ending
3 JULY 2009

Turning a Nickelate Fermi Surface into a Cupratelike One through Heterostructuring

P. Hansmann,^{1,2} Xiaoping Yang,¹ A. Toschi,^{1,2} G. Khaliullin,¹ O. K. Andersen,¹ and K. Held²

¹Max-Planck-Institut für Festkörperforschung, Heisenbergstrasse 1, D-70569 Stuttgart, Germany

²Institute for Solid State Physics, Vienna University of Technology, 1040 Vienna, Austria



Creating a High-Tc Cuprate like fermi-surface in Nickelates

PRL 103, 016401 (2009)

PHYSICAL REVIEW LETTERS

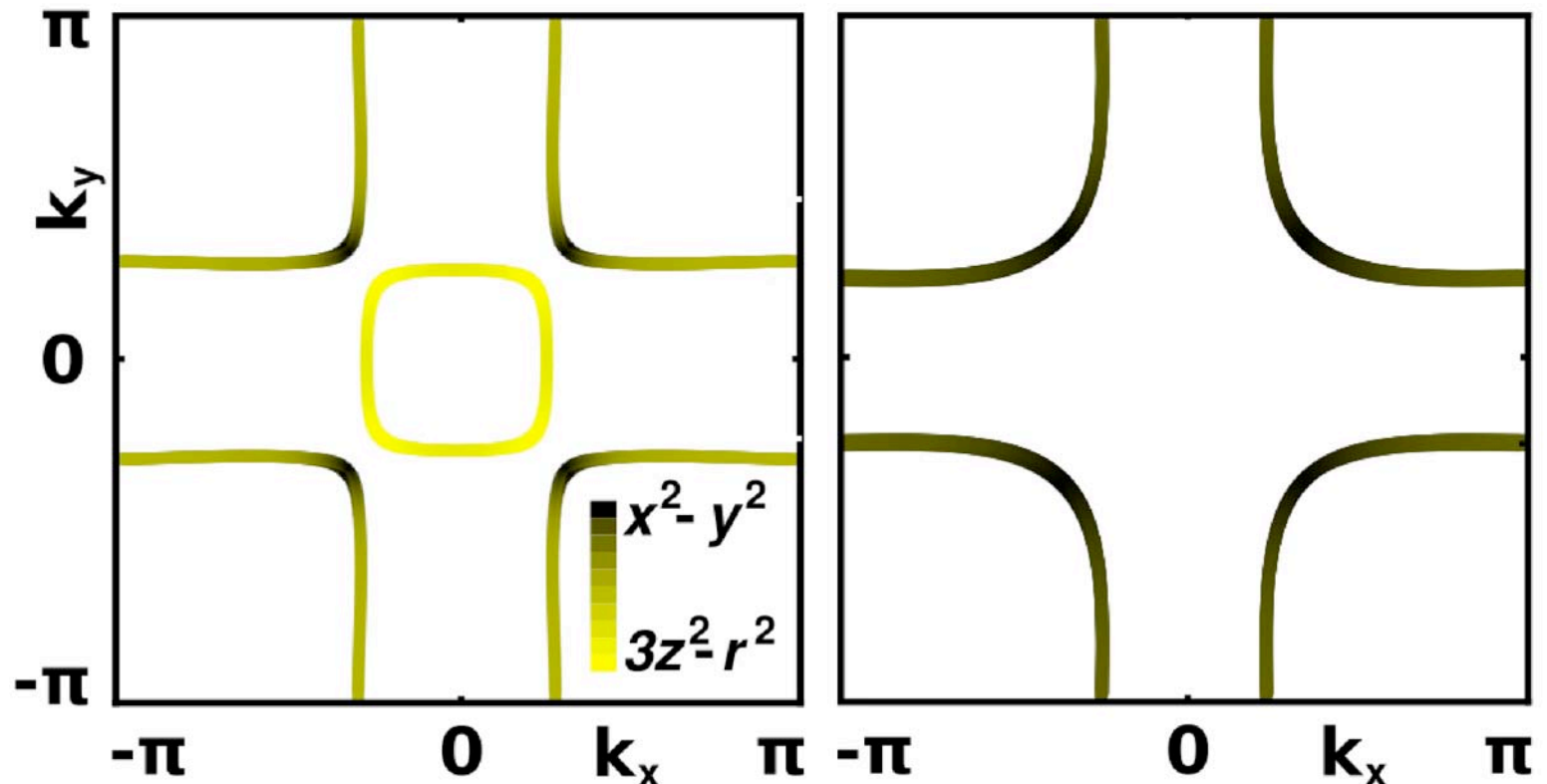
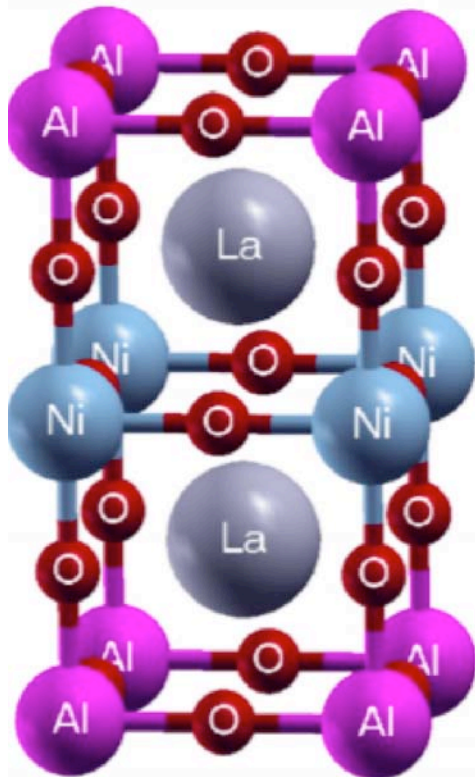
week ending
3 JULY 2009

Turning a Nickelate Fermi Surface into a Cupratelike One through Heterostructuring

P. Hansmann,^{1,2} Xiaoping Yang,¹ A. Toschi,^{1,2} G. Khaliullin,¹ O. K. Andersen,¹ and K. Held²

¹Max-Planck-Institut für Festkörperforschung, Heisenbergstrasse 1, D-70569 Stuttgart, Germany

²Institute for Solid State Physics, Vienna University of Technology, 1040 Vienna, Austria



Do we have an orbitally polarized Ni Film

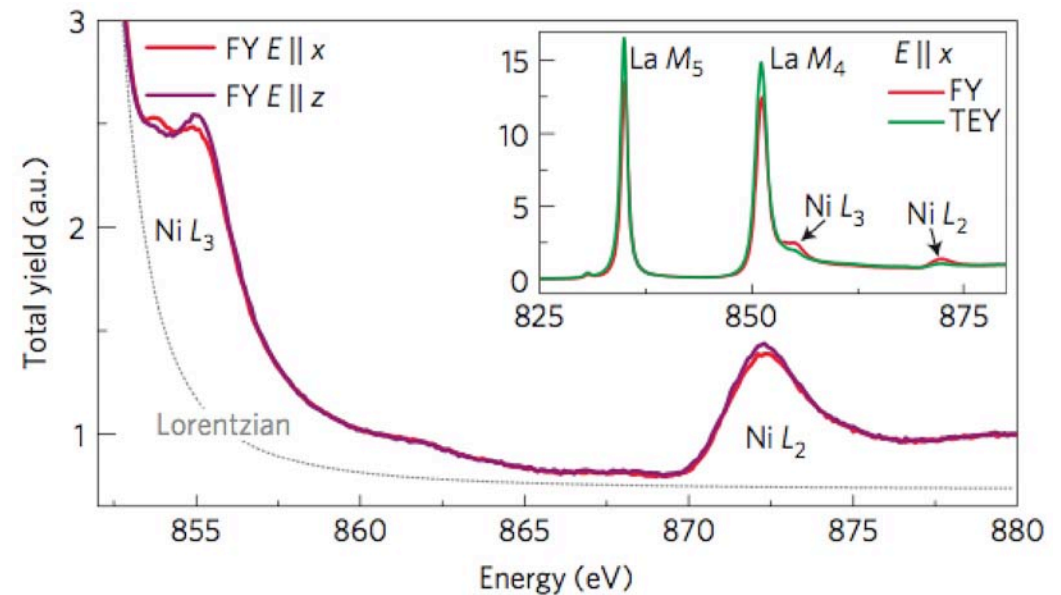
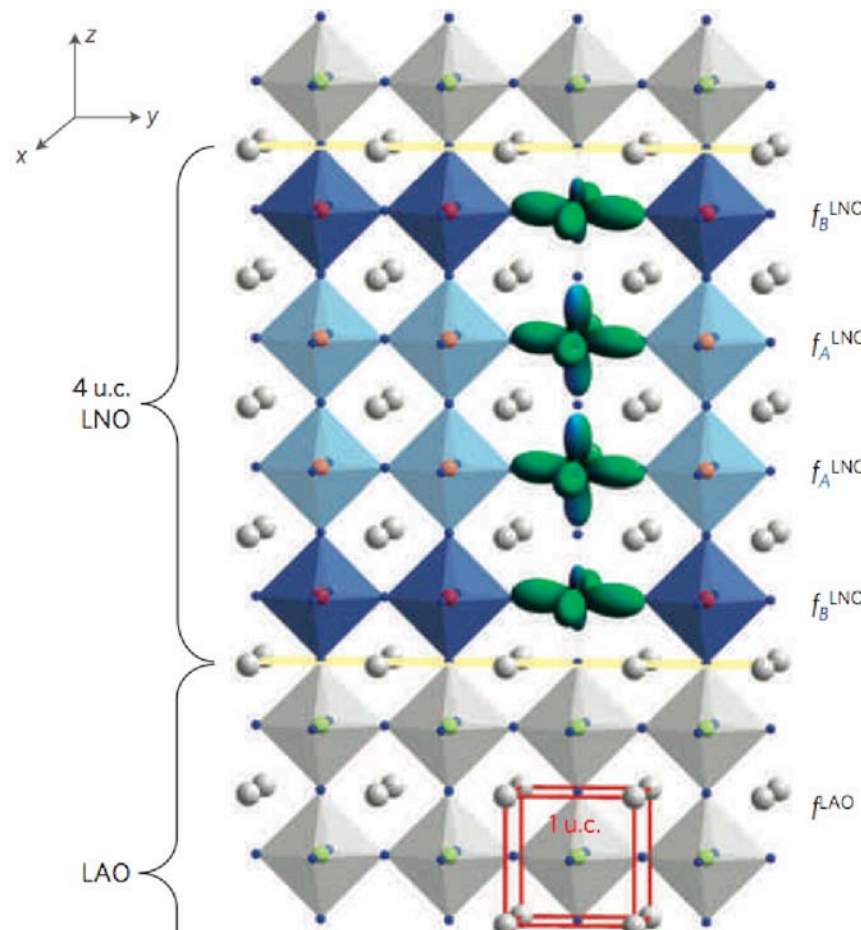
nature
materials

LETTERS

PUBLISHED ONLINE: 6 FEBRUARY 2011 | DOI: 10.1038/NMAT2958

Orbital reflectometry of oxide heterostructures

Eva Benckiser¹, Maurits W. Haverkort¹, Sebastian Brück^{2,3}, Eberhard Goering², Sebastian Macke², Alex Frañó¹, Xiaoping Yang^{1,4}, Ole K. Andersen¹, Georg Cristiani¹, Hanns-Ulrich Habermeier¹, Alexander V. Boris¹, Ioannis Zegkinoglou¹, Peter Wochner², Heon-Jung Kim^{1,5}, Vladimir Hinkov^{1*} and Bernhard Keimer^{1*}



Do we have an orbitally polarized Ni Film

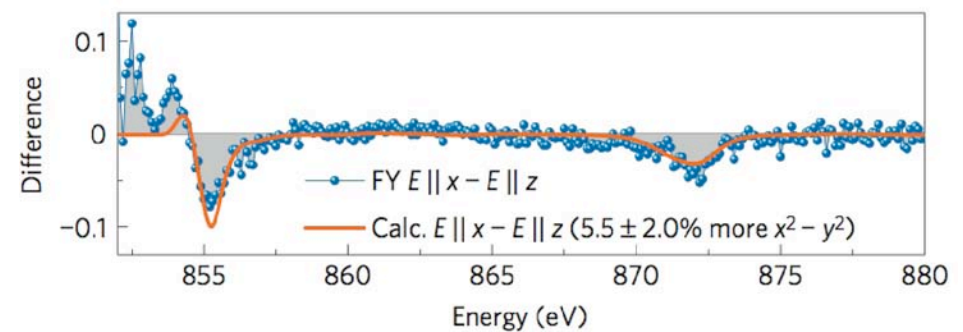
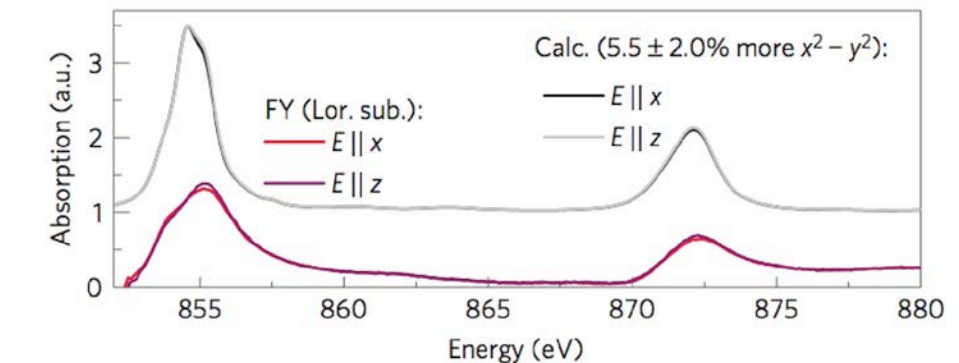
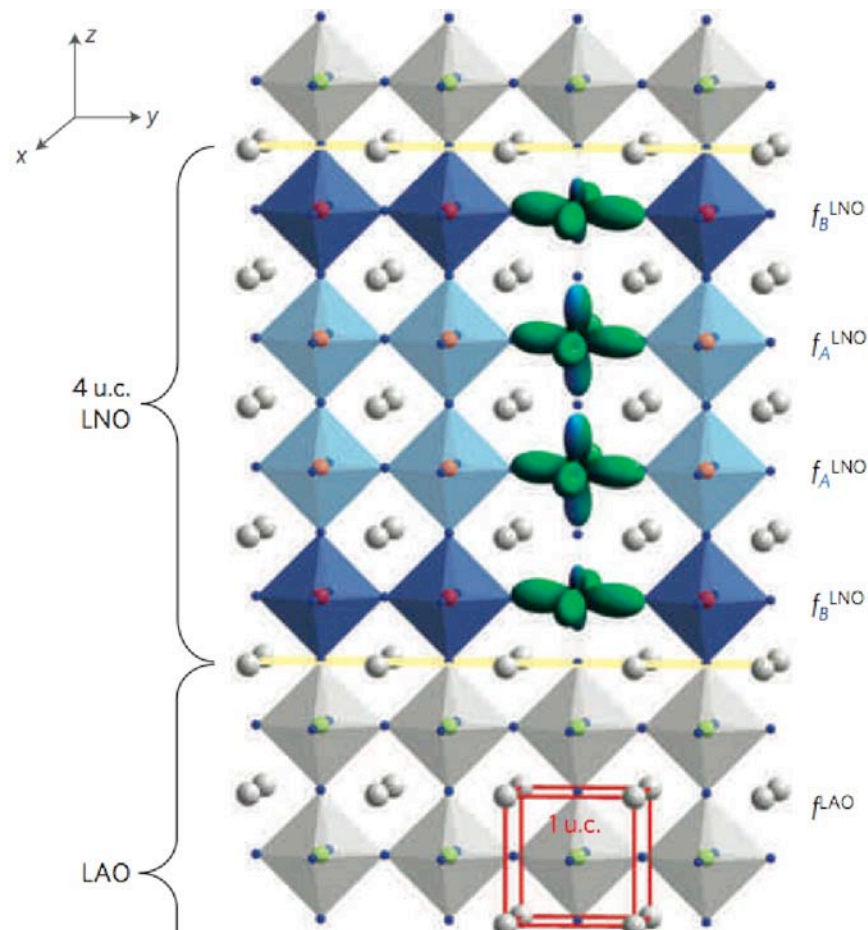
nature
materials

LETTERS

PUBLISHED ONLINE: 6 FEBRUARY 2011 | DOI: 10.1038/NMAT2958

Orbital reflectometry of oxide heterostructures

Eva Benckiser¹, Maurits W. Haverkort¹, Sebastian Brück^{2,3}, Eberhard Goering², Sebastian Macke², Alex Frañó¹, Xiaoping Yang^{1,4}, Ole K. Andersen¹, Georg Cristiani¹, Hanns-Ulrich Habermeier¹, Alexander V. Boris¹, Ioannis Zegkinoglou¹, Peter Wochner², Heon-Jung Kim^{1,5}, Vladimir Hinkov^{1*} and Bernhard Keimer^{1*}



Do we have an orbitally polarized Ni Film

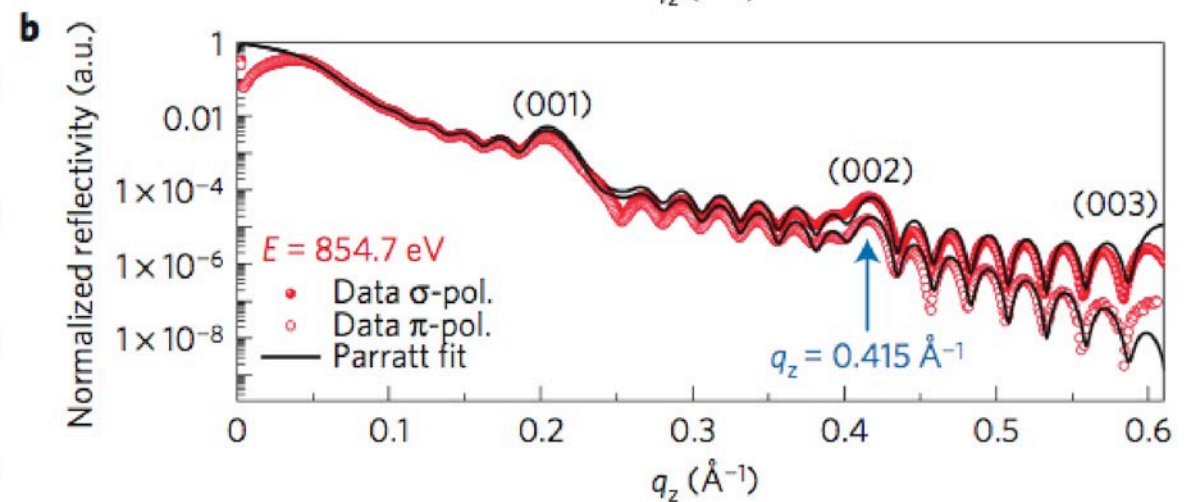
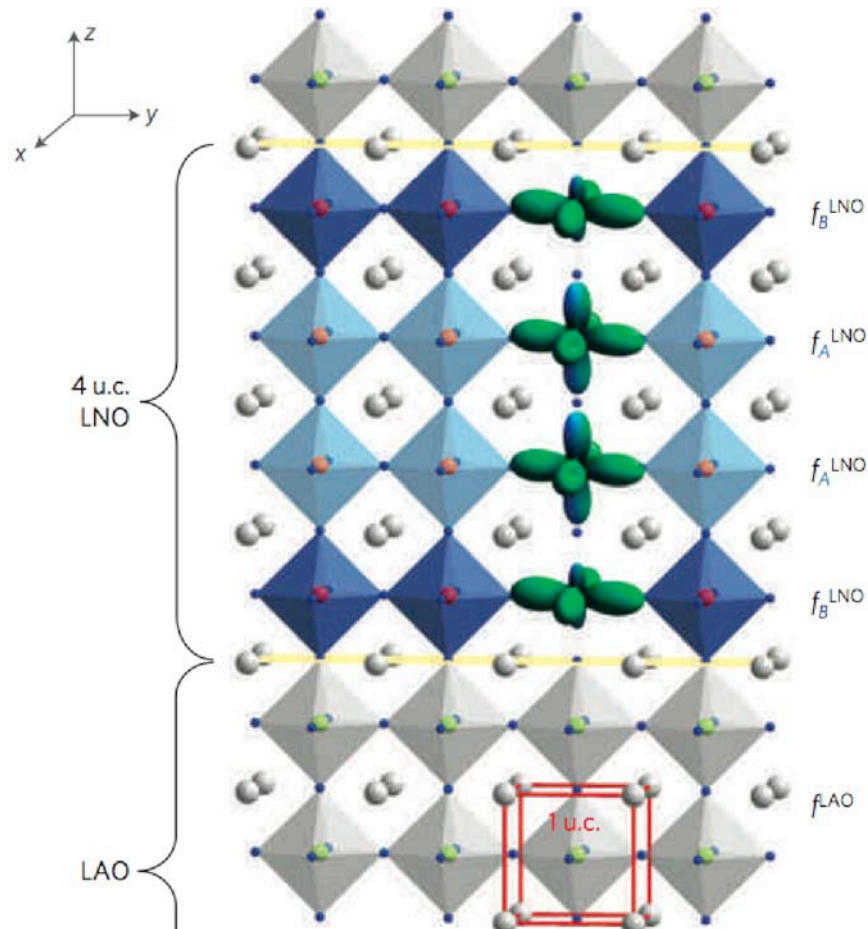
nature
materials

LETTERS

PUBLISHED ONLINE: 6 FEBRUARY 2011 | DOI: 10.1038/NMAT2958

Orbital reflectometry of oxide heterostructures

Eva Benckiser¹, Maurits W. Haverkort¹, Sebastian Brück^{2,3}, Eberhard Goering², Sebastian Macke², Alex Frañó¹, Xiaoping Yang^{1,4}, Ole K. Andersen¹, Georg Cristiani¹, Hanns-Ulrich Habermeier¹, Alexander V. Boris¹, Ioannis Zegkinoglou¹, Peter Wochner², Heon-Jung Kim^{1,5}, Vladimir Hinkov^{1*} and Bernhard Keimer^{1*}



Do we have an orbitally polarized Ni Film

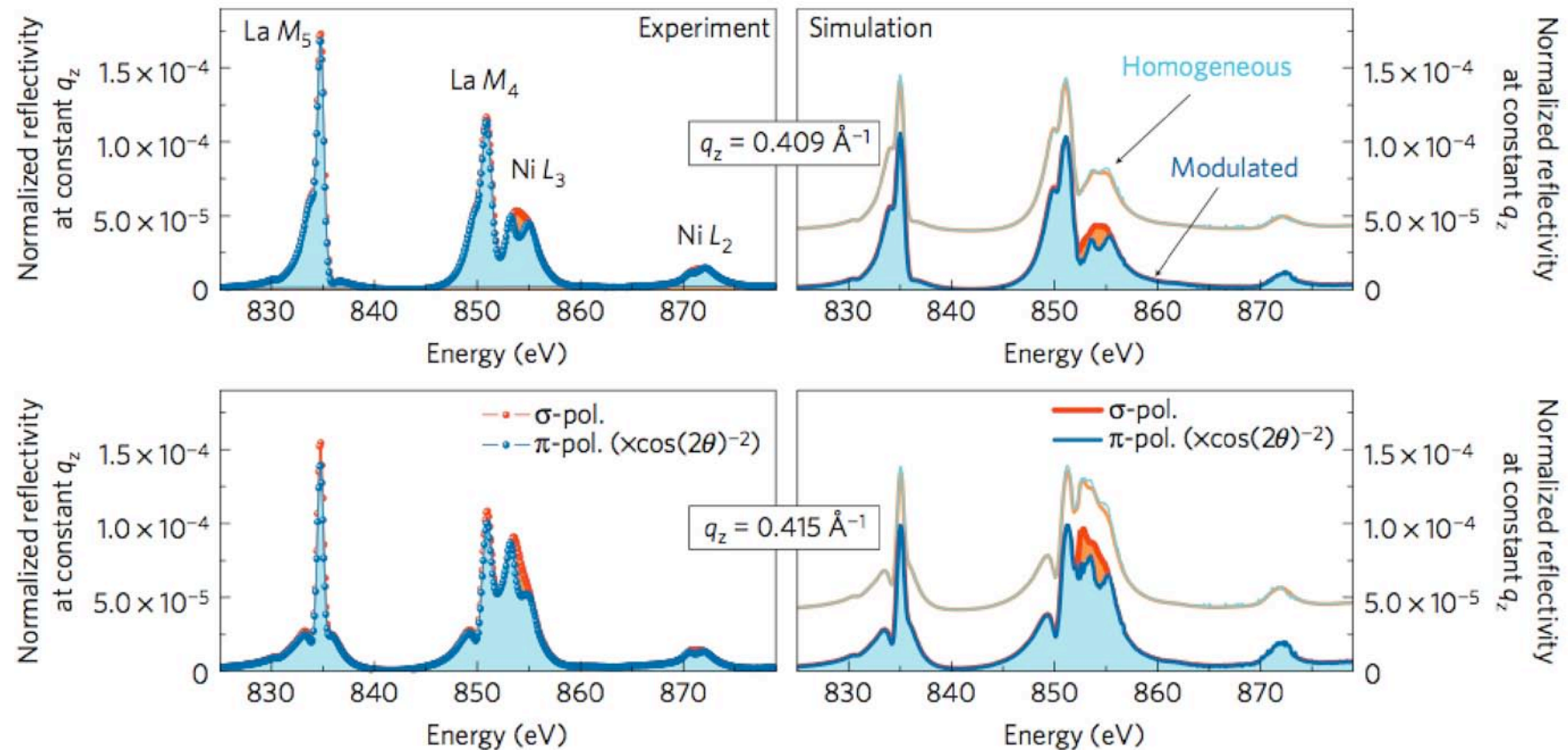
nature
materials

LETTERS

PUBLISHED ONLINE: 6 FEBRUARY 2011 | DOI: 10.1038/NMAT2958

Orbital reflectometry of oxide heterostructures

Eva Benckiser¹, Maurits W. Haverkort¹, Sebastian Brück^{2,3}, Eberhard Goering², Sebastian Macke², Alex Frañó¹, Xiaoping Yang^{1,4}, Ole K. Andersen¹, Georg Cristiani¹, Hanns-Ulrich Habermeier¹, Alexander V. Boris¹, Ioannis Zegkinoglou¹, Peter Wochner², Heon-Jung Kim^{1,5}, Vladimir Hinkov^{1*} and Bernhard Keimer^{1*}



Do we have an orbitally polarized Ni Film

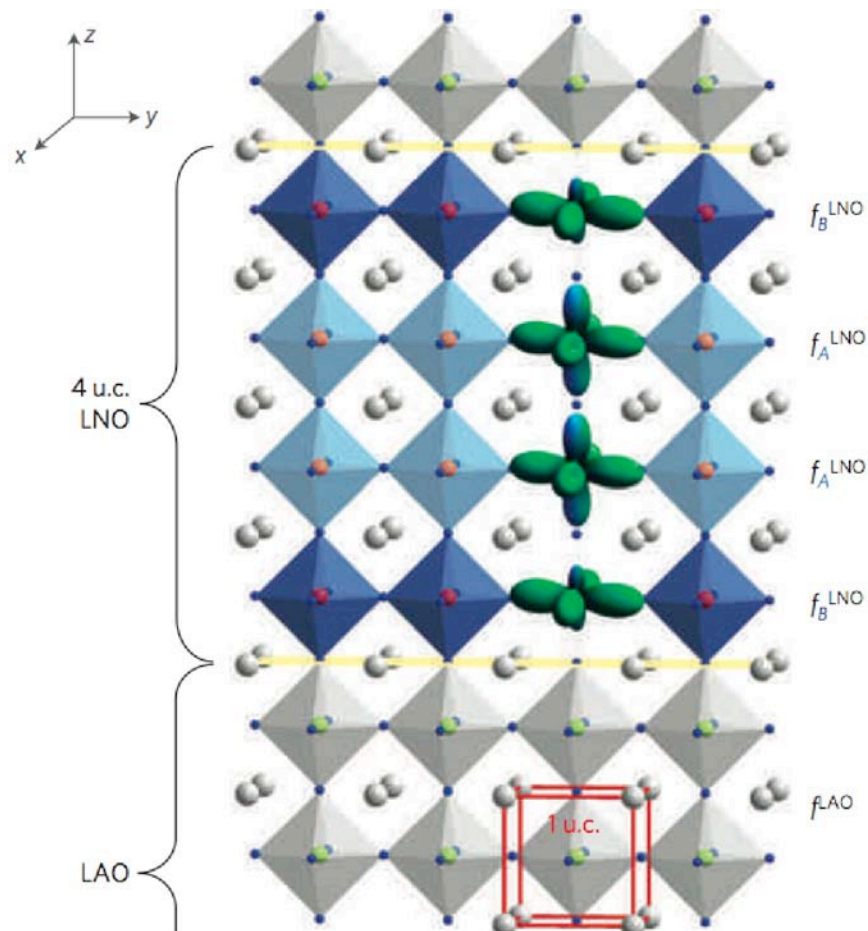
nature
materials

LETTERS

PUBLISHED ONLINE: 6 FEBRUARY 2011 | DOI: 10.1038/NMAT2958

Orbital reflectometry of oxide heterostructures

Eva Benckiser¹, Maurits W. Haverkort¹, Sebastian Brück^{2,3}, Eberhard Goering², Sebastian Macke², Alex Frañó¹, Xiaoping Yang^{1,4}, Ole K. Andersen¹, Georg Cristiani¹, Hanns-Ulrich Habermeier¹, Alexander V. Boris¹, Ioannis Zegkinoglou¹, Peter Wochner², Heon-Jung Kim^{1,5}, Vladimir Hinkov^{1*} and Bernhard Keimer^{1*}



Do we have an orbitally polarized Ni Film

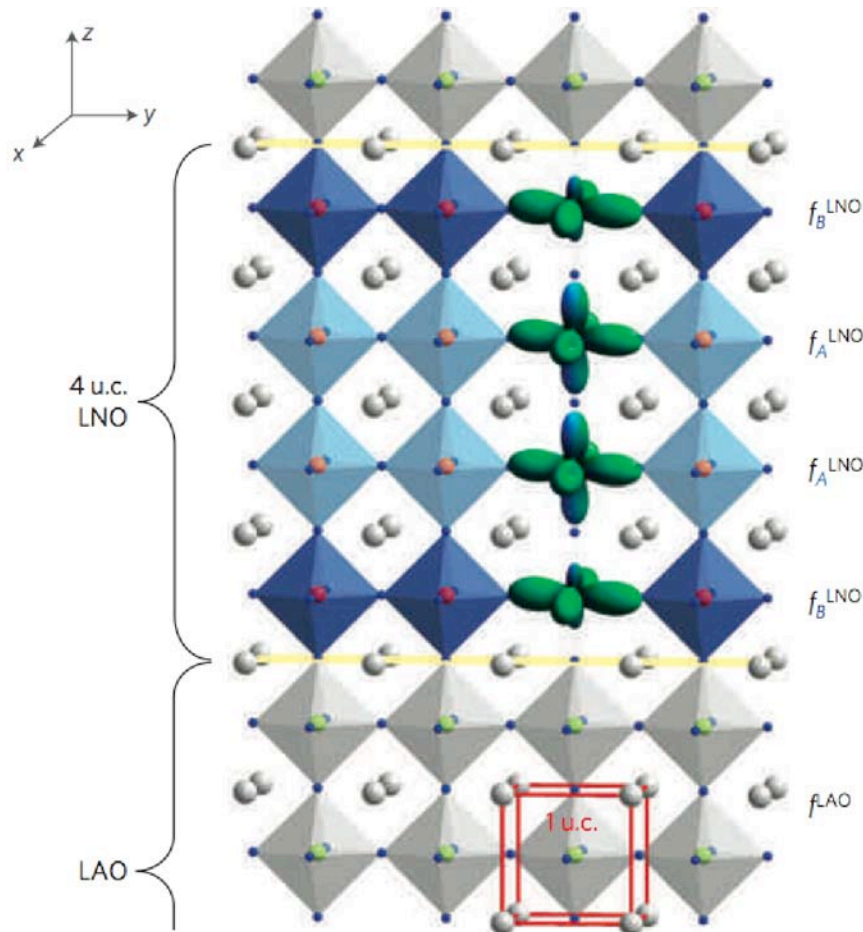
nature
materials

LETTERS

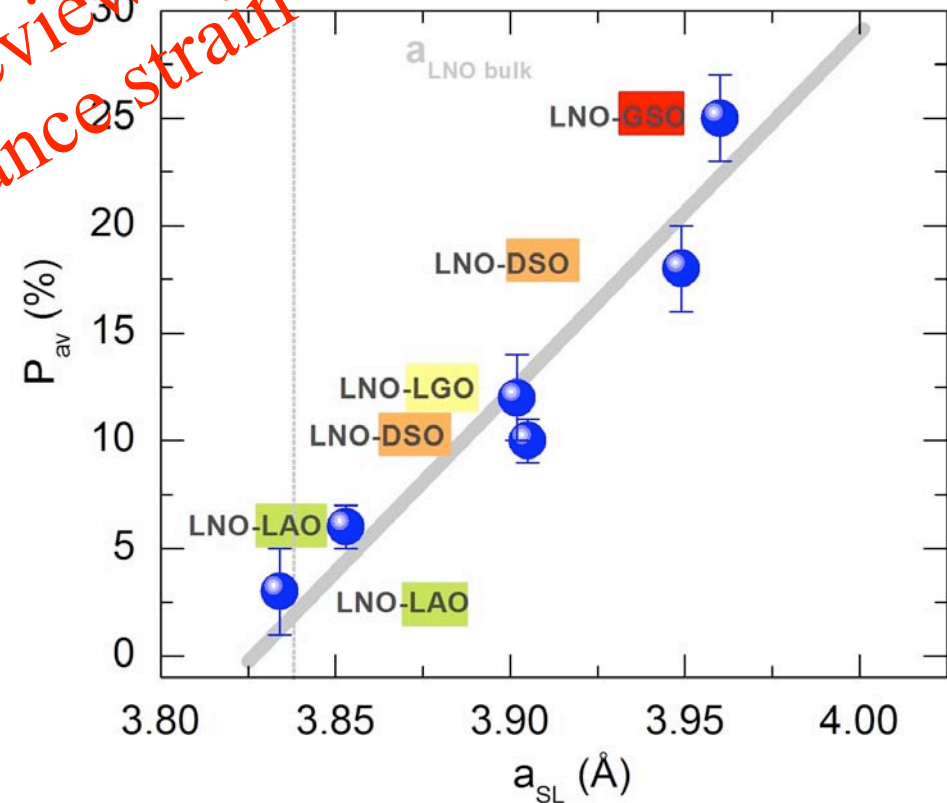
PUBLISHED ONLINE: 6 FEBRUARY 2011 | DOI: 10.1038/NMAT2958

Orbital reflectometry of oxide heterostructures

Eva Benckiser¹, Maurits W. Haverkort¹, Sebastian Brück^{2,3}, Eberhard Goering², Sebastian Macke², Alex Frañó¹, Xiaoping Yang^{1,4}, Ole K. Andersen¹, Georg Cristiani¹, Hanns-Ulrich Habermeier¹, Alexander V. Boris¹, Ioannis Zegkinoglou¹, Peter Wochner², Heon-Jung Kim^{1,5}, Vladimir Hinkov^{1*} and Bernhard Keimer^{1*}

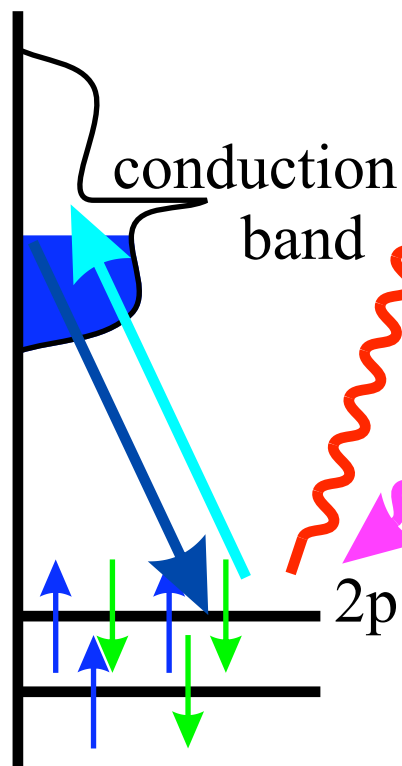
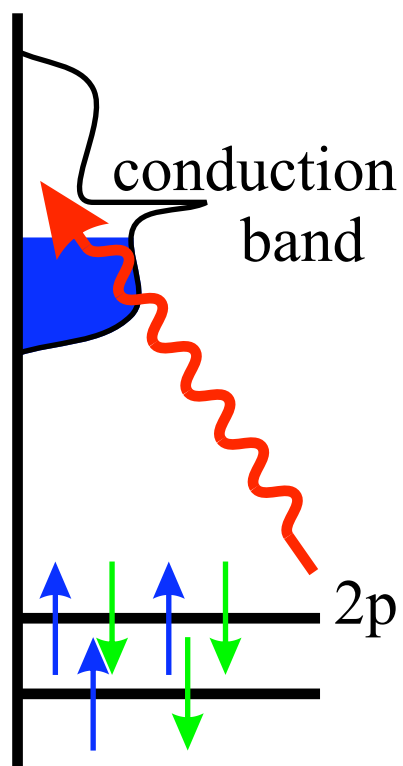


preview
Enhance strain

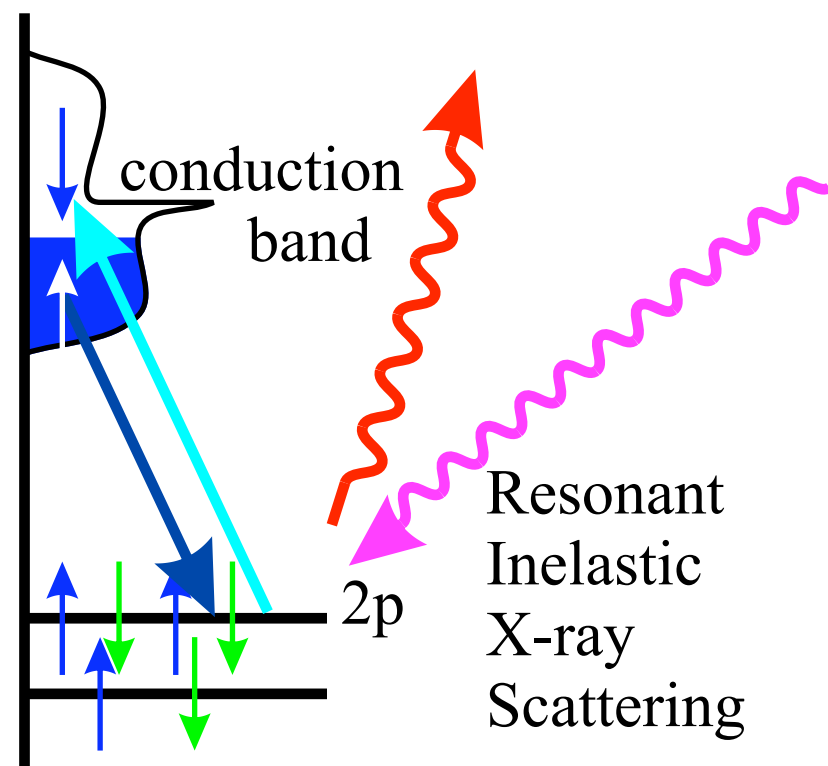


Theory of XAS, RXD, and RIXS

X-ray
Absorption
Spectroscopy



Resonant
X-ray
Diffraction



Resonant
Inelastic
X-ray
Scattering

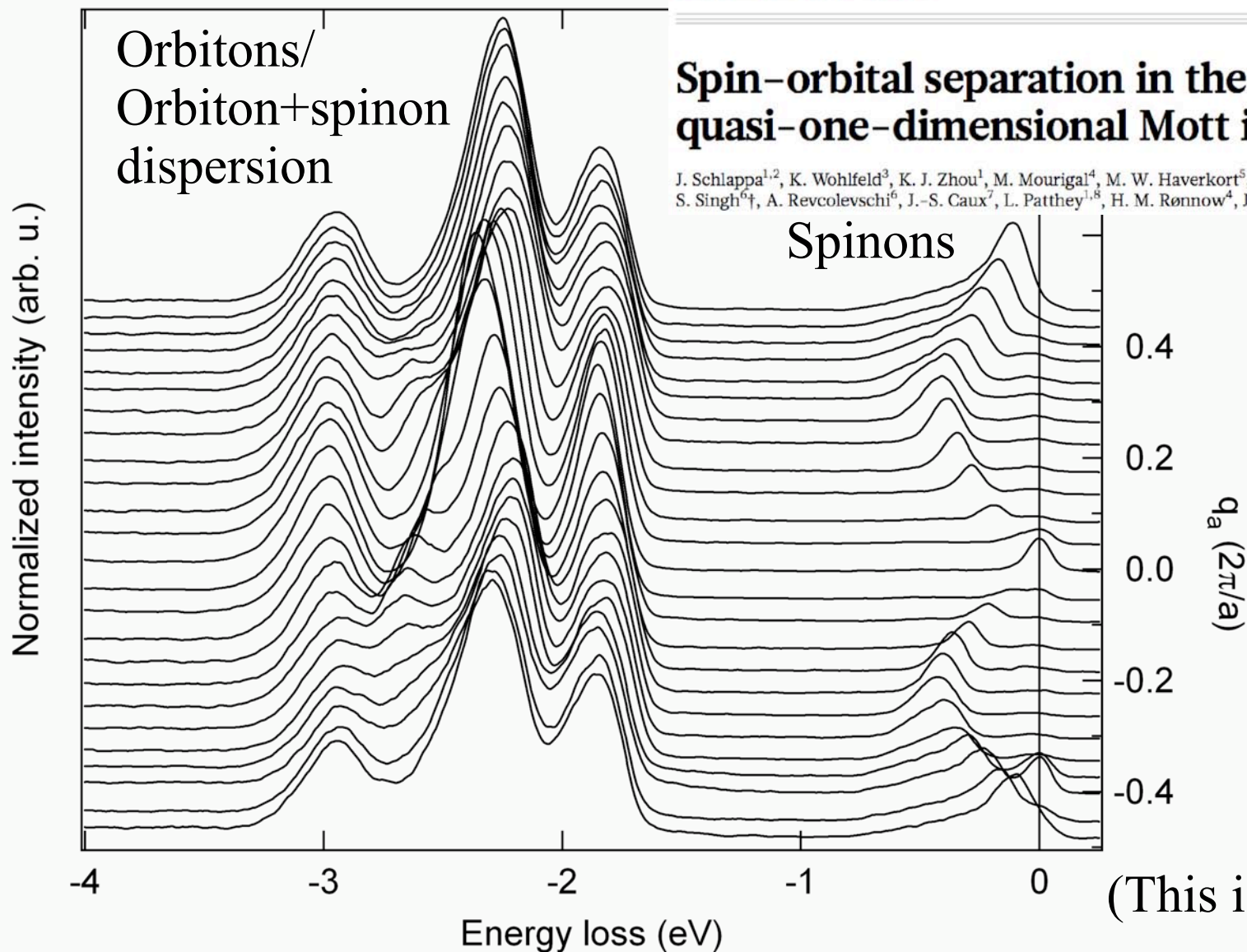
RIXS: Spin and Orbital (dispersing) excitations

LETTER

doi:10.1038/nature10974

Spin-orbital separation in the quasi-one-dimensional Mott insulator Sr_2CuO_3

J. Schlappa^{1,2}, K. Wohlfeld³, K. J. Zhou¹, M. Mourigal⁴, M. W. Haverkort⁵, V. N. Strocov¹, L. Hozoi³, C. Monney¹, S. Nishimoto³, S. Singh^{6†}, A. Revcolevschii⁶, J.-S. Caux⁷, L. Patthey^{1,8}, H. M. Rønnow⁴, J. van den Brink³ & T. Schmitt¹

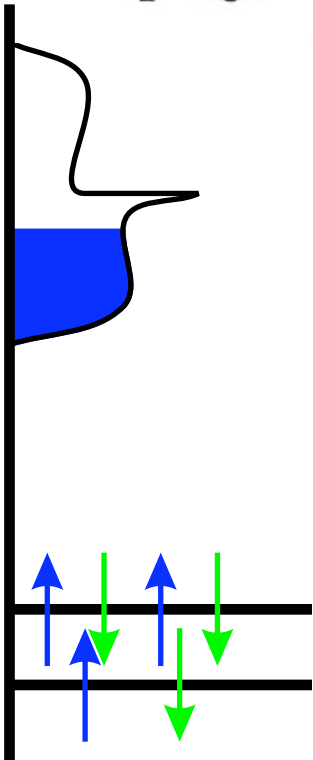


(This is raw data!!)

Spin excitations with resonant x-rays: Theory

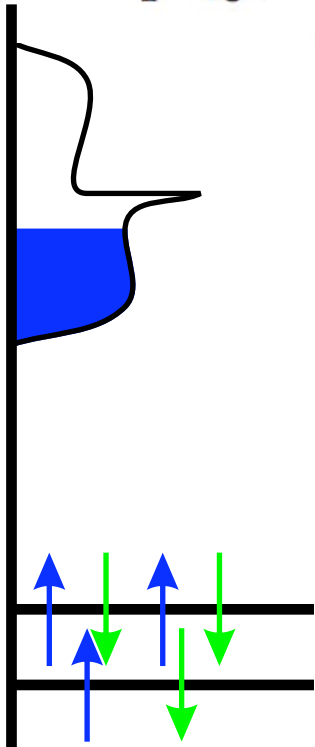
- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\epsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$



Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\epsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$


The diagram shows a vertical black line representing a crystal surface. To the right of the surface, a blue shaded region indicates a spin excitation. Below the surface, a lattice of atoms is shown with blue and green arrows representing spin orientations. A large cyan bracket connects the diagram to the equation above, highlighting the transition matrix elements and the energy levels involved in the scattering process.

Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\epsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$

The diagram illustrates the theory of spin excitations with resonant x-rays. It shows two energy level diagrams for two different states, with a red wavy arrow representing an x-ray photon and a red arrow pointing from the equation to the diagram.

The diagram shows two energy level diagrams for two different states, with a red wavy arrow representing an x-ray photon and a red arrow pointing from the equation to the diagram. The energy levels are represented by horizontal lines, and the spin states are indicated by blue and green arrows. The red wavy arrow represents the x-ray photon, and the red arrow points from the equation to the diagram.

Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\epsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$

Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\varepsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\varepsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$

Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\varepsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\varepsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$

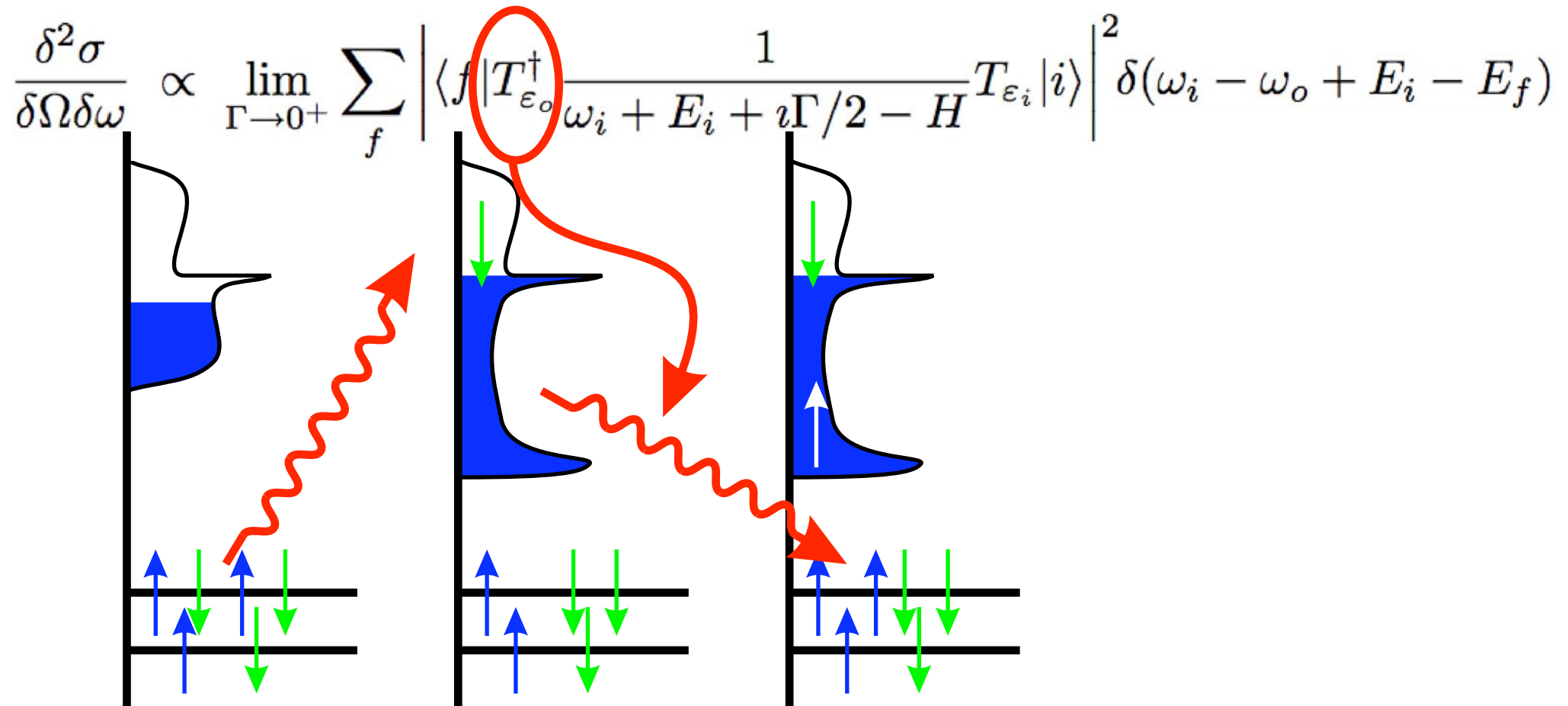
Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\varepsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\varepsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$

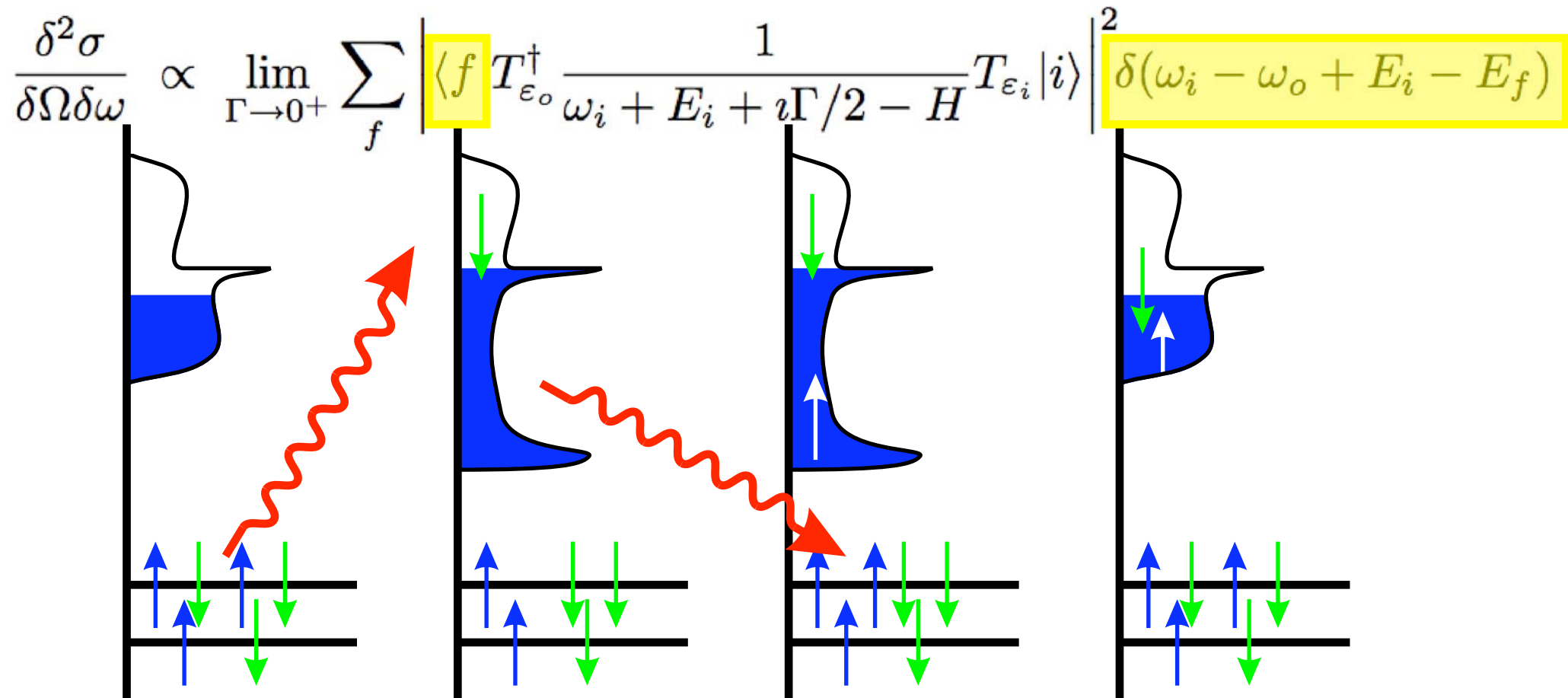
Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...



Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...



Spin excitations with resonant x-rays: Theory


- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...


$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\epsilon_i}^\dagger \frac{R_{\omega_i} \epsilon_i \cdot \epsilon_o}{H} T_{\epsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_o + E_i - E_f)$$

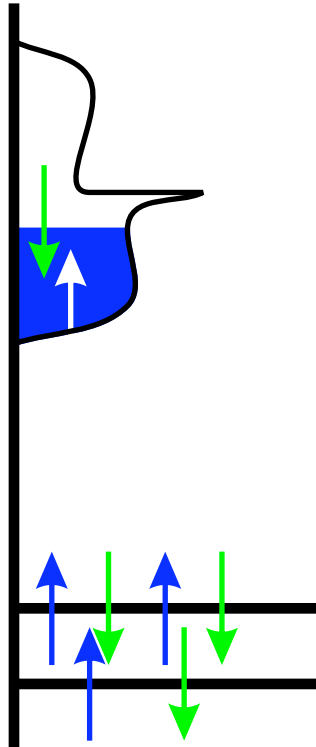
Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \lim_{\Gamma \rightarrow 0^+} \sum_f \left| \langle f | T_{\epsilon_0}^\dagger \frac{R_{\omega_i}^{\epsilon_i \epsilon_0}}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i} | i \rangle \right|^2 \delta(\omega_i - \omega_0 + E_i - E_f)$$







$$R_{\omega_i}^{\epsilon_i \epsilon_0} = T_{\epsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i}$$

Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto -\text{Im} \lim_{\Gamma \rightarrow 0^+} \langle i | (R_{\omega_i, Q}^{\epsilon_i \epsilon_o})^\dagger \frac{1}{\omega + E_i + i\Gamma/2 - H} R_{\omega_i, Q}^{\epsilon_i \epsilon_o} | i \rangle$$

$$R_{\omega_i}^{\epsilon_i \epsilon_o} = T_{\epsilon_o}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i}$$

Spin excitations with resonant x-rays: Theory

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\frac{\delta^2 \sigma}{\delta \Omega \delta \omega} \propto \text{Re} \int_0^\infty e^{i\omega t} \langle i | R_{\omega_i, Q}^{\varepsilon_i \varepsilon_o}(t)^\dagger R_{\omega_i, Q}^{\varepsilon_i \varepsilon_o}(t=0) | i \rangle dt$$

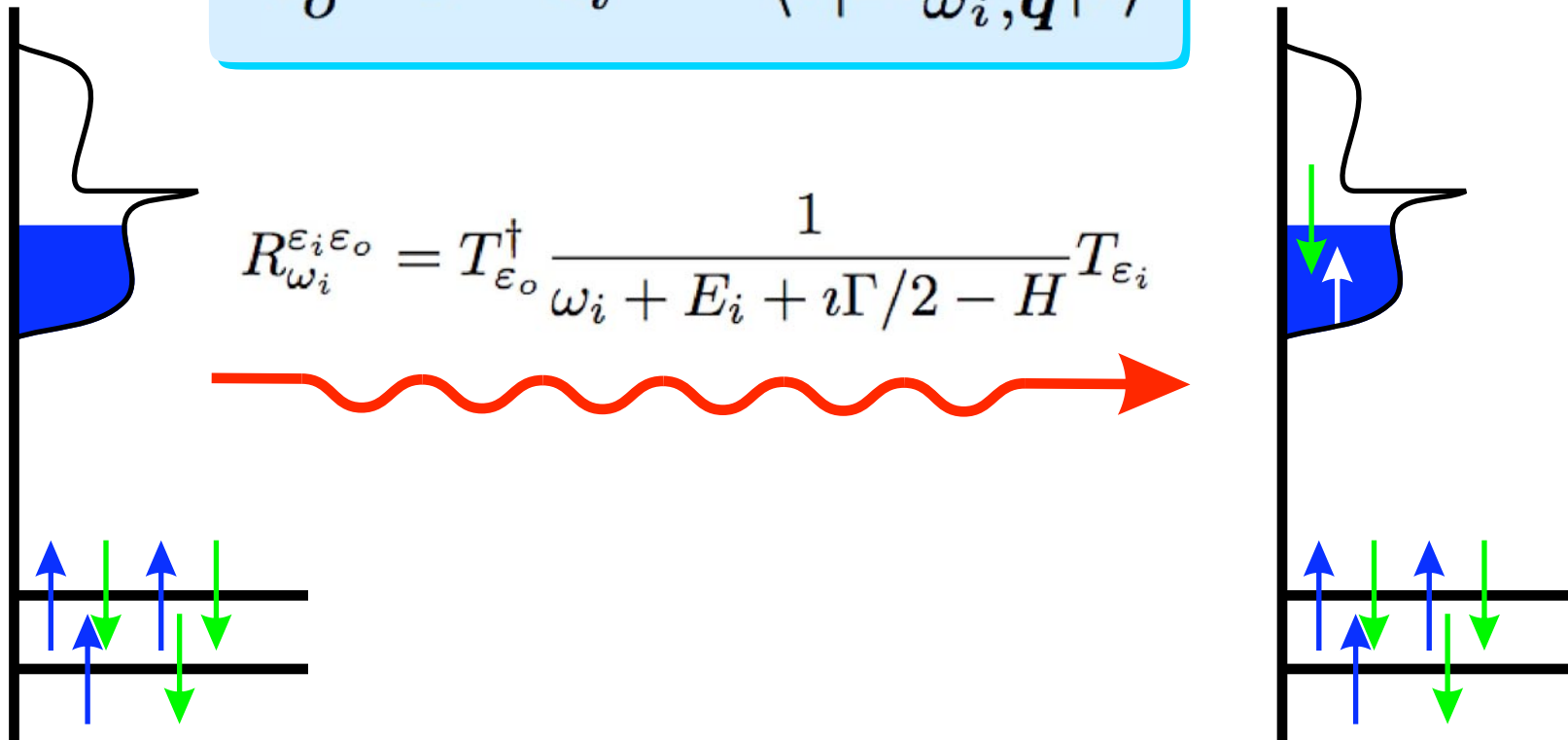
$$R_{\omega_i}^{\varepsilon_i \varepsilon_o} = T_{\varepsilon_o}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\varepsilon_i}$$

Relation to absorption

- Interaction of Neutrons with matter is easy (magnetic dipole) but for resonant x-ray scattering nothing is obvious. A lot is known though...

$$\epsilon_0^* \cdot \sigma \cdot \epsilon_i = \langle i | R_{\omega_i, \mathbf{q}}^{\epsilon_i \epsilon_0} | i \rangle$$

$$R_{\omega_i}^{\epsilon_i \epsilon_0} = T_{\epsilon_0}^\dagger \frac{1}{\omega_i + E_i + i\Gamma/2 - H} T_{\epsilon_i}$$



Effective scattering operator for magnetic spin excitations

M. W. Haverkort

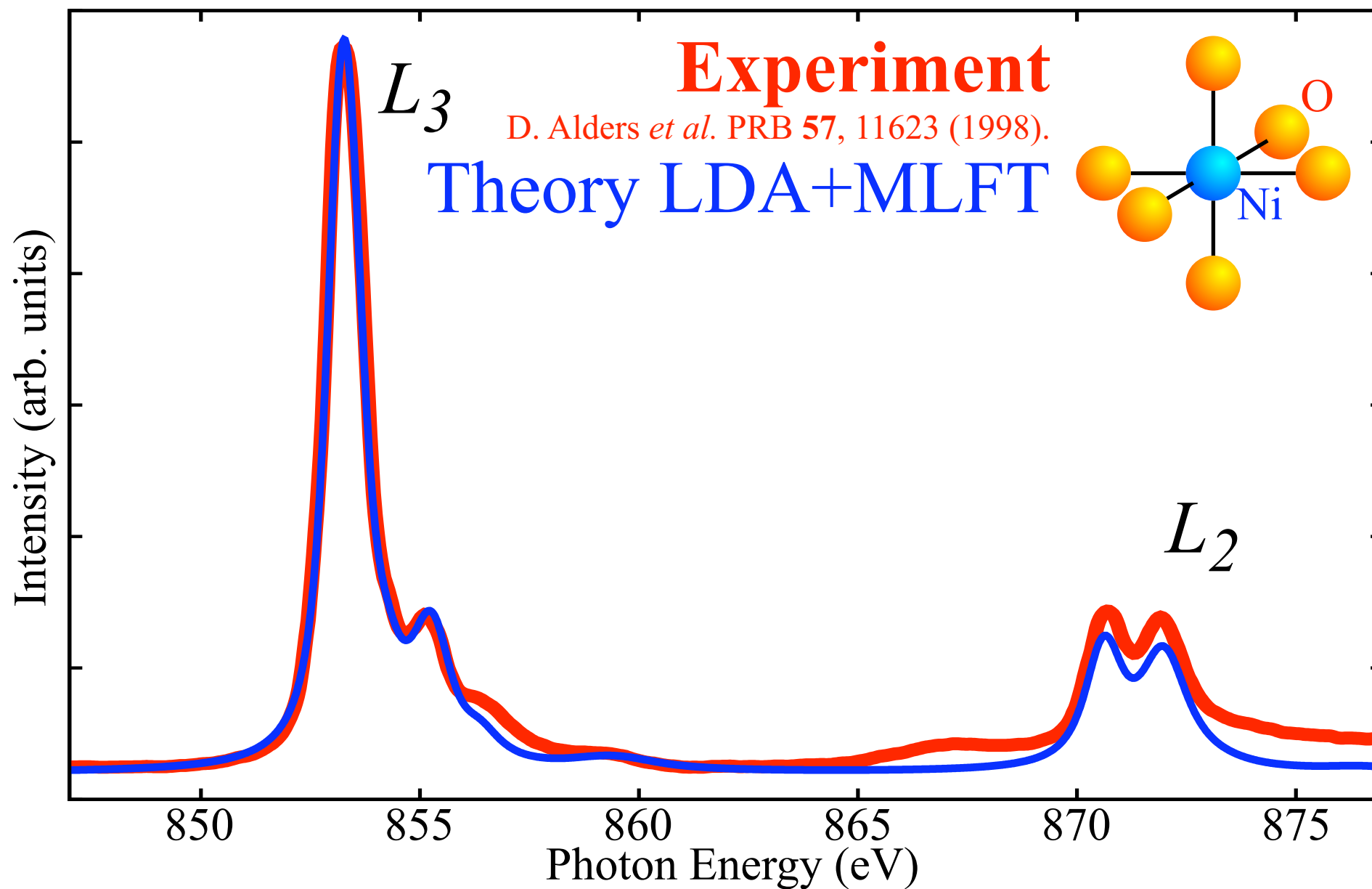
Phys. Rev. Lett. **105**, 167404 (2010)

One needs to calculate x-ray absorption spectra

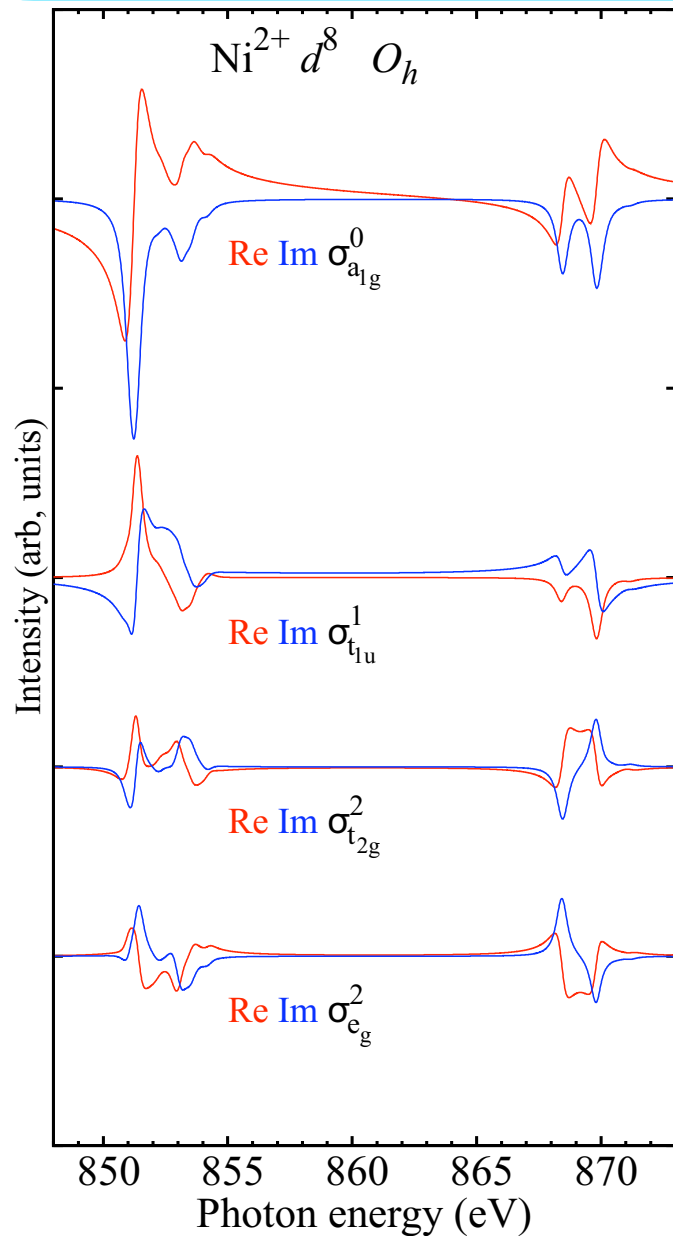
... and spin susceptibilities

$$\begin{aligned} R_{\omega_i, j}^{\varepsilon_i \varepsilon_o} &= \sigma^{(0)} \varepsilon_i \cdot \varepsilon_o^* + \frac{\sigma^{(1)}}{s} \varepsilon_o^* \times \varepsilon_i \cdot S_j \\ &+ \frac{\sigma^{(2)}}{s(2s-1)} (\varepsilon_i \cdot S_j \varepsilon_o^* \cdot S_j + \varepsilon_o^* \cdot S_j \varepsilon_i \cdot S_j - \frac{2}{3} \varepsilon_i \cdot \varepsilon_o^* S_j^2) \end{aligned}$$

2p X-ray Absorption Spectroscopy of NiO



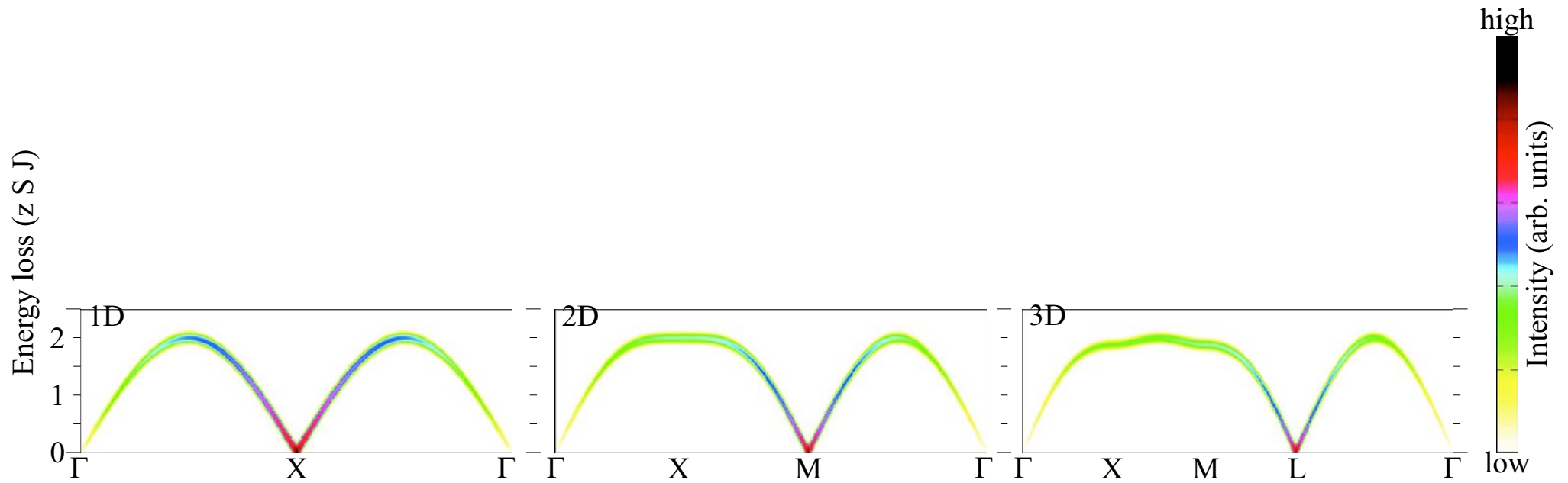
σ for Ni^{2+} in O_h



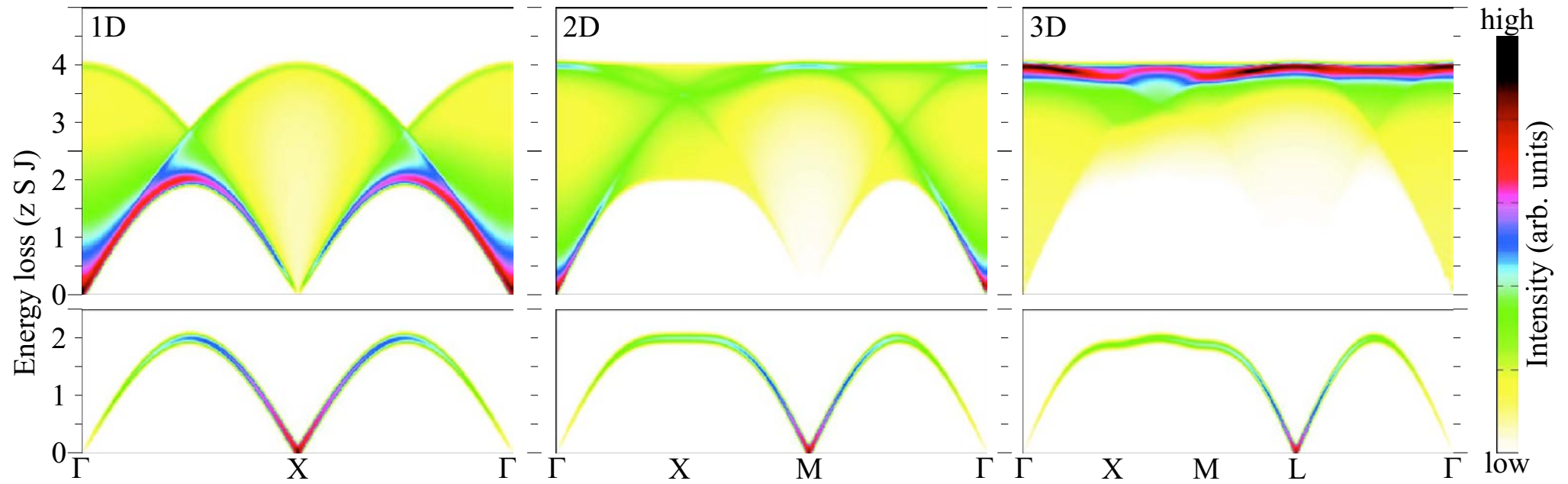
$$R_j^{\epsilon_i \epsilon_o} = \sigma^{(0)} \epsilon_i \cdot \epsilon_o^* + \frac{\sigma^{(1)}}{s} \epsilon_i \times \epsilon_o^* \cdot S$$

$$+ \frac{\sigma^{(2)}}{s(2s-1)} (\epsilon_i \cdot S \epsilon_o^* \cdot S + \epsilon_o^* \cdot S \epsilon_i \cdot S - \frac{2}{3} \epsilon_i \cdot \epsilon_o^* S^2)$$

Single spin-flips resonant enhancement proportional to XMCD

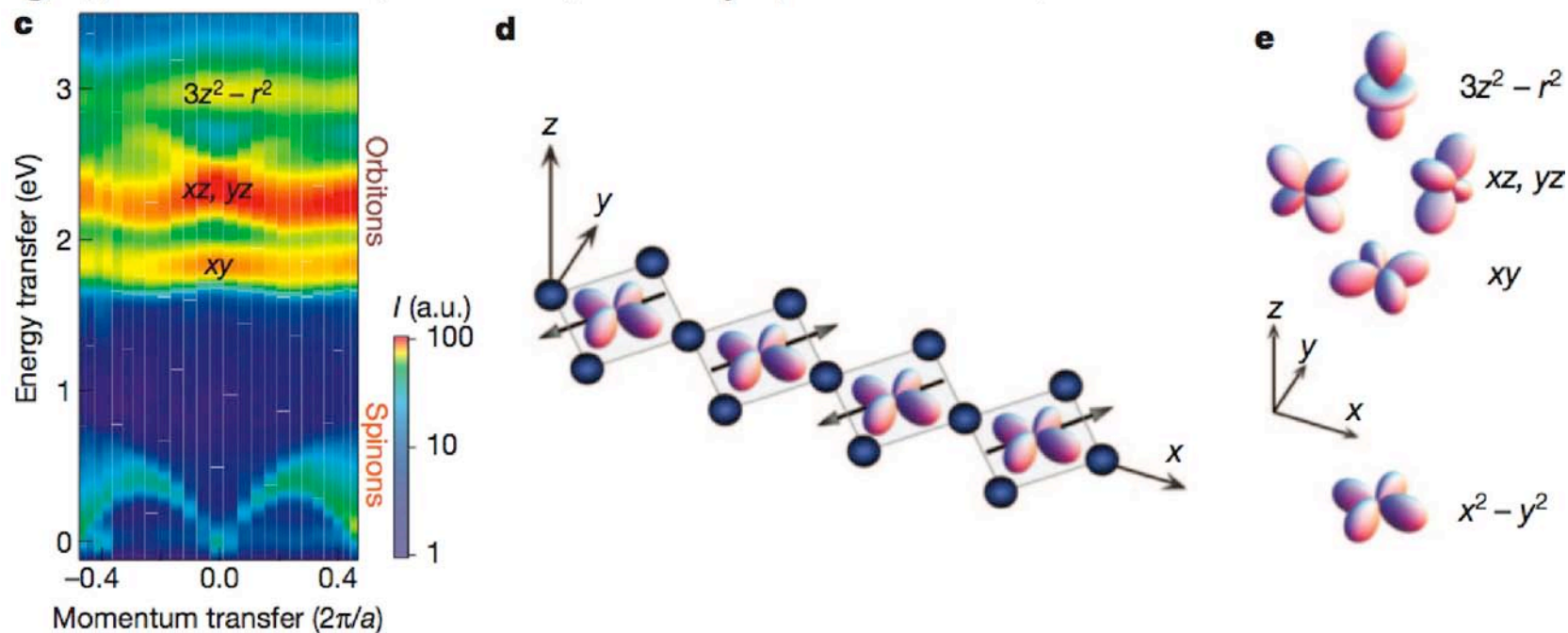


Two spin-flips resonant enhancement proportional to XMLD



Spin-orbital separation in the quasi-one-dimensional Mott insulator Sr_2CuO_3

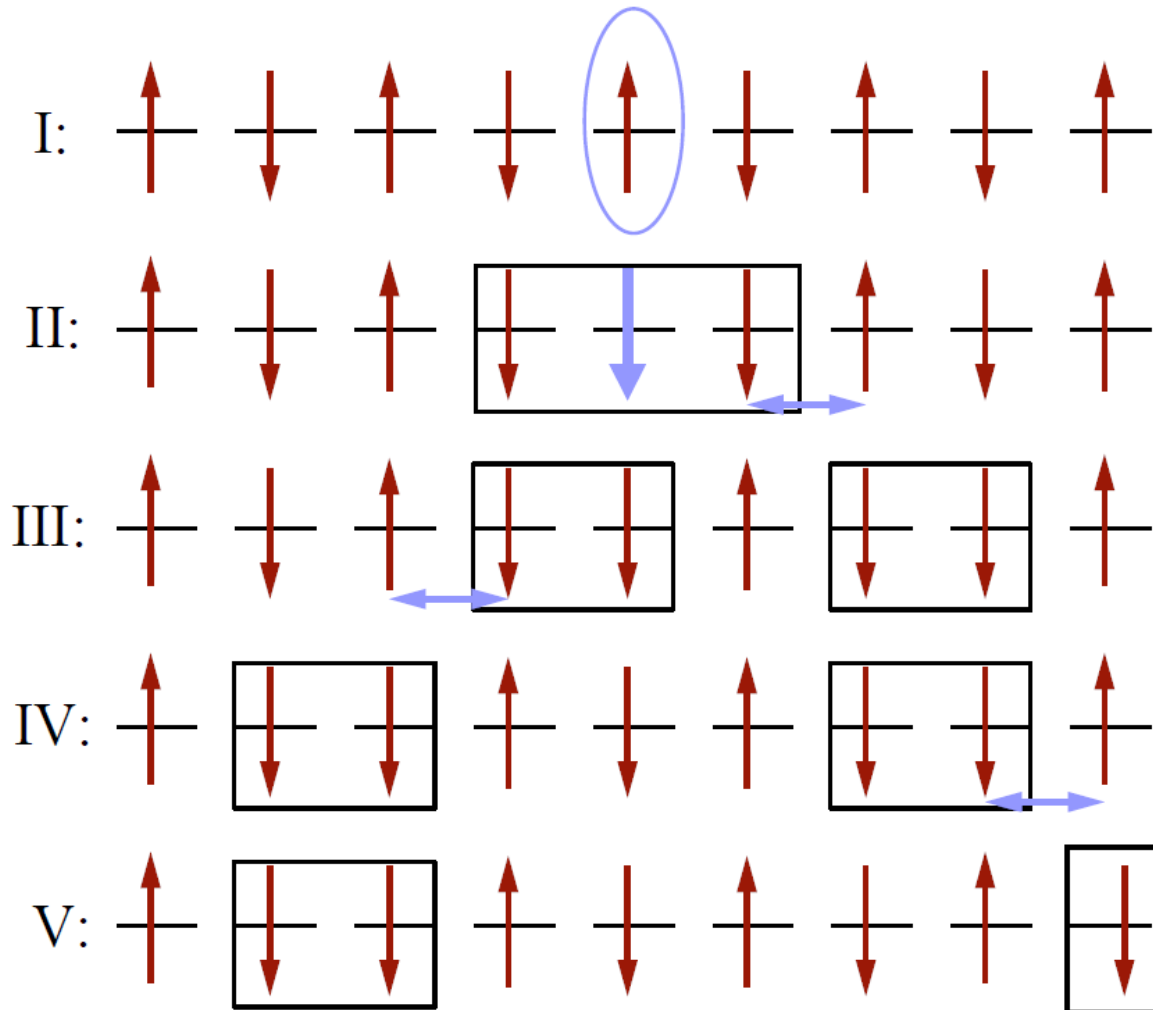
J. Schlappa^{1,2}, K. Wohlfeld³, K. J. Zhou¹, M. Mourigal⁴, M. W. Haverkort⁵, V. N. Strocov¹, L. Hozoi³, C. Monney¹, S. Nishimoto³, S. Singh^{6†}, A. Revcolevschi⁶, J.-S. Caux⁷, L. Patthey^{1,8}, H. M. Rønnow⁴, J. van den Brink³ & T. Schmitt¹



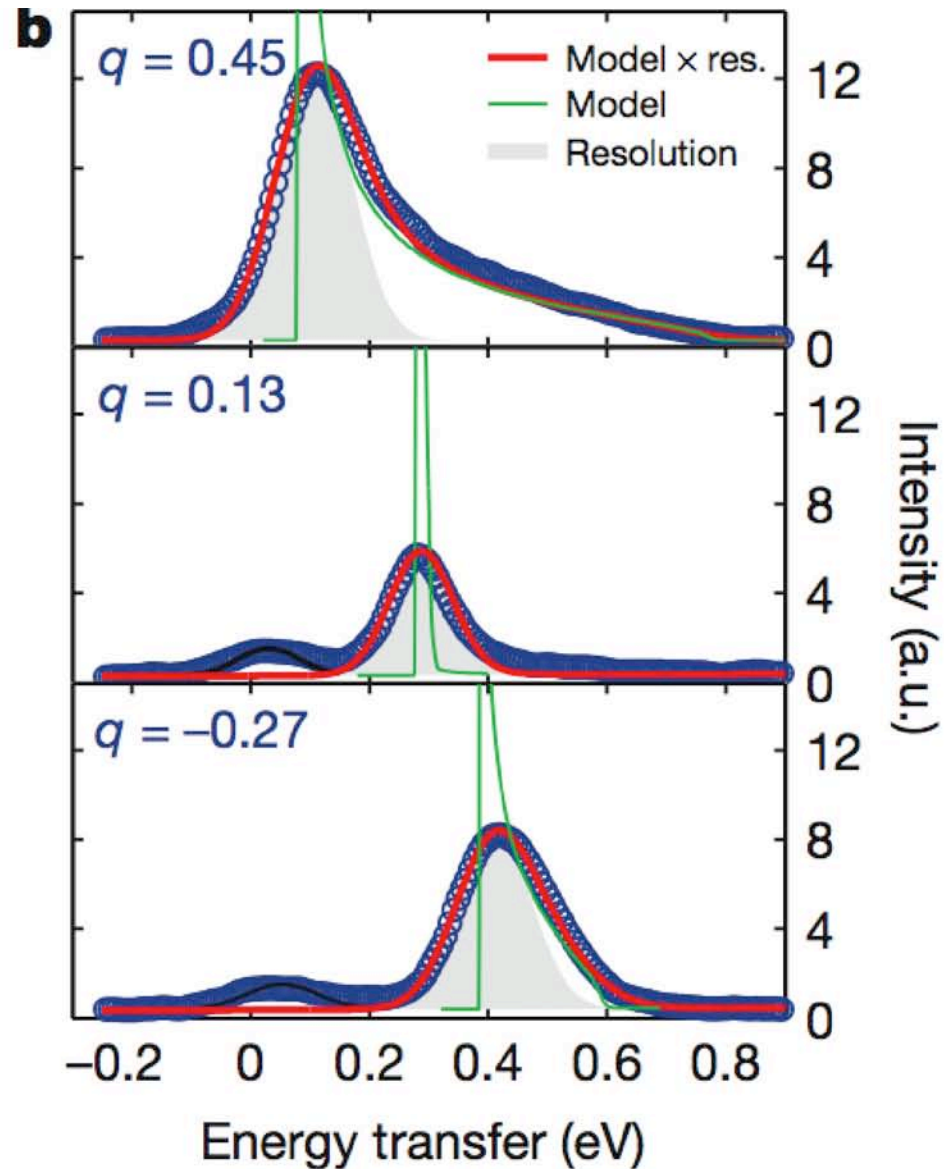
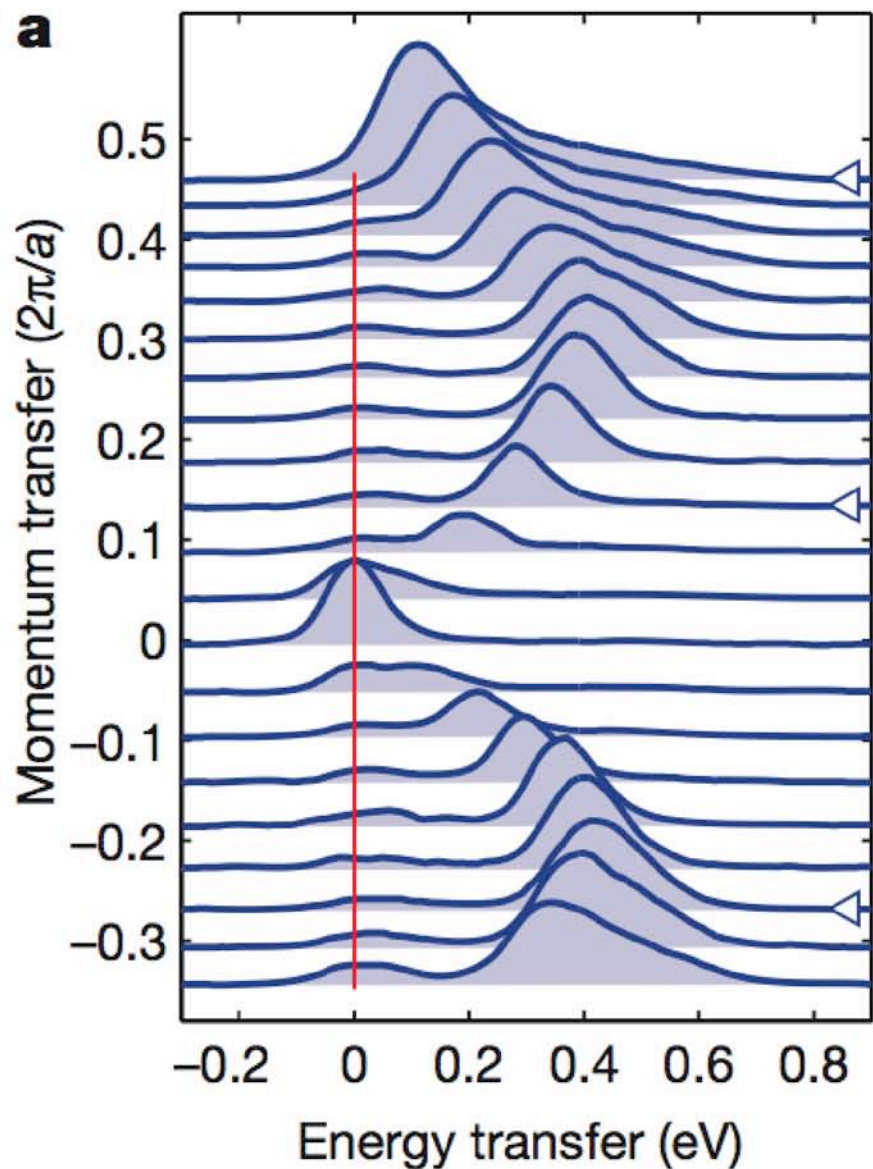
Spin and orbital excitations in Sr_2CuO_3

Dynamical spin structure factor

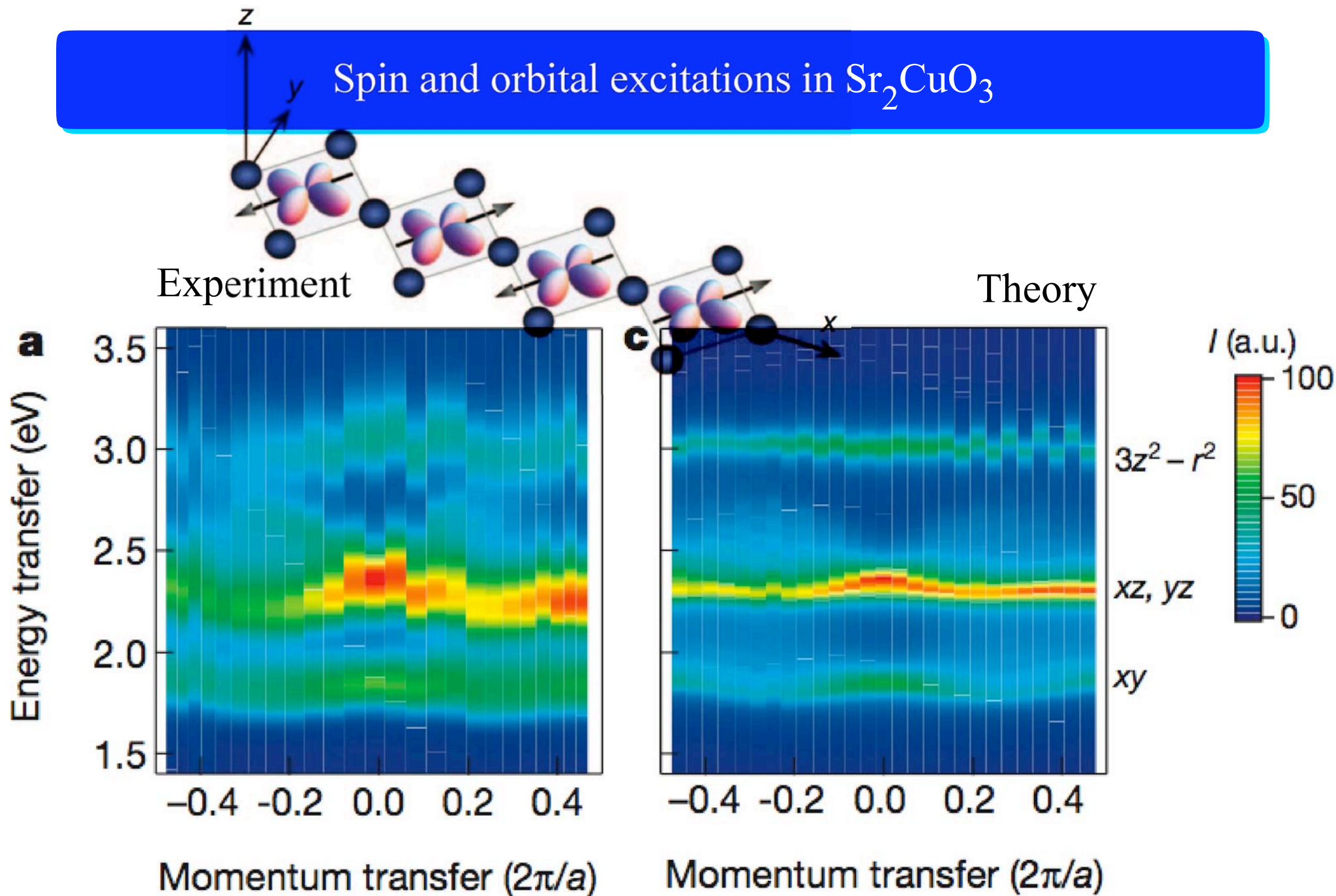
For a 1-dimensional spin chain the fundamental excitations are 2 spinon excitations



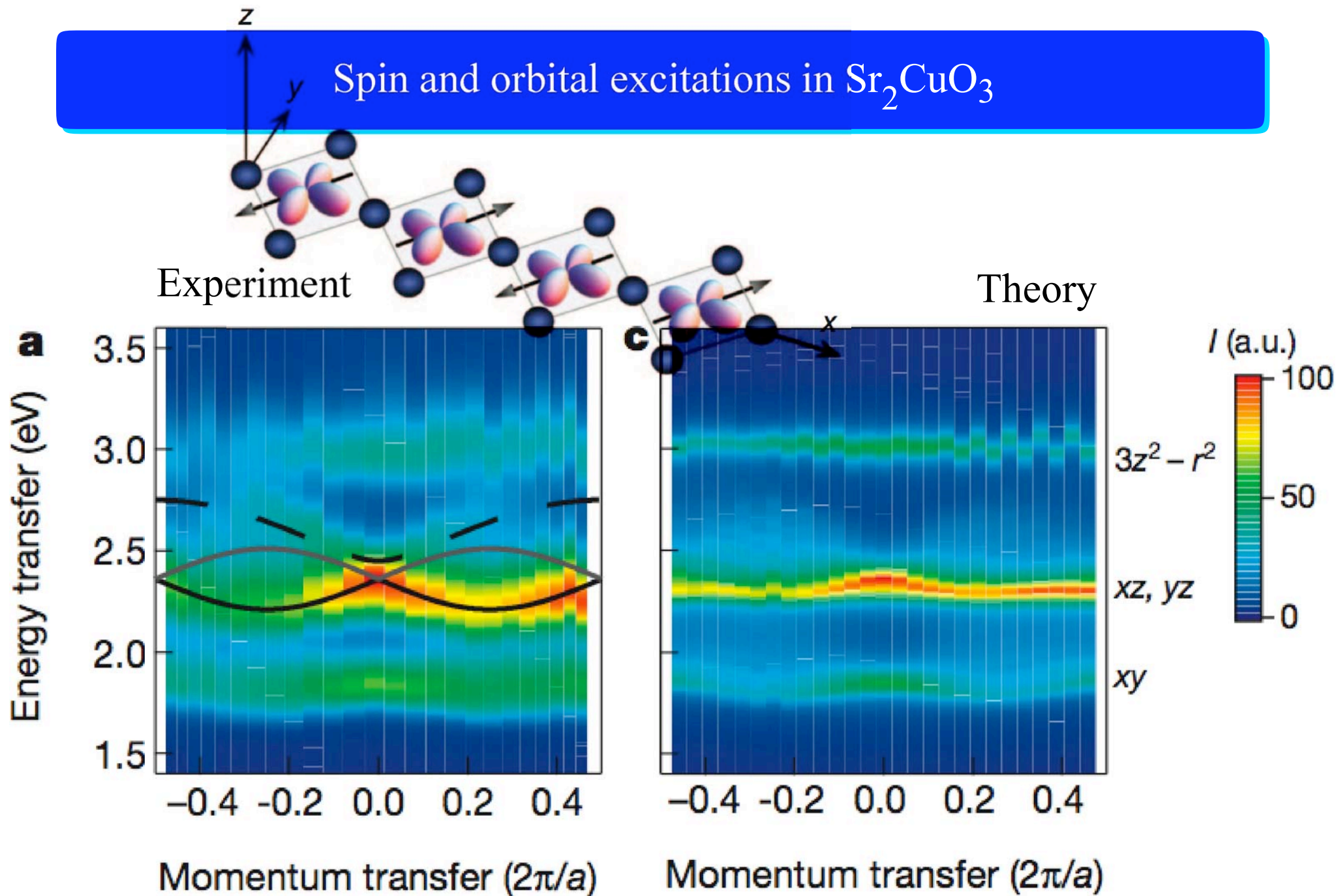
Spin and orbital excitations in Sr_2CuO_3



Spin and orbital excitations in Sr_2CuO_3



Spin and orbital excitations in Sr_2CuO_3



Additional literature

- ❑ *Core Level Spectroscopy of Solids*
Frank de Groot and Akio Kotani
- ❑ *NEXAFS Spectroscopy*
Joachim Stöhr
- ❑ *Journal of Electron Spectroscopy and Related Phenomena*
Vol **86** (1997).
- ❑ *Introduction to ligand field theory*
Carl Johan Ballhausen (1962)