



Dark Matter: Candidates I

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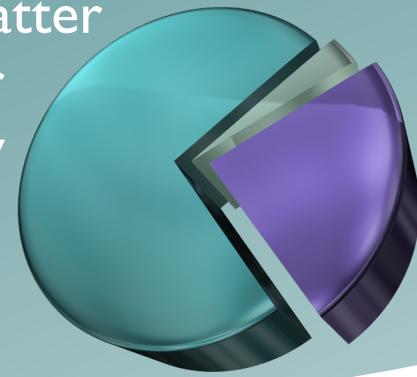
SLAC Summer Institute
July 28, 2011

Outline of Lecture I

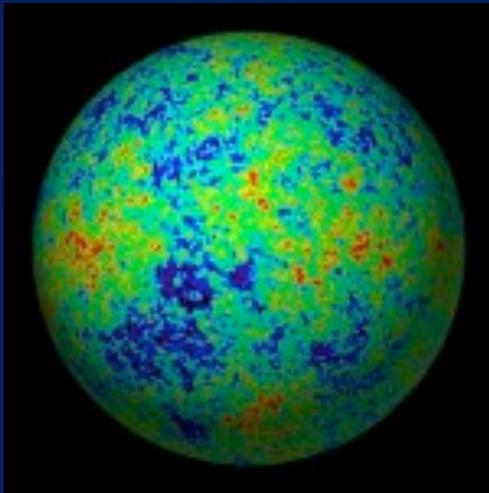
- Characteristics of a Dark Matter Candidate Particle
 - Stabilization
 - Relic Density
- WIMPs
 - R-parity: The SUSY WIMP
 - Relic Density
 - Indirect Signals
 - Direct Detection
 - Colliders

Dark Matter

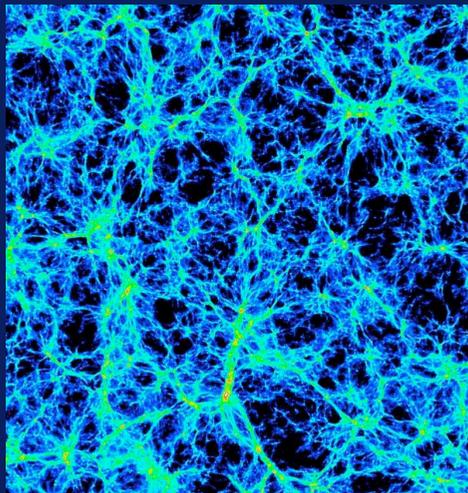
- Ordinary Matter
- Dark Matter
- Dark Energy



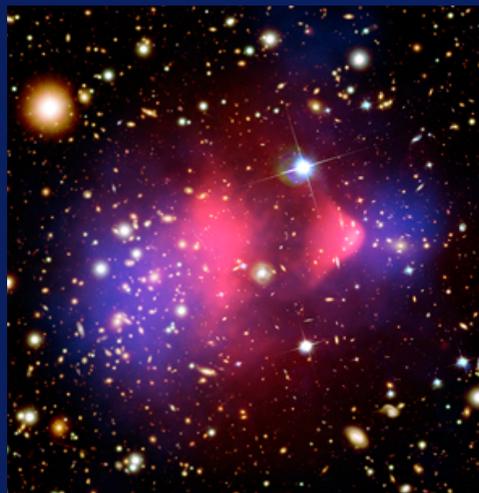
CMB



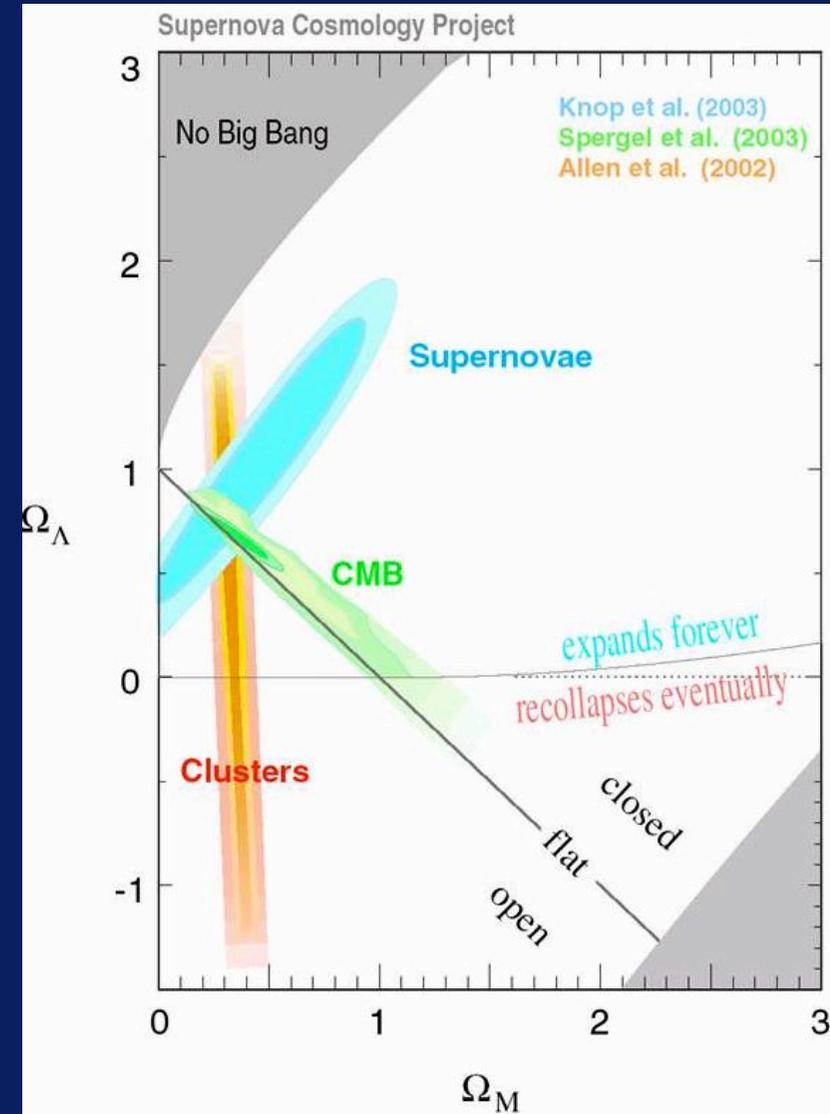
Supernova



Structure



Lensing



So what is this stuff?



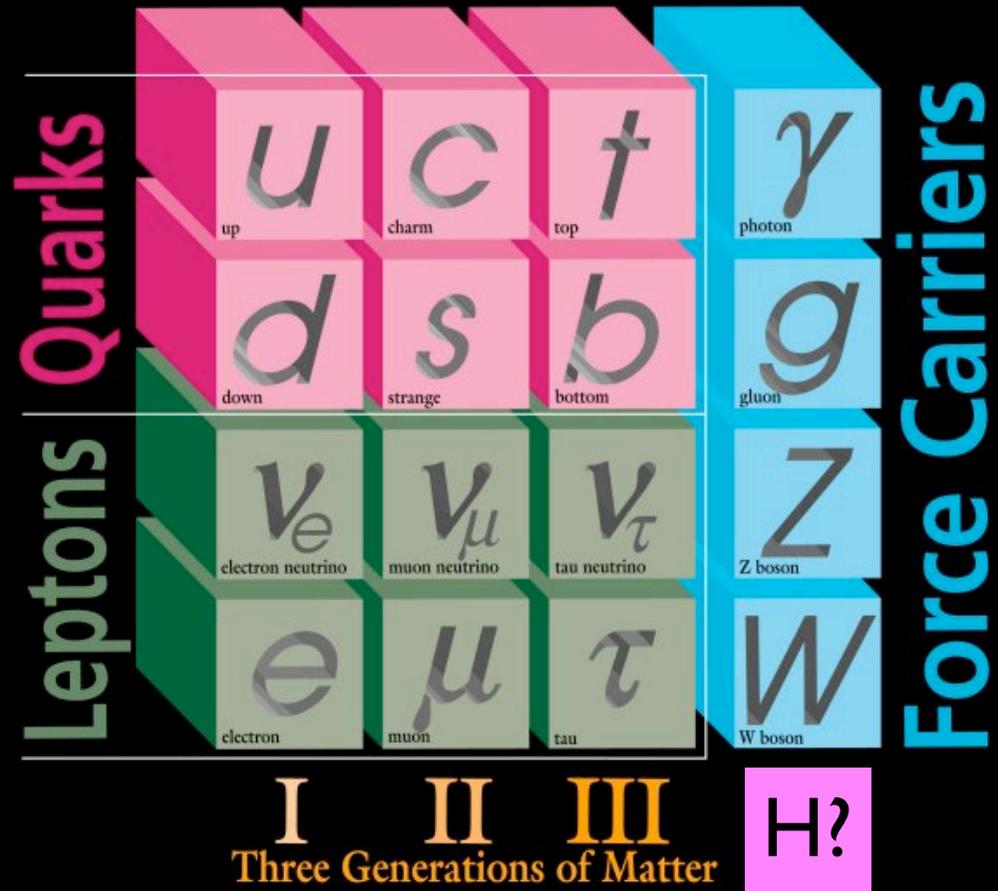
“Cold Dark Matter: An Exploded View” by Cornelia Parker

- As a particle physicist, my job is to explore how dark matter fits into the bigger picture of particles.
- What do we know about dark matter?
 - Dark (neutral)
 - Massive
 - Still around today
 - Stable or with a lifetime of the order of the age of the Universe itself).
- Nothing in the Standard Model of particle physics fits the description.

Physics Beyond the SM

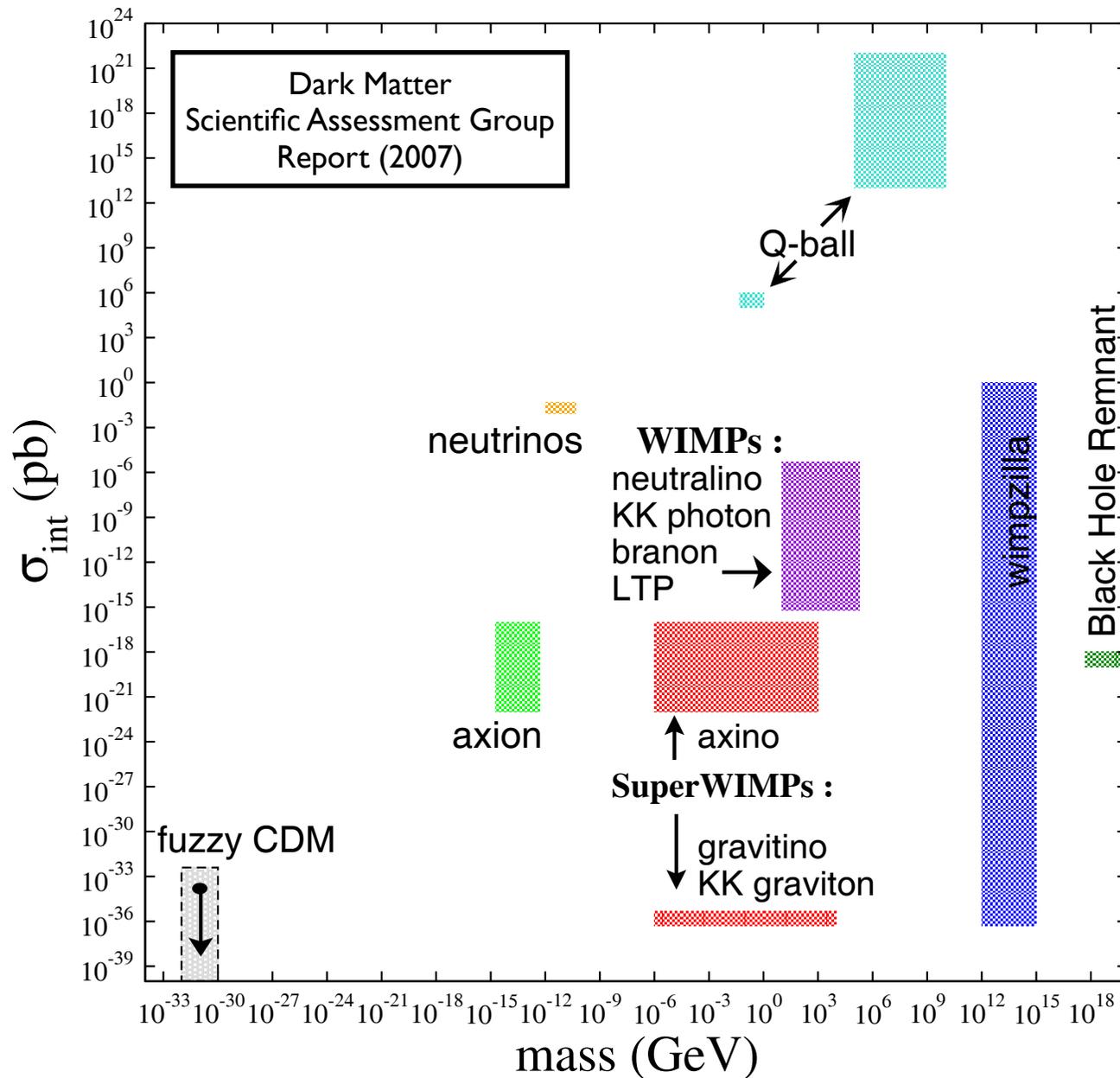
- The Standard Model of particle physics has nothing with the right properties to be dark matter:
- Photons, leptons, hadrons, and W bosons all shine too brightly.
- Neutrinos are too light.
- Z and Higgs bosons are too short-lived.
- Dark matter is a manifestation of physics beyond the Standard Model.

ELEMENTARY PARTICLES



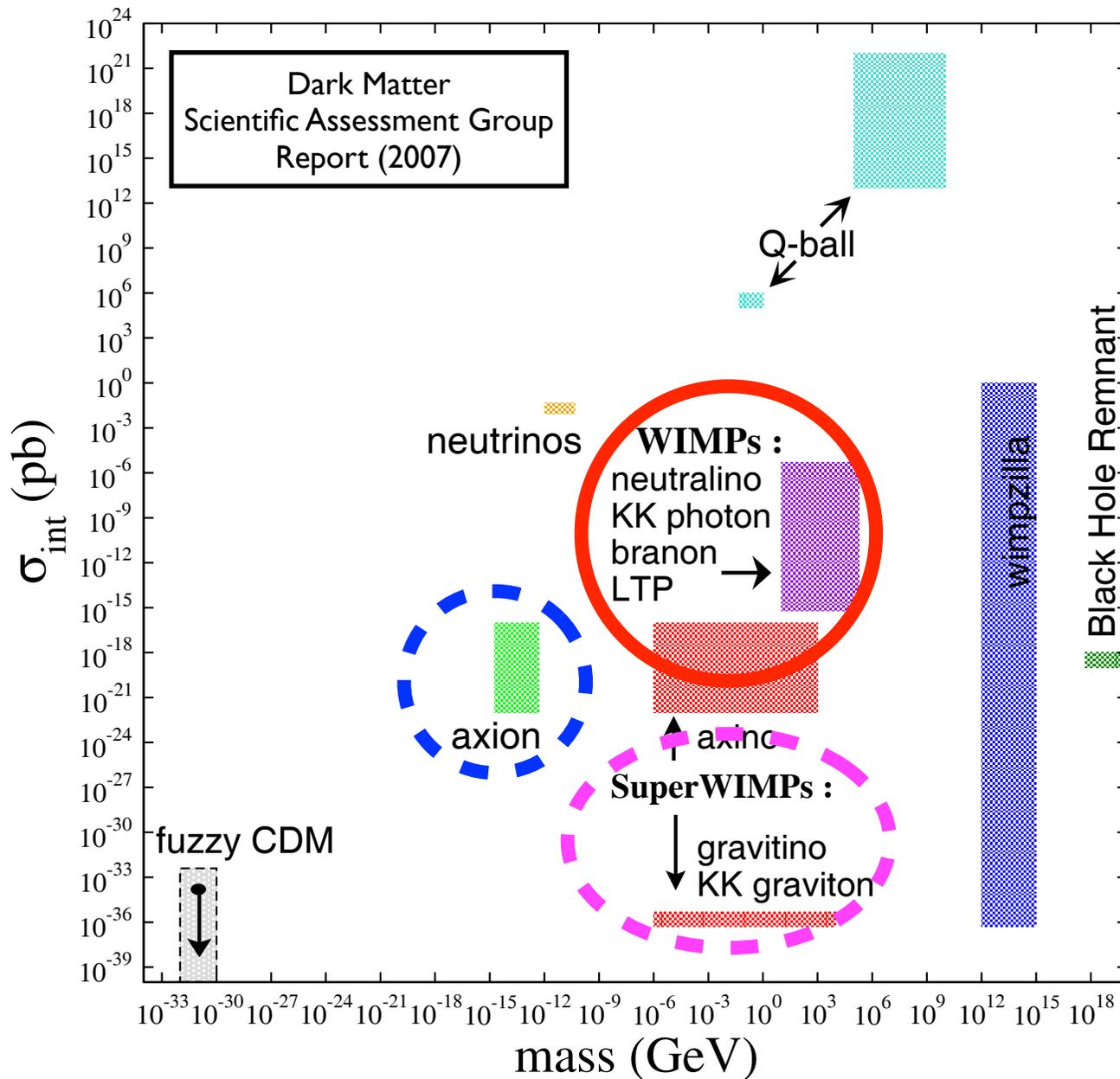
Many Ideas...

Some Dark Matter Candidate Particles



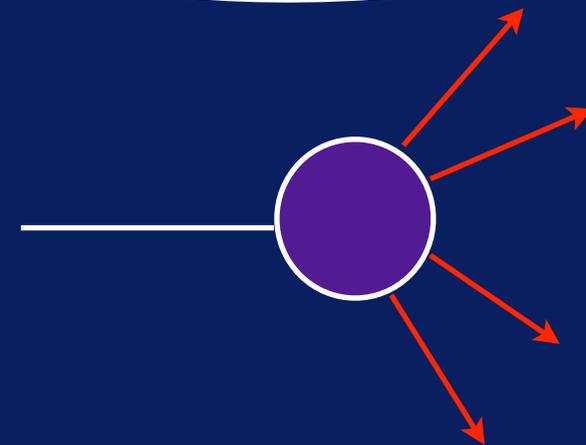
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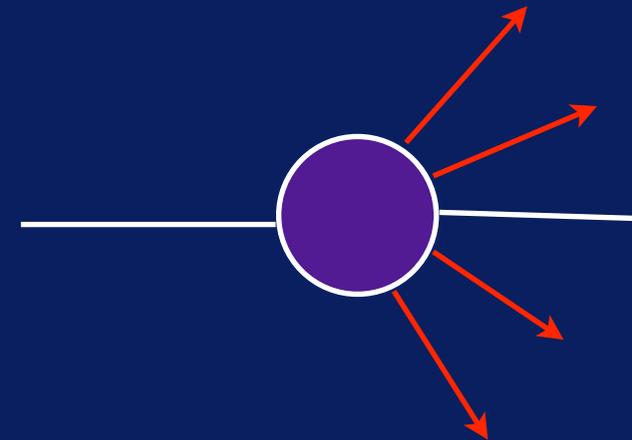


(Quasi) Stable

- One of the mysteries of dark matter is why it is very massive but (at least to very good approximation) stable.
- This is actually telling us something very important about how it can interact with the Standard Model.
- We need a symmetry (at least approximately) to prevent dark matter particles from decaying.
- The simplest example is a new kind of parity (a Z_2 discrete symmetry) which forces them to couple in pairs to SM fields.
- We could explore larger (and continuous) symmetries as well.



χ decays.



The number of χ 's is conserved.

WIMPs

One of the most attractive proposals for dark matter is that it is a **W**eakly **I**nteracting **M**assive **P**article.

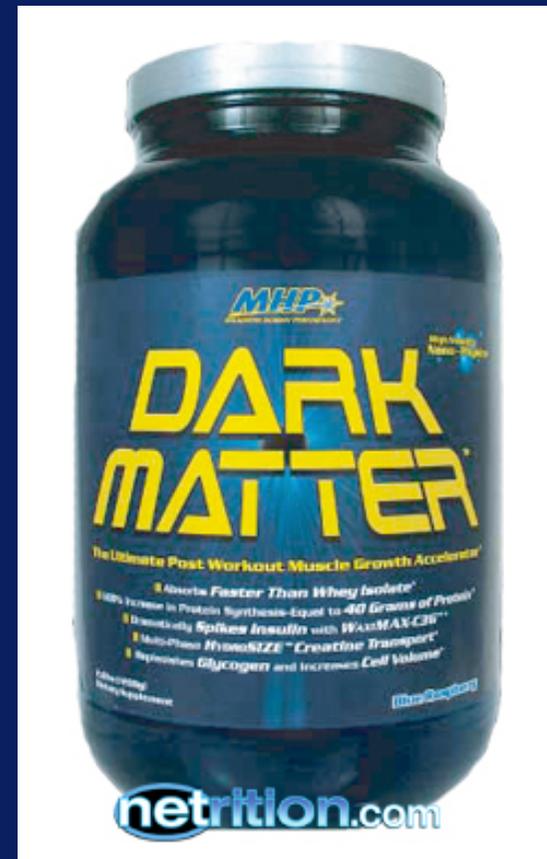
WIMPs naturally can account for the amount of dark matter we observe in the Universe.

WIMPs automatically occur in many models of physics beyond the Standard Model, such as i.e. supersymmetric extensions.

WIMPs are a vision of dark matter for which we can use particle physics experimental techniques to search very effectively.

Are we looking under the lamp post?

We will classify different WIMPs based on which symmetry allows them to be stable.



\$59.99 for 20 servings

Available in Blue Raspberry, Fruit Punch, and Grape flavors....

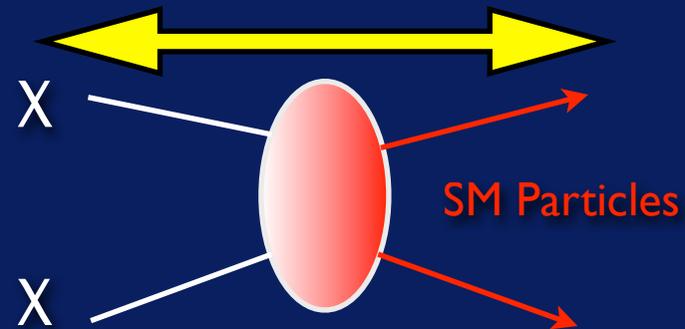
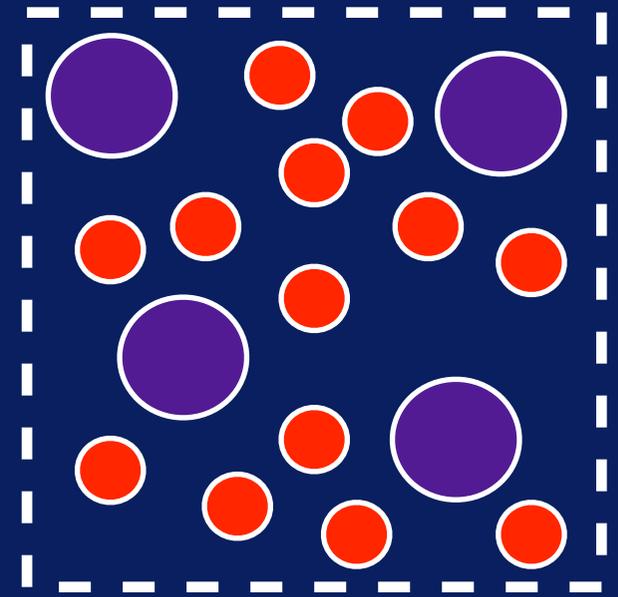
The WIMP Miracle

One of the primary motivations for WIMPs is the “WIMP miracle”, an attractive picture explaining the density of dark matter in the Universe today.

While not strictly a requirement for a successful theory of dark matter, this picture is very attractive [meaning: we think it is likely that things work this way], and so it is worth understanding the argument.

The picture starts out with the WIMP in chemical equilibrium with the Standard Model plasma at early times.

Equilibrium is maintained by scattering of WIMPs into SM particles, $\chi\chi \rightarrow \text{SM}$ and vice-versa.



Boltzmann Equation

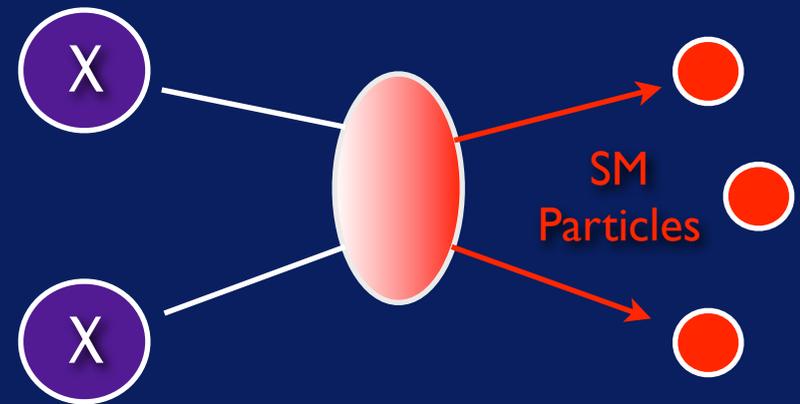
The evolution of the dark matter number density (n) is controlled by a Boltzmann equation, which tracks the effect of the expansion of the Universe (H) and the creation and destruction of dark matter.

A Universe where WIMPs stayed in equilibrium would be pretty boring.

As the temperature falls, there will be fewer and fewer WIMPs present, since the fraction of the plasma with enough energy to produce them will become smaller and smaller.

(Almost) Nothing would be left!

$$\frac{dn}{dt} + 3Hn = -\langle\sigma v\rangle [n^2 - n_{eq}^2]$$



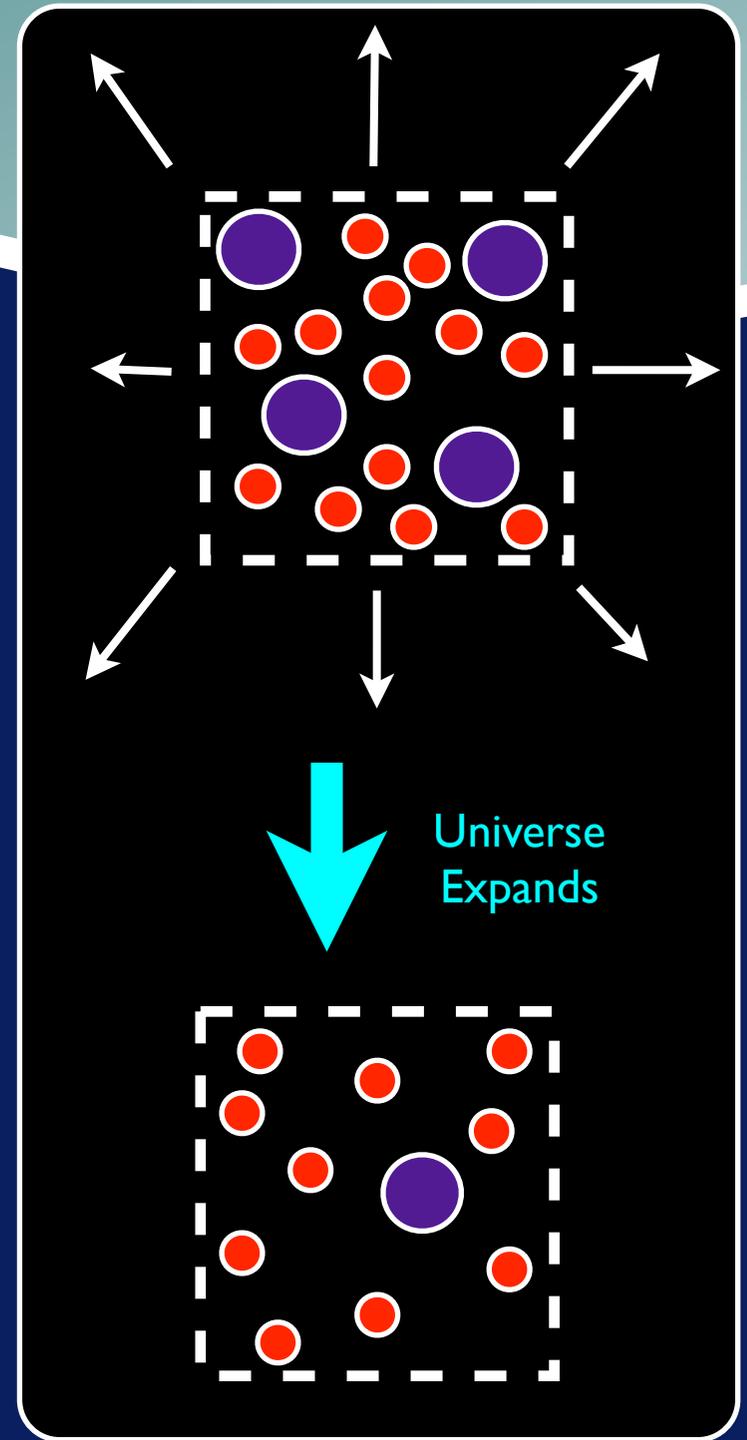
$$n_{eq} = g \left(\frac{mT}{2\pi} \right)^{3/2} \text{Exp} [-m/T]$$

Freeze-Out

- However, the expansion of the Universe eventually results in a loss of equilibrium.

$$\frac{dn}{dt} + 3Hn = -\langle\sigma v\rangle [n^2 - n_{eq}^2]$$

- When $(n_{eq} \langle\sigma v\rangle) \ll H$, the scattering that maintains equilibrium can't keep up with the expansion.
- The WIMPs become sufficiently diluted that they can no longer find each other to annihilate and they cease tracking the Boltzmann distribution.
- Where they “freeze out” obviously depends on how big $\langle\sigma v\rangle$ is.



Relic Density

So the basic picture is:

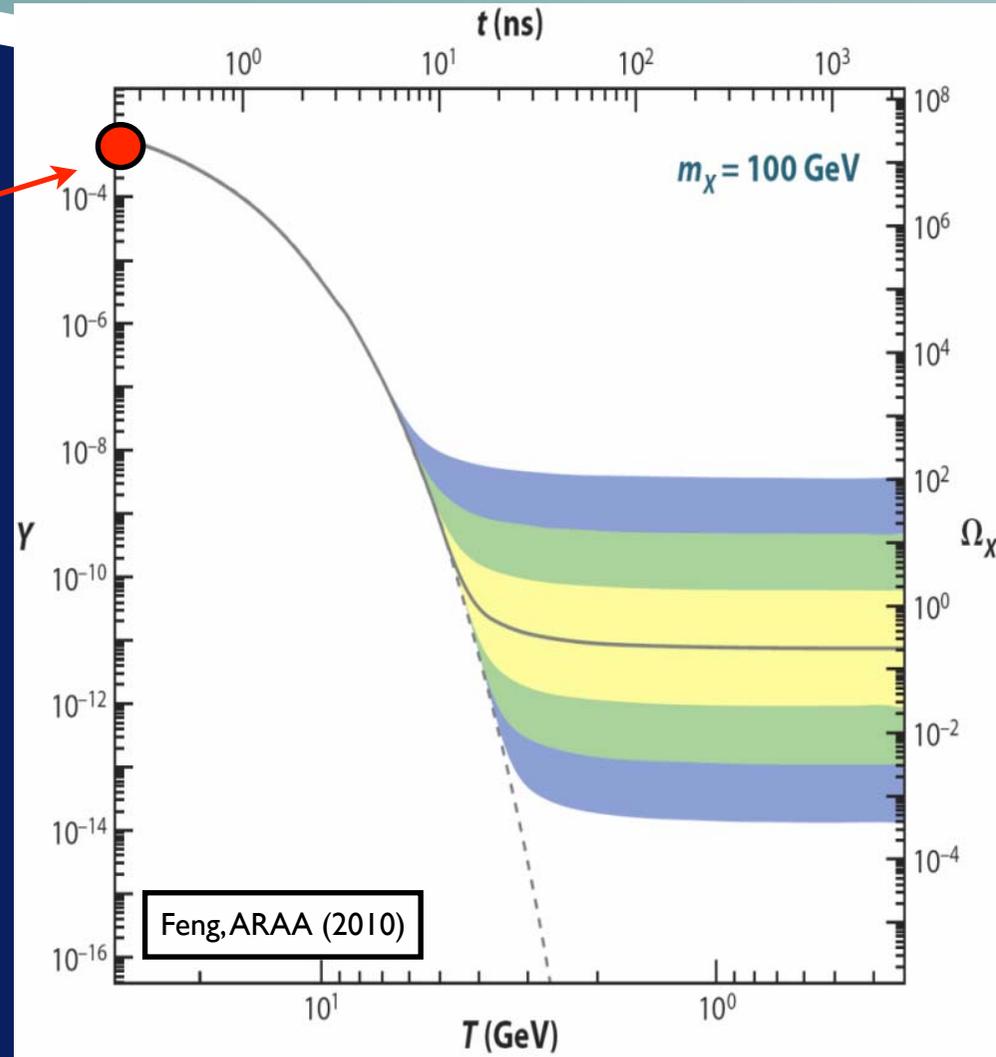
We start out with dark matter in equilibrium with the SM plasma.

As the temperature falls, the number of WIMPs does too.

We track the equilibrium density until freeze-out:

$$n_{eq} \langle \sigma v \rangle \sim H$$

$$\begin{matrix} \swarrow & \searrow & \searrow \\ (mT)^{3/2} e^{-m/T} & \frac{g^4}{m^2} & \frac{T^2}{M_{Pl}} \end{matrix}$$



$$\frac{m}{T} \sim \log \left[\frac{M_{Pl}}{m} \right] \quad m \sim 100 \text{ GeV} : \frac{m}{T} \sim 40$$

...and that's how much dark matter we get!

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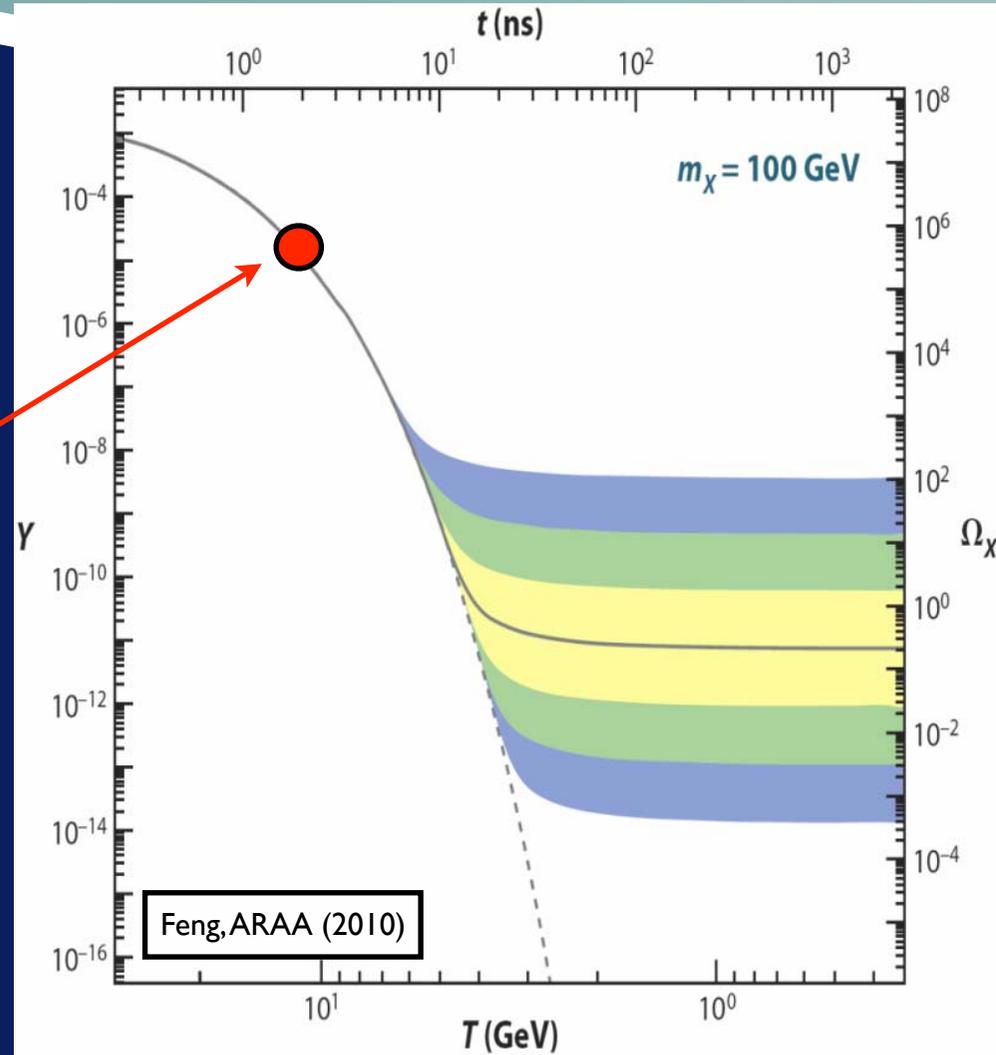
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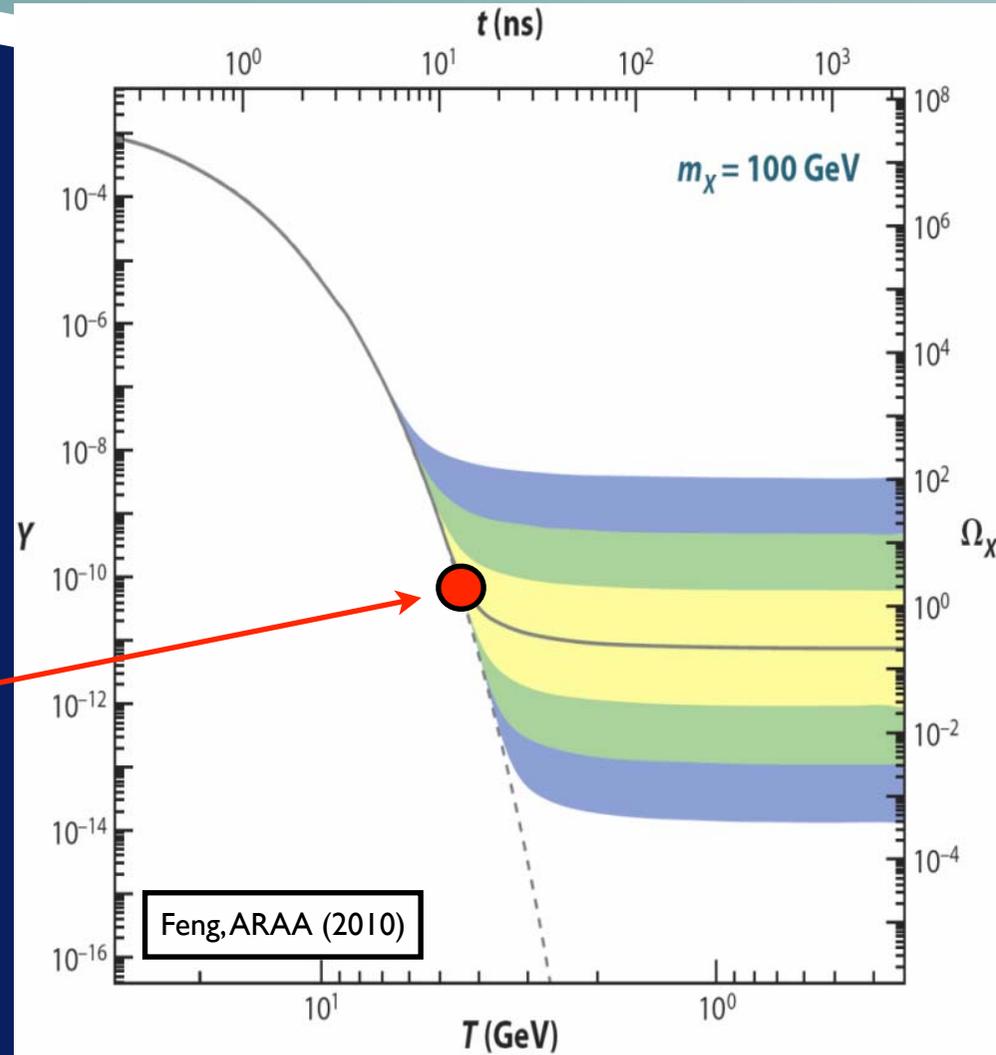
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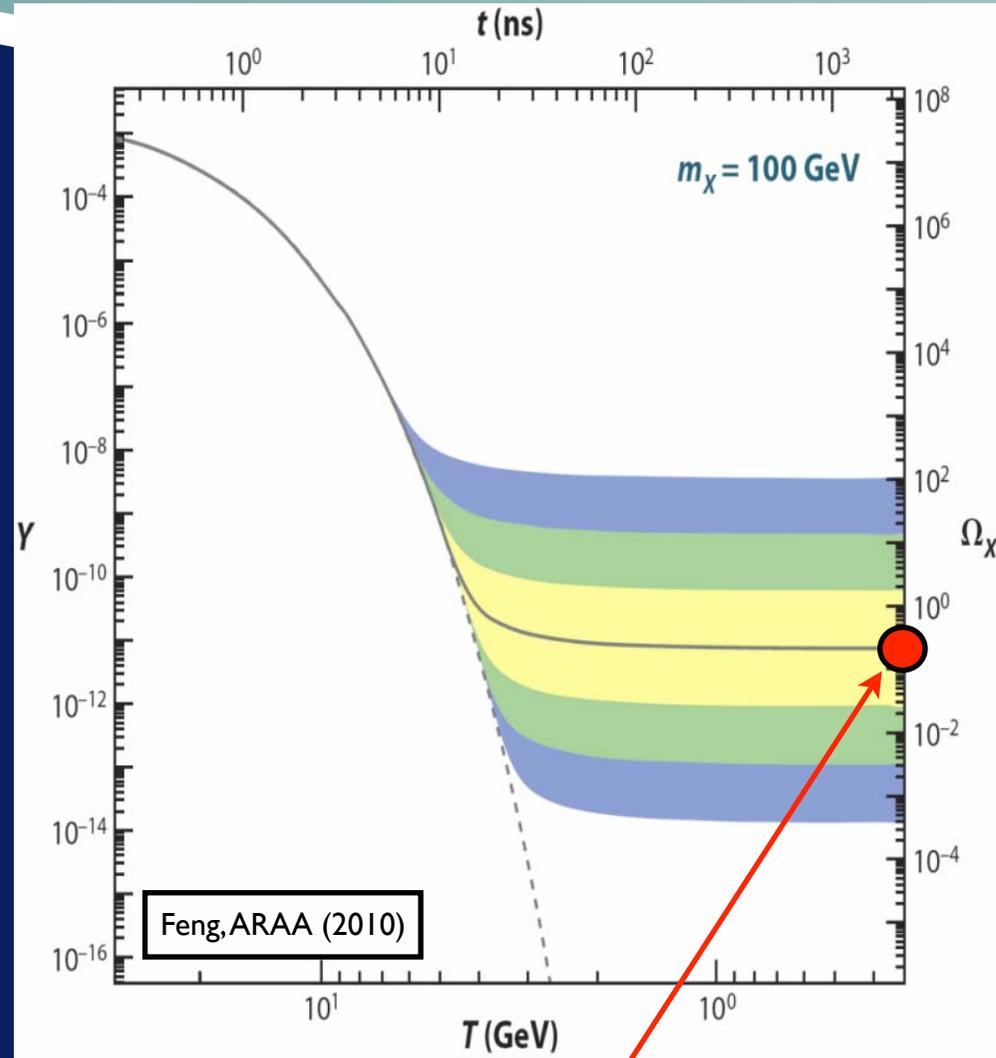
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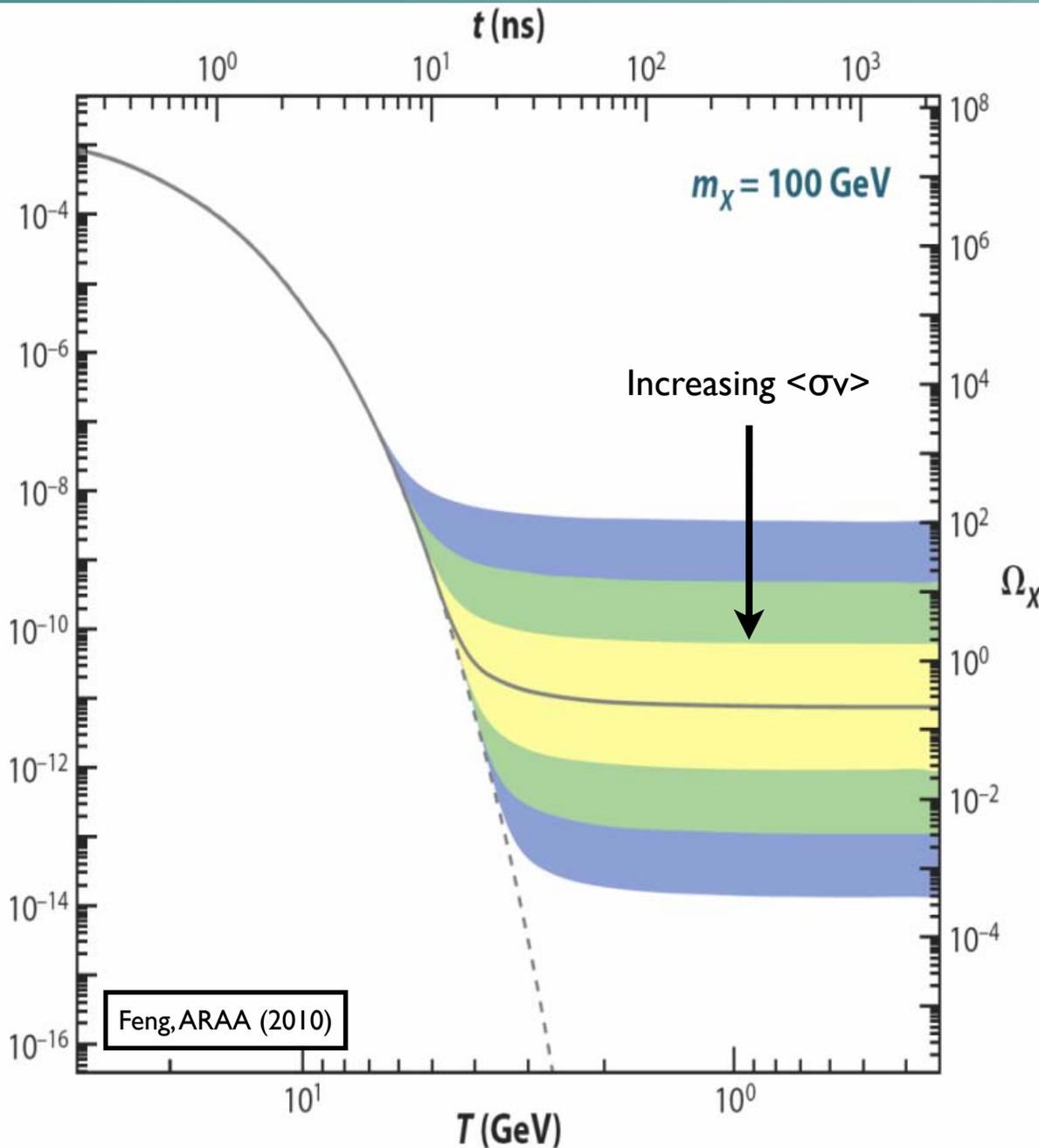
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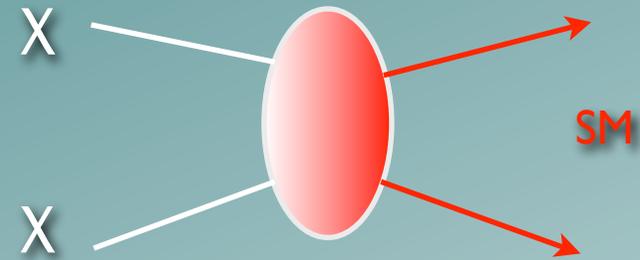
Relic Density



For a WIMP, once we know its mass and cross section into SM particles, we can predict its relic density.

I find it remarkable that one simple, reasonable assumption (DM in equilibrium with the SM at early times) is enough to predict the dark matter density today in terms of the particle physics properties of DM.

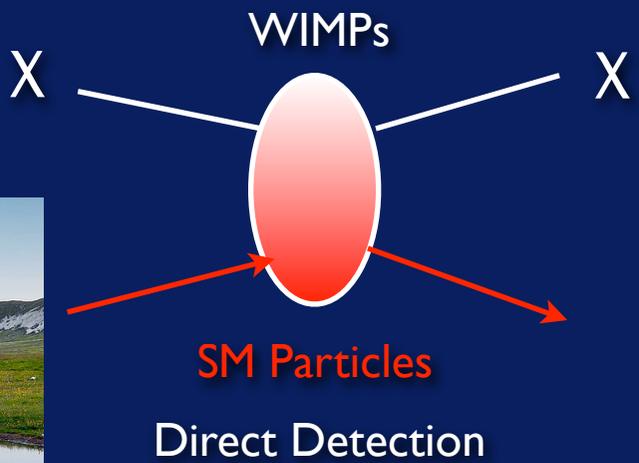
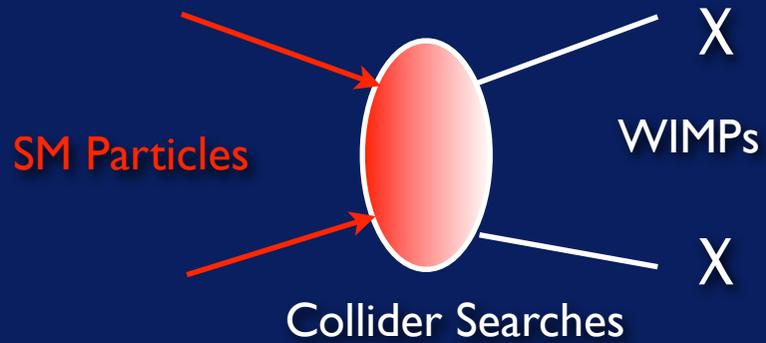
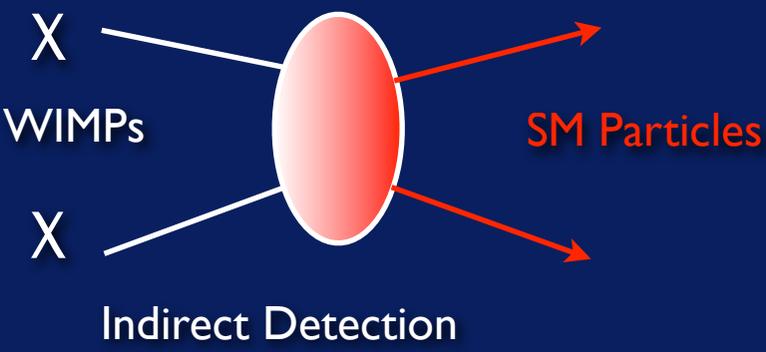
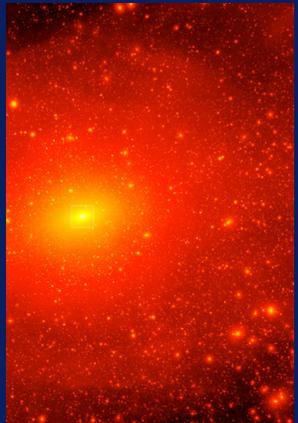
WIMP Interactions



- Ideally, we would like to measure WIMP interactions with the Standard Model, allowing us to compute $\sigma(\chi\chi \rightarrow \text{SM particles})$ and check the relic density.
- If our predictions “check out” we have indirect evidence that our extrapolation backward to higher temperatures is working.
- If not, we will look for signs of new physical processes to make up the difference.
- The first step is to actually rediscover dark matter by seeing it interact through some force other than gravitational.
- That tells us which SM particles it likes to talk to and in some cases something about its spin, mass, etc.



WIMP Searches



WIMPs interacting with SM particles allow indirect searches for annihilation products, direct scattering searches, and production at colliders.

Catalogue of Candidates

So here is how we'll catalogue WIMPs:

Stability Mechanism

How they interact with the SM.

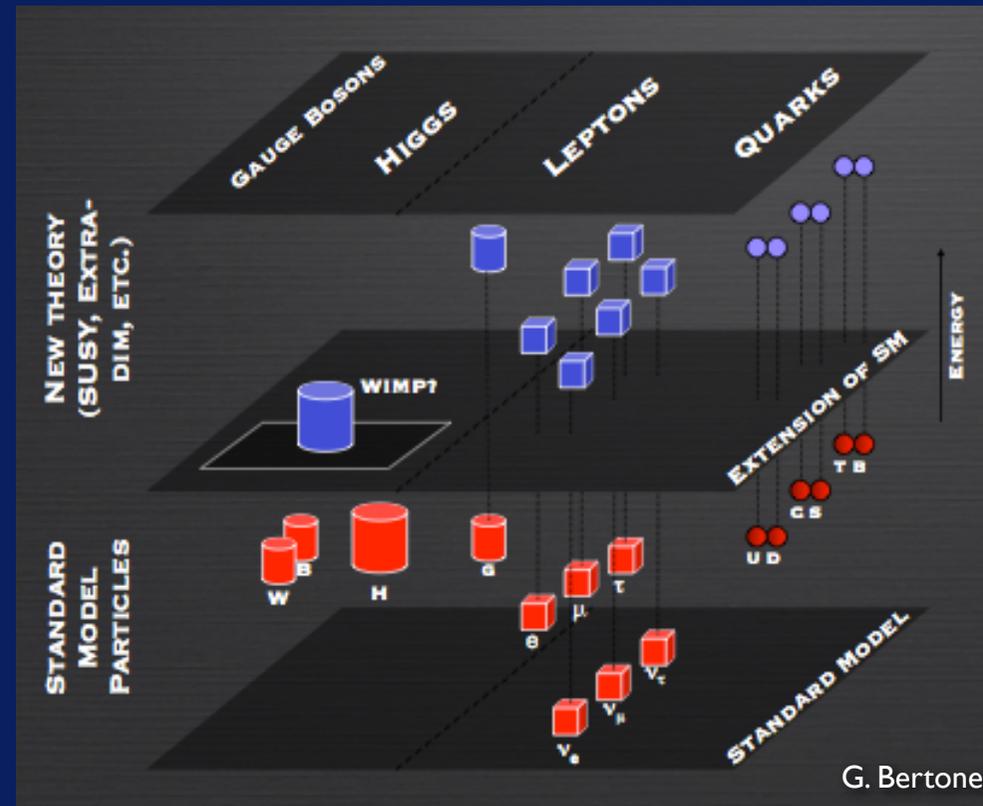
Relic density

Detection prospects

Direct

Indirect

Collider



G. Bertone

The picture that emerges will be that there are a lot of interesting ideas for DM -- and we can test them!

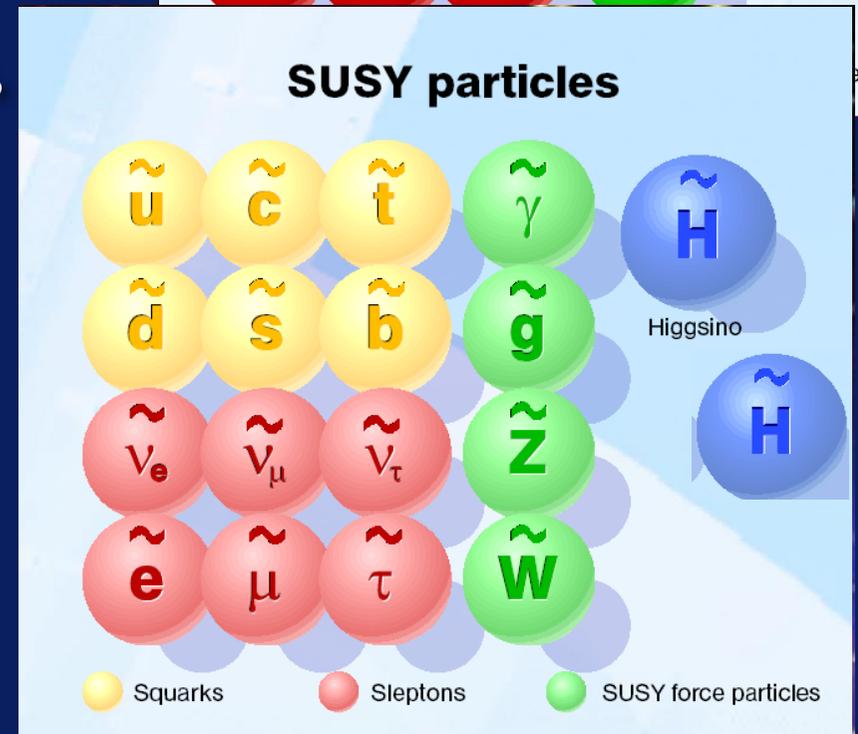
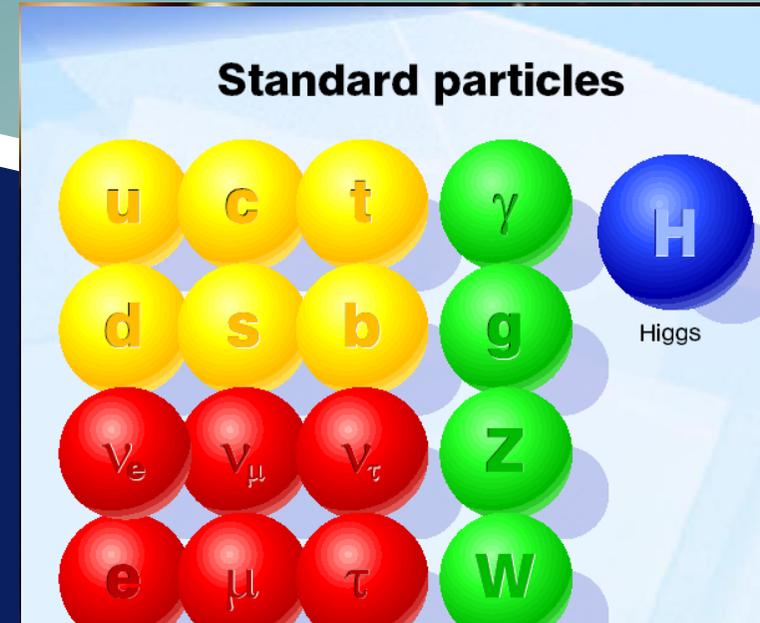
Supersymmetry (SUSY)

- The most famous candidate for dark matter is a supersymmetric particle.

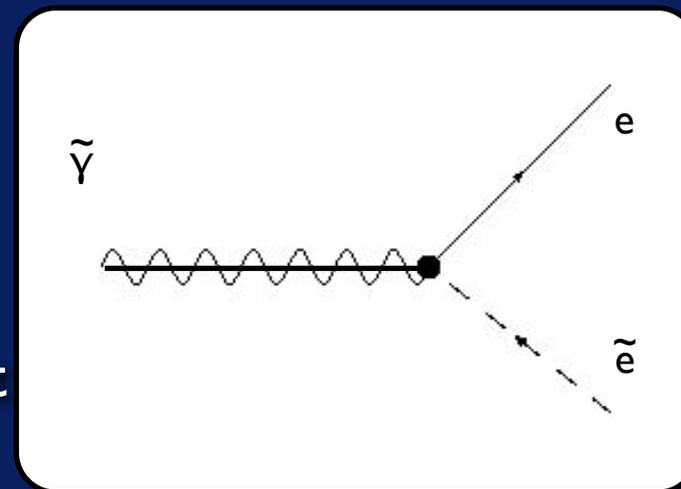
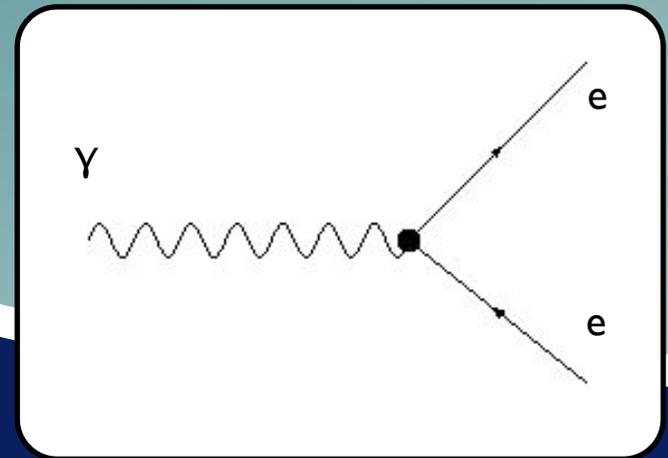
- Supersymmetry is an extension of Poincare invariance to include fermionic generators

- So in addition to the usual translations, rotations, and boosts, we have new transformations that 'rotate' fermions into bosons and vice versa.

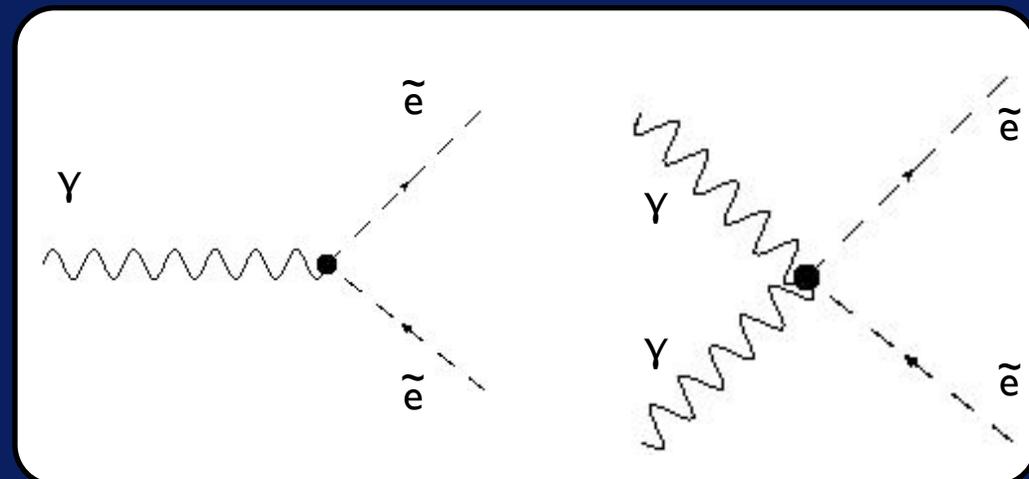
- Every particle needs a superpartner whose spin differs by 1/2. If SUSY were exact, the superpartners would have the same masses as their SM counterparts.



SUSY Interactions



$\propto \alpha_{EM}$



- If we break supersymmetry “softly”, the masses of the superpartners will separate, but the interactions remain fixed by supersymmetry.
- Despite having many, many new parameters, SUSY theories inherit a huge structure from the SM.
- This implies that many things can be calculated in supersymmetric theories in terms of the masses of the superpartners.
- See: Martin, hep-ph/9709356 for a complete introduction to SUSY.

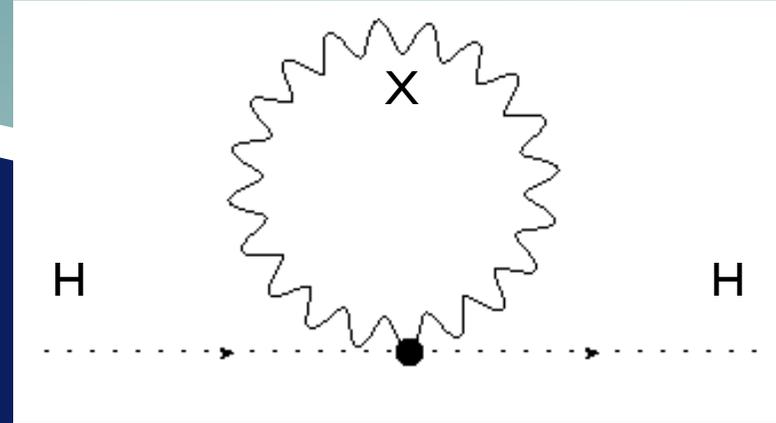
Motivations for SUSY

- SUSY has a lot of good motivation independently from its being a theory of dark matter.

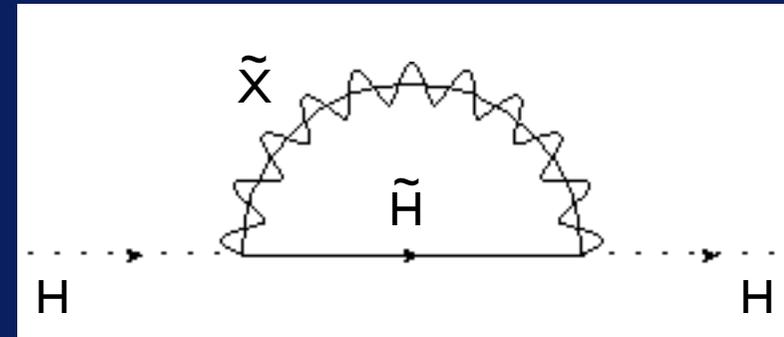
- The original motivation was to address the gauge hierarchy problem: the fact that the Higgs VEV is very sensitive to quantum corrections.

- If there is some heavy particle which couples to the Higgs (i.e. a GUT gauge boson), quantum corrections try to drive the Higgs mass to its own mass.

- SUSY cancels these large corrections by adding additional diagrams containing the superpartners.



$$\sim \frac{g_{\text{GUT}}^2}{16\pi^2} M_X^2$$

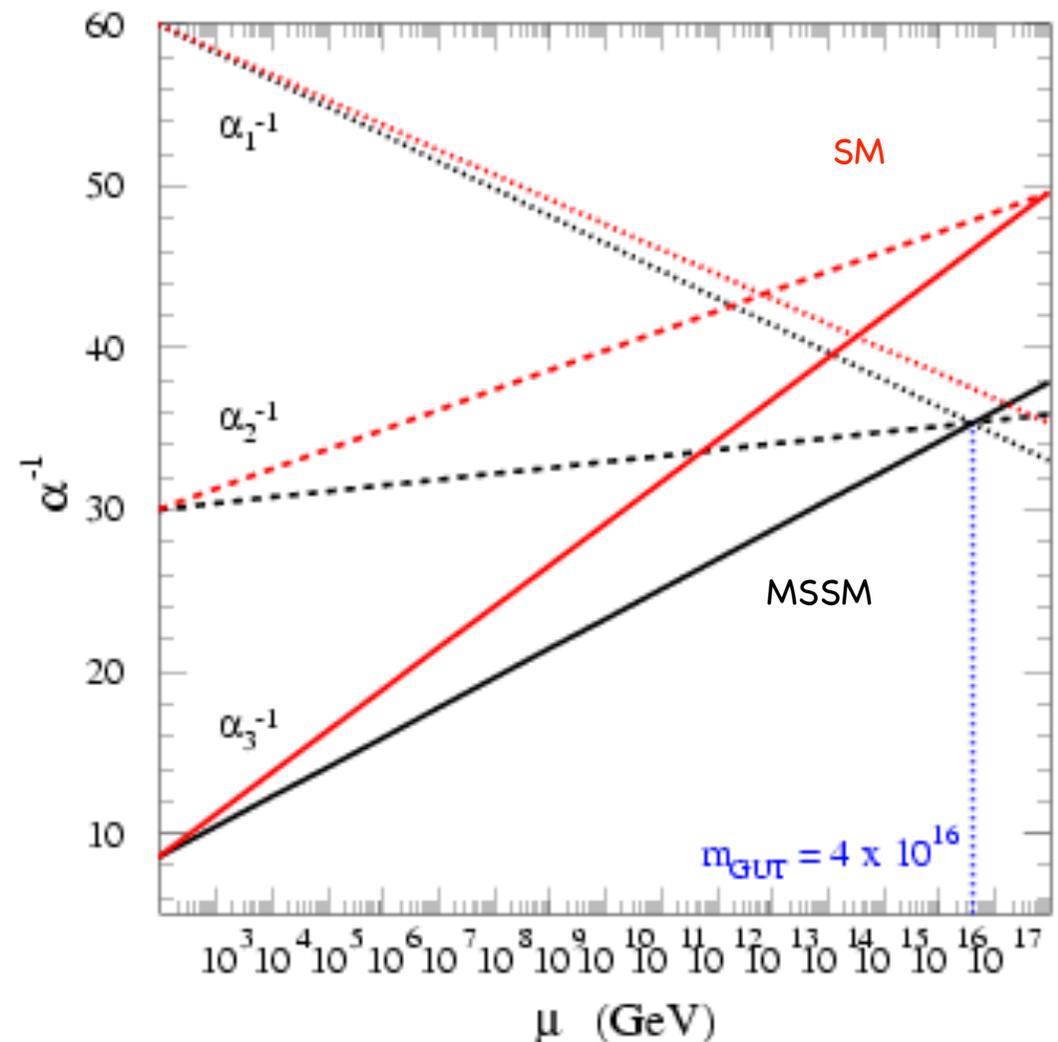


$$\sim -\frac{g_{\text{GUT}}^2}{16\pi^2} M_{\tilde{X}}^2$$

$$|M_X^2 - M_{\tilde{X}}^2| \lesssim v^2: \text{No fine-tuning!}$$

Gauge Coupling Unification

- A by-product of the Minimal Supersymmetric Standard Model is improved gauge coupling unification.
- In the SM, the gauge couplings approach each other, but don't actually converge.
- Somewhat miraculously, adding the MSSM particle content at the $\sim \text{TeV}$ scale deflects the running and produces a convergence of the couplings.
- Is this a mirage or nature telling us something?



R-Parity

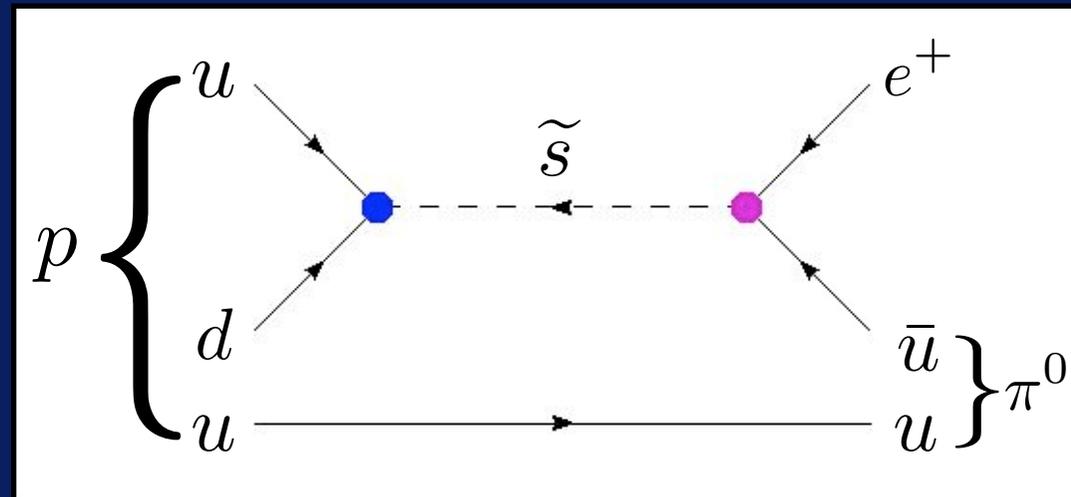
By itself, supersymmetry does not imply a stable massive particle.

It has interactions which would naively violate baryon and lepton number, and do scary things like make protons decay.

The usual take on this is to simply forbid all of these interactions by invoking a symmetry: R-parity.

R-parity insures that the superpartners only couple in pairs to the SM.

It produces a stable particle!



$$R_P \equiv (-1)^{3(B-L)+2S}$$

SM particles: +1

Superpartners: -1

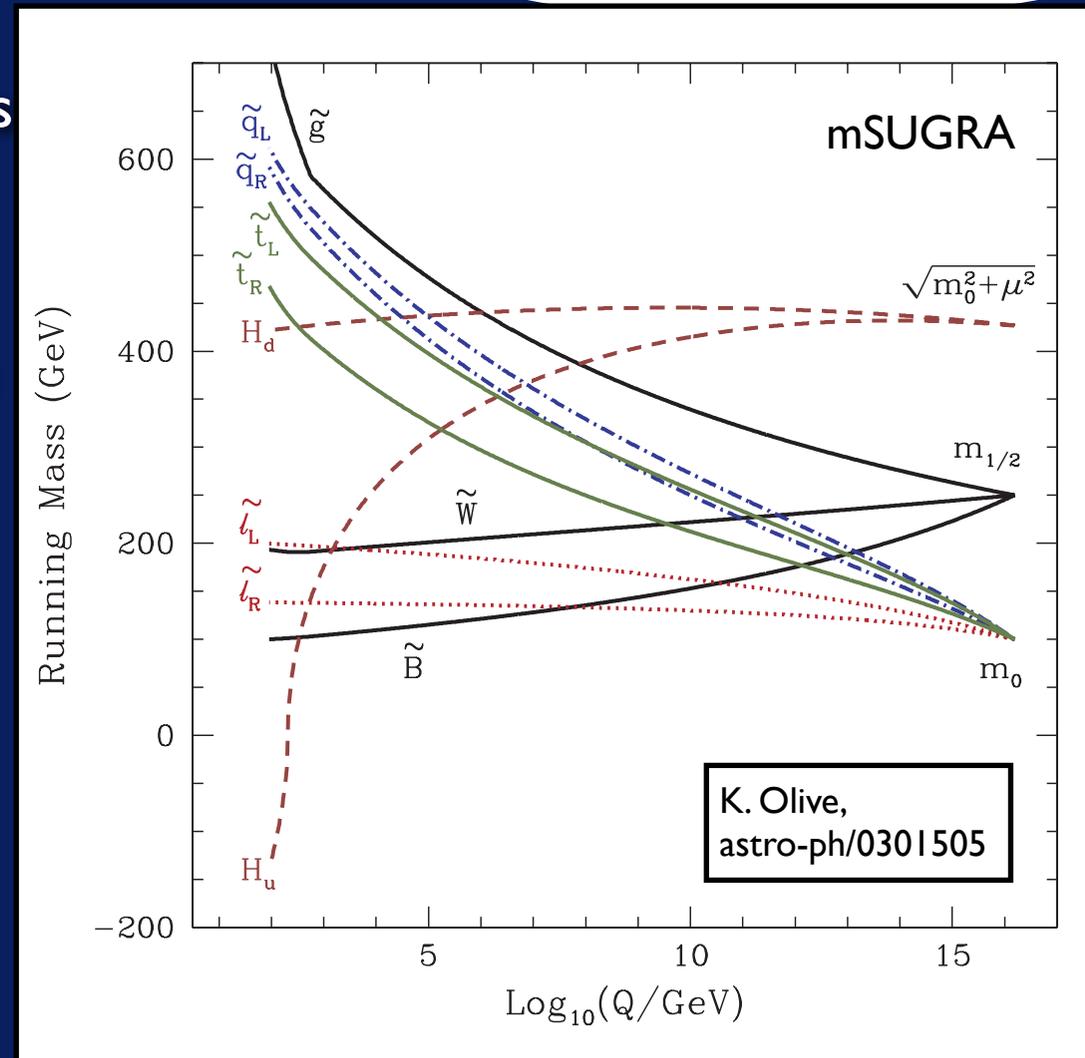
Identity of the LSP

If the Lightest Supersymmetric Particle is stable, any superpartners present in the early universe will eventually decay into them.

The LSP had better turn out to be neutral if we would like it to play the role of dark matter.

For a given model of SUSY breaking, we can calculate the spectrum and determine which particle is the lightest.

In fact, there are some generic trends that come about from the renormalization group.



Neutralino Dark Matter

- In the MSSM, the 4 neutralinos are Majorana fermions which are mixtures of the superpartners of W_3 , B, and the two neutral Higgses.

- As a result, their interactions are a little complicated: it depends on what admixture of each state is present.

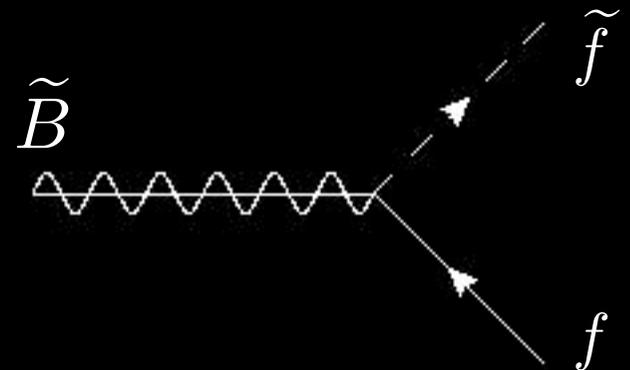
- The RGEs typically result in an LSP which is mostly Bino, with a small amount of Higgsino and W_3 ino.

- Specific models of SUSY breaking may upset these expectations.

- AMSB: W_3 ino WIMP

$$\tilde{\chi}_1^0 = N_{11}\tilde{B} + N_{12}\tilde{W}_3 + N_{13}\tilde{H}_1^0 + N_{14}\tilde{H}_2^0$$

Bino: Couples to $g_1 Y$
(interactions with the SM involve the sfermions)



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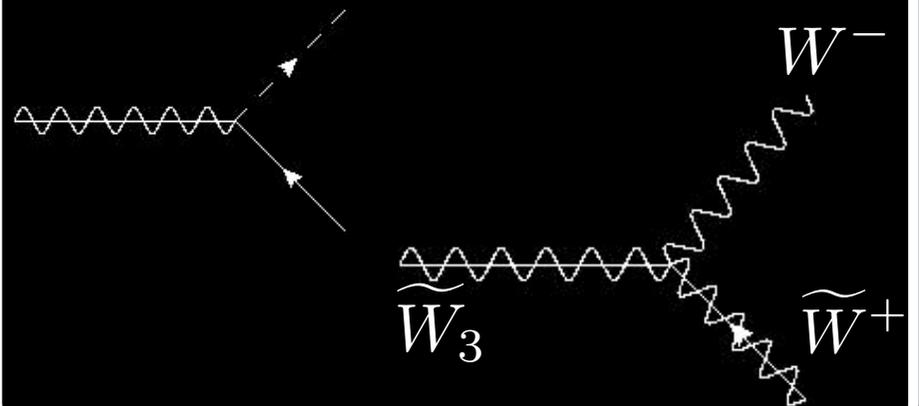
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W_3 ino: Couples to $g_2 T_3$
(interactions with sfermions and W -- not Z ! -- bosons)



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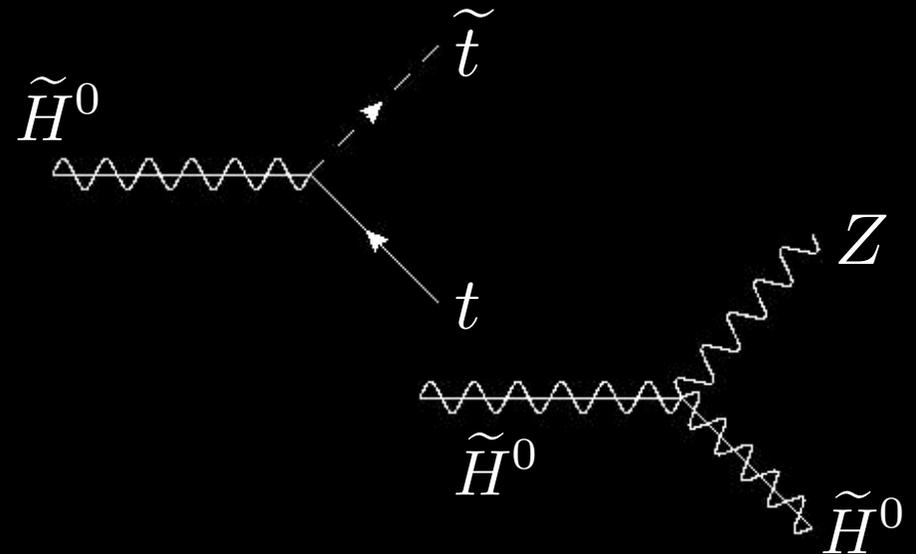
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Higgsino: Couples to massive particles



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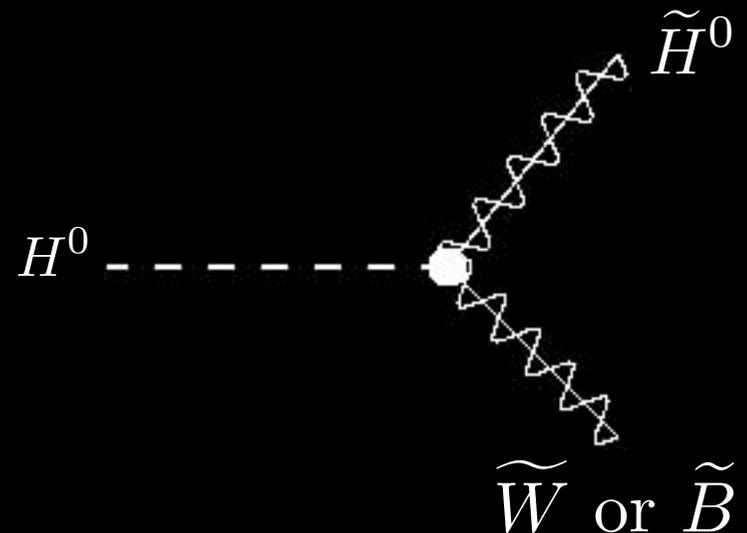
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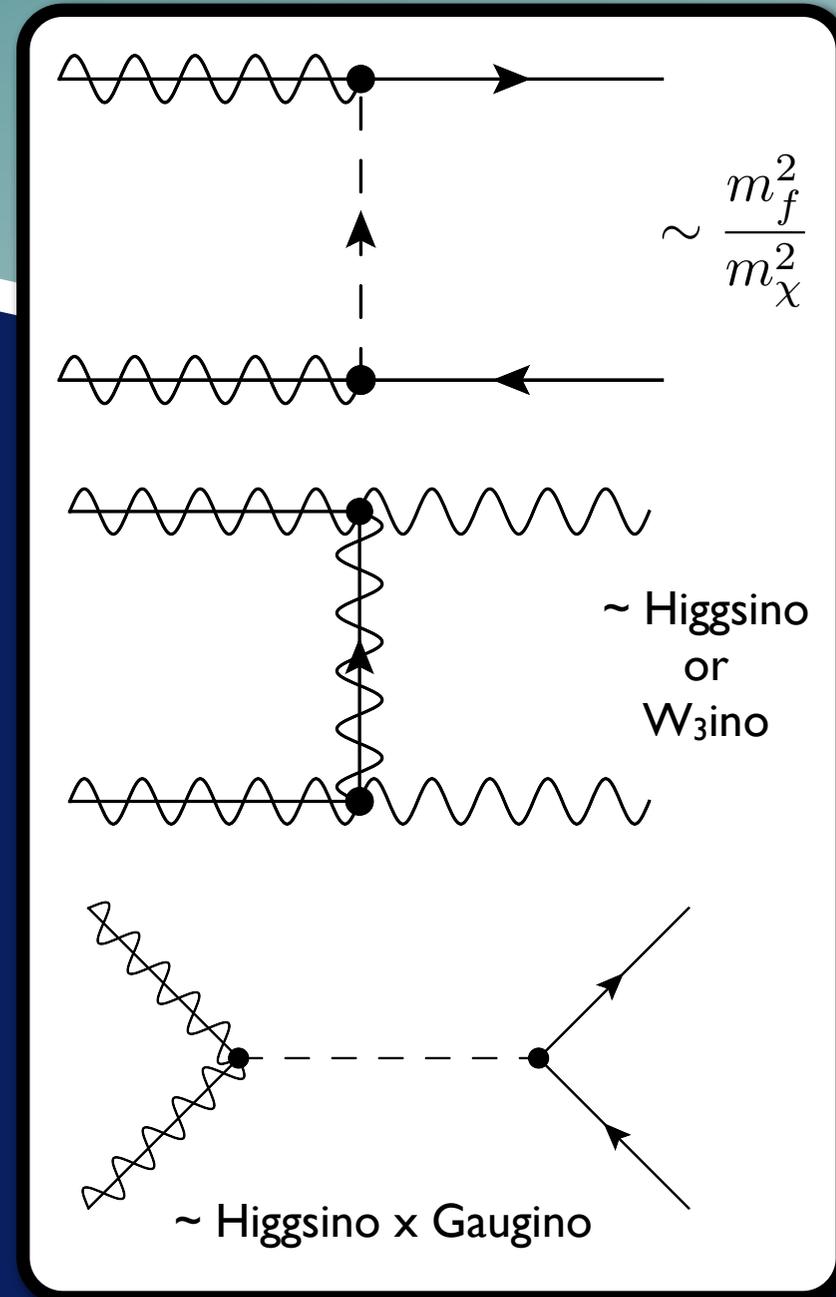
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Higgs interactions are hybrids...



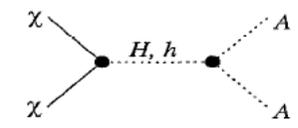
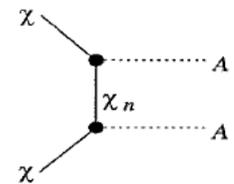
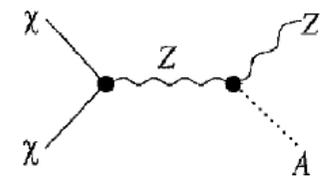
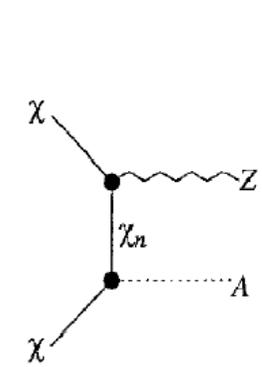
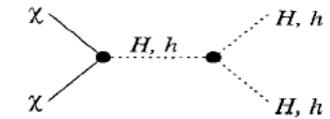
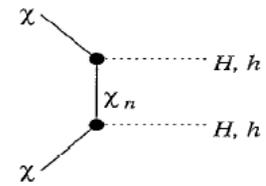
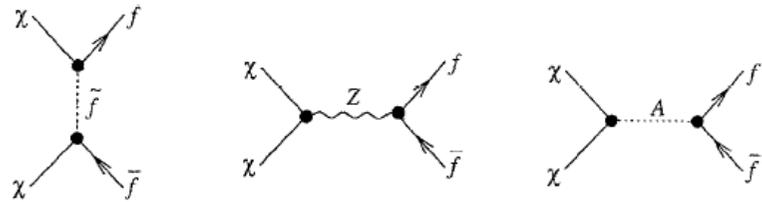
Annihilation

- Now we have everything we need to look at neutralino annihilations. This is a complicated process... but we can understand some general features.
- Neutralinos are Majorana fermions.
- In the non-relativistic limit, they are Pauli-blocked from an initial $S=1$ state.
- No annihilation through an s-channel vector particle.
- Sfermion exchange likes to produce SM fermions of like-chirality, ($S=1$) and is suppressed by m_f for an $S=0$ initial state.

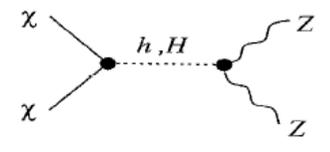
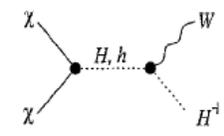
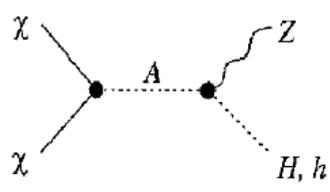
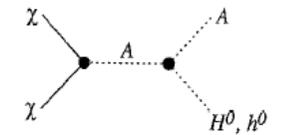
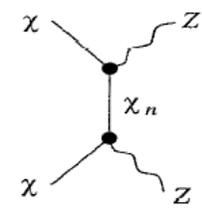
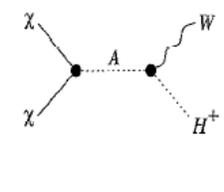
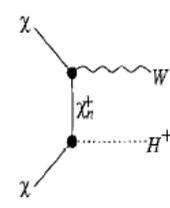
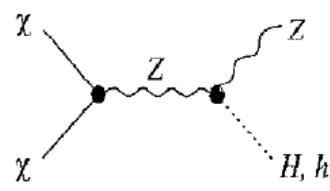
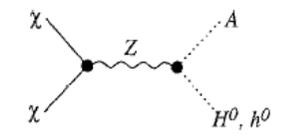
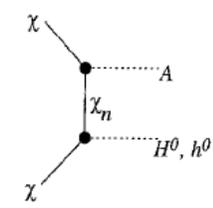
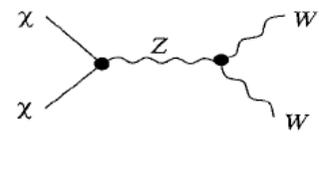
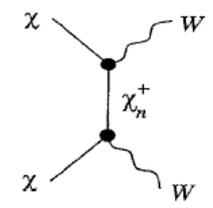
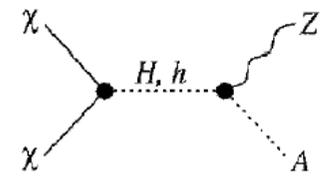
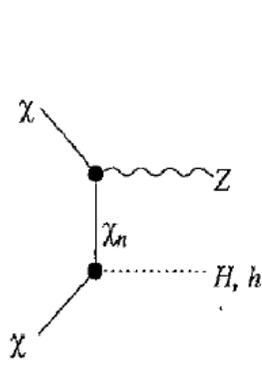


Bottom Line: Suppressed σv leads to generically too many Binos.

A Plethora of Processes

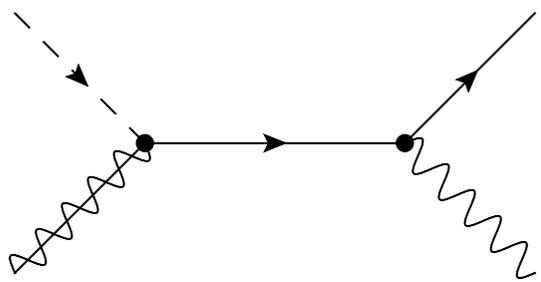


Jungman, Kamnionkowski, Griest,
Physics Reports'95

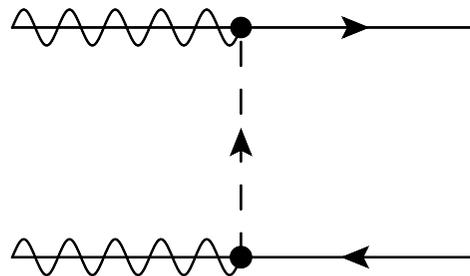


Relic Density: Small $\tan \beta$

mSUGRA

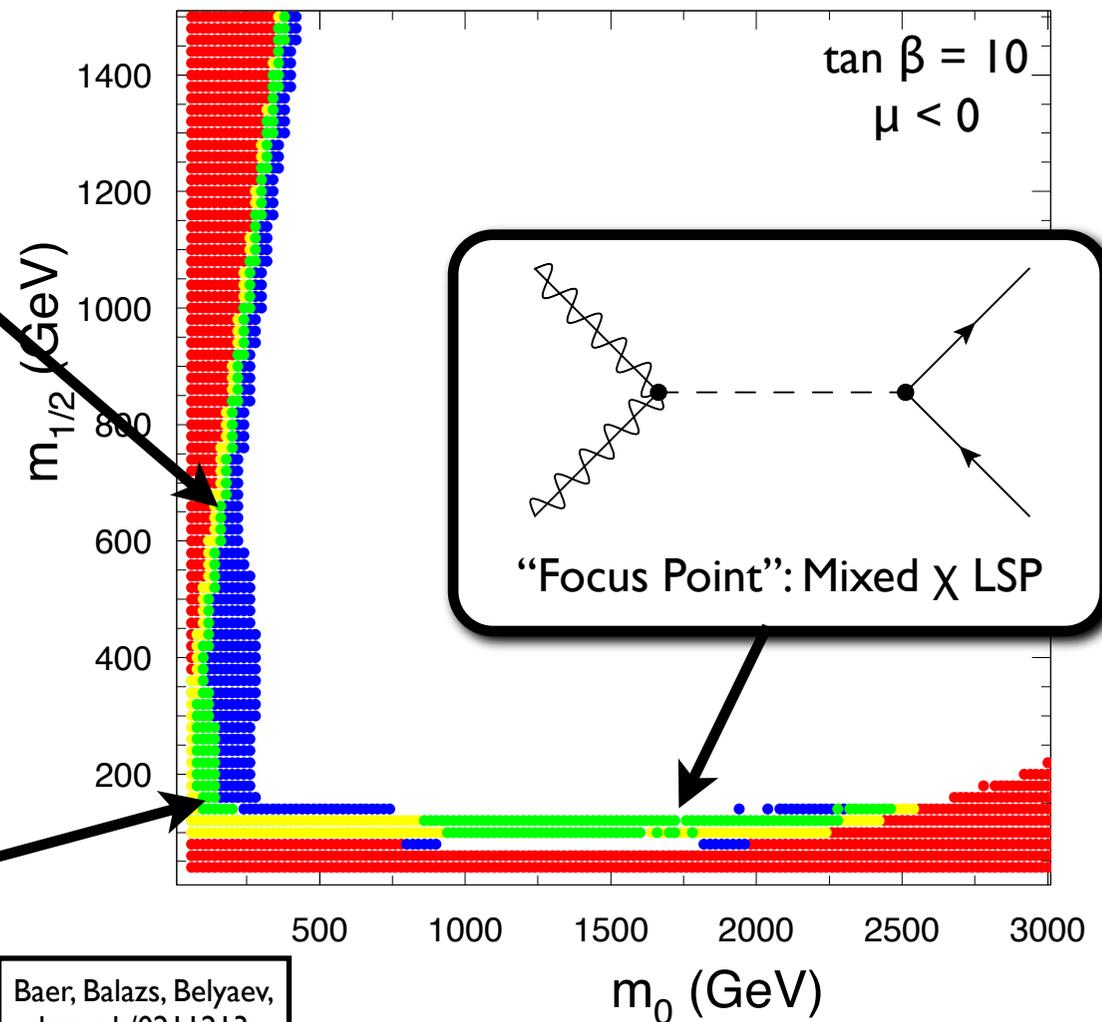


“Coannihilation Region”:
Degenerate stau active during
freeze-out



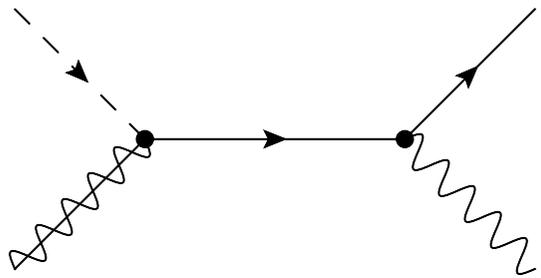
“Bulk Region”: Light sfermions
(Under assault by LHC)

- $0.3 < \Omega h^2 < 1.$
- $\Omega h^2 < 0.1$
- $0.1 < \Omega h^2 < 0.3$
- excluded by theory

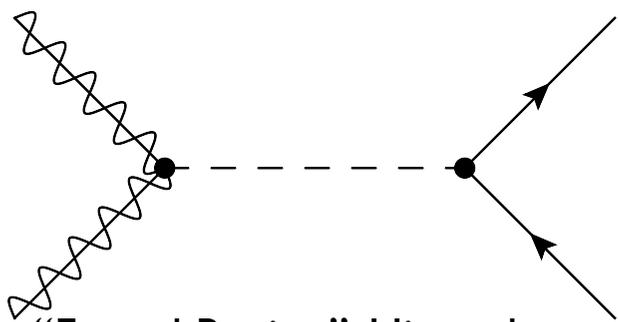


Large $\tan \beta$

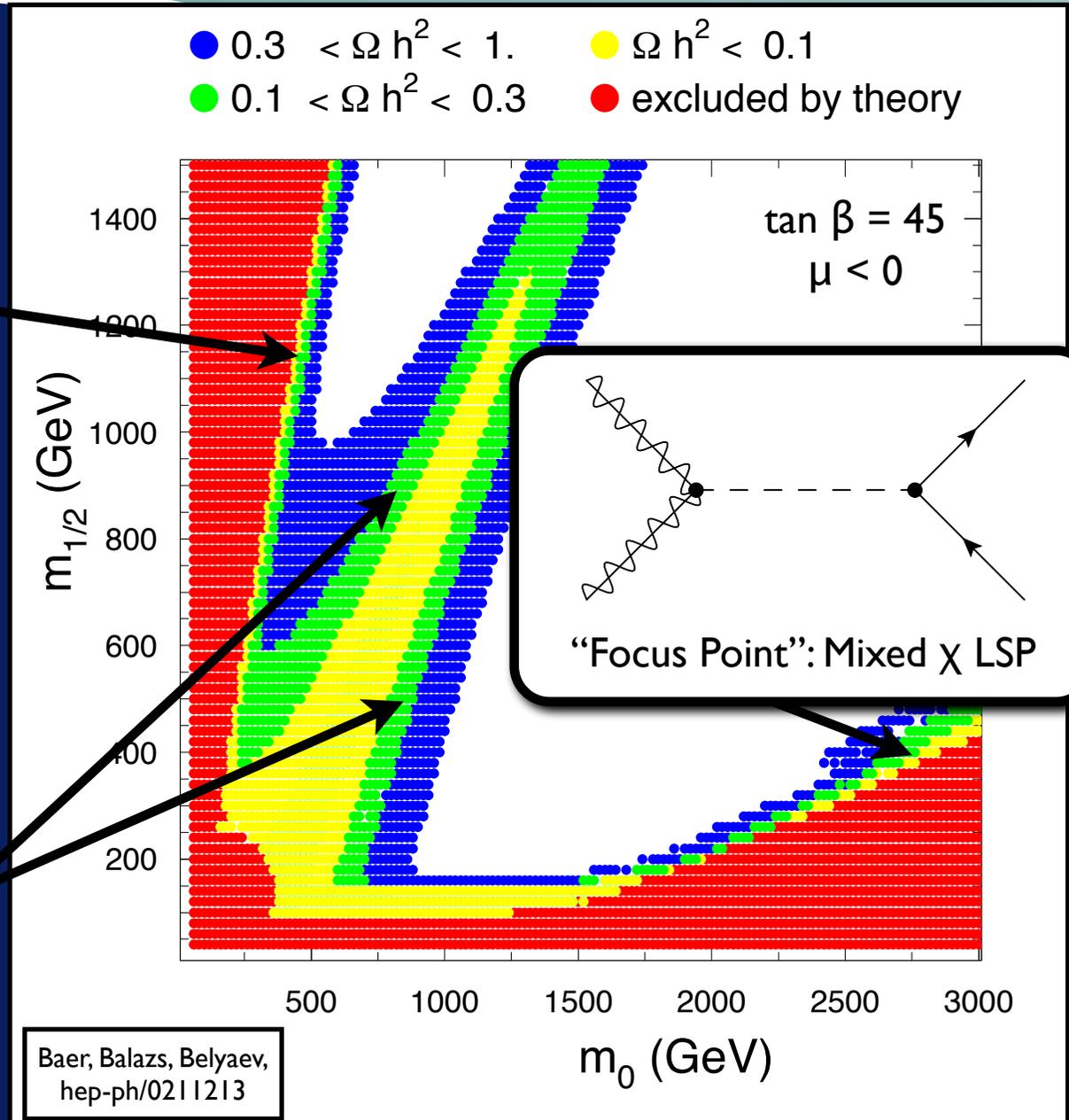
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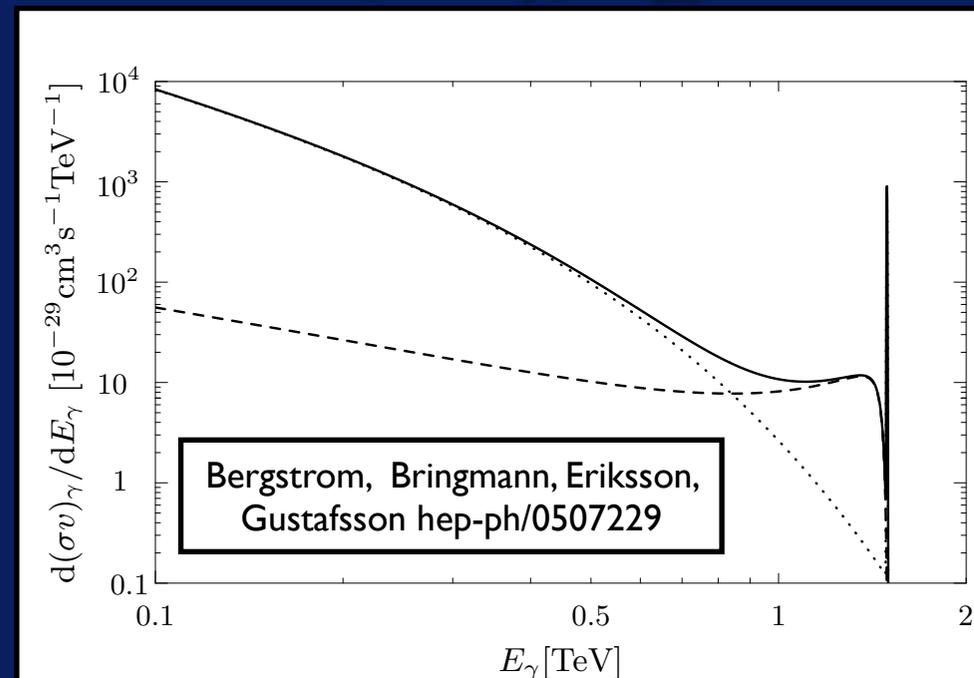
“Funnel Region”:
Higgs close
to on-shell in decay



Cosmic Neutralino Signals

- We've already learned a fair amount about how neutralinos annihilate by studying the relic density.
- The same physics controls the search for them annihilating in the halo.
- As Majorana particles, they tend to annihilate into heavier fermions and/or W bosons.
- Fermi searches for bb spectra...
- Loops of charged particles allow them to annihilate into $\gamma\gamma$ or γZ .
- A “smoking gun” signal!

1.5 TeV (Mostly) Higgsino LSP



Cosmic Neutralino Signals

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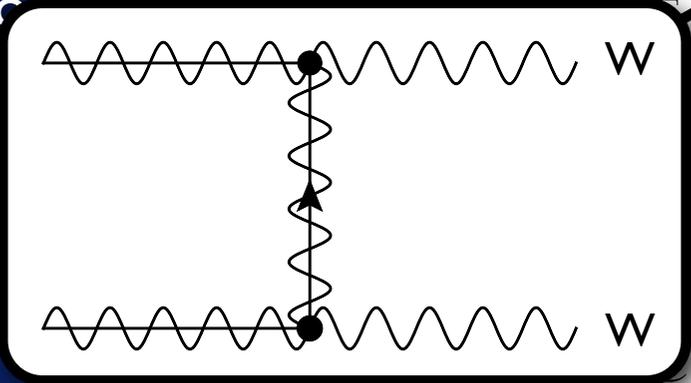
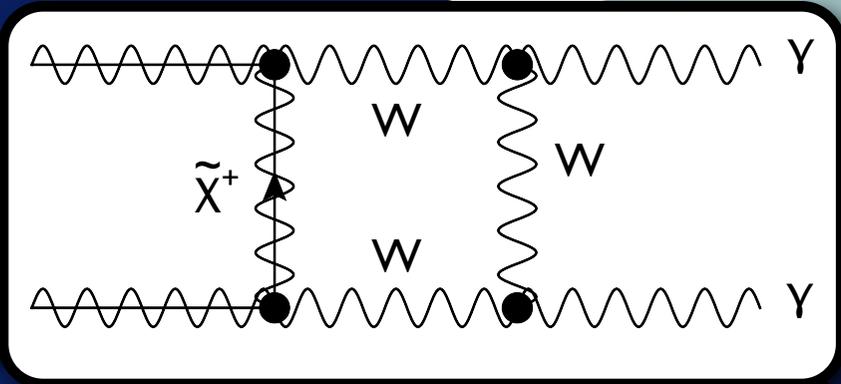
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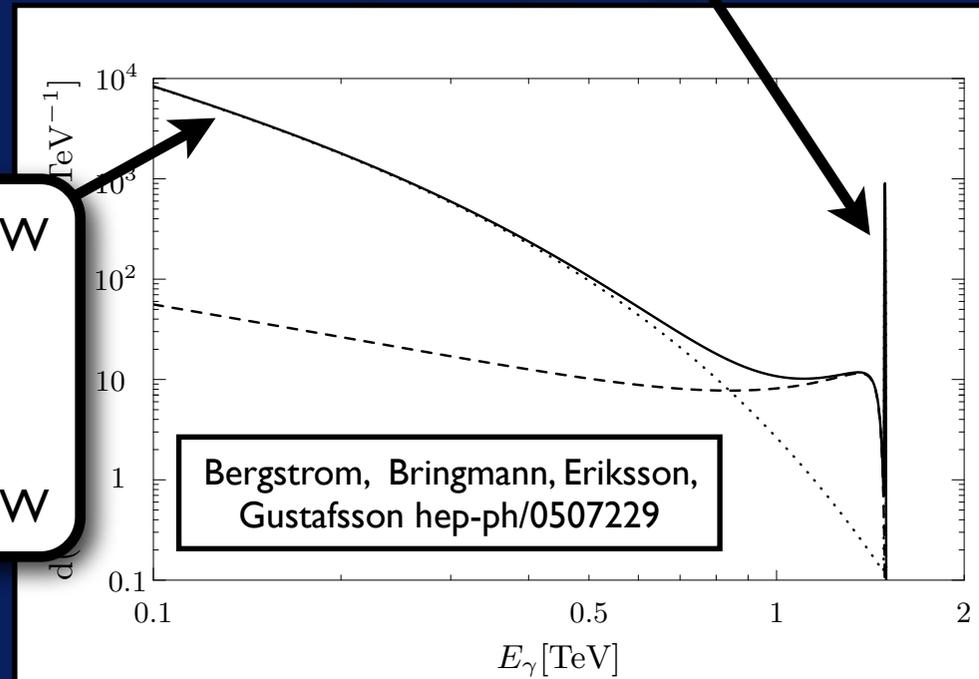
Fermi selection rules

Loops of charged particles can allow them to annihilate into photons.

A "smoking gun" signal!



1.5 TeV (Mostly) Higgsino LSP



Direct Detection

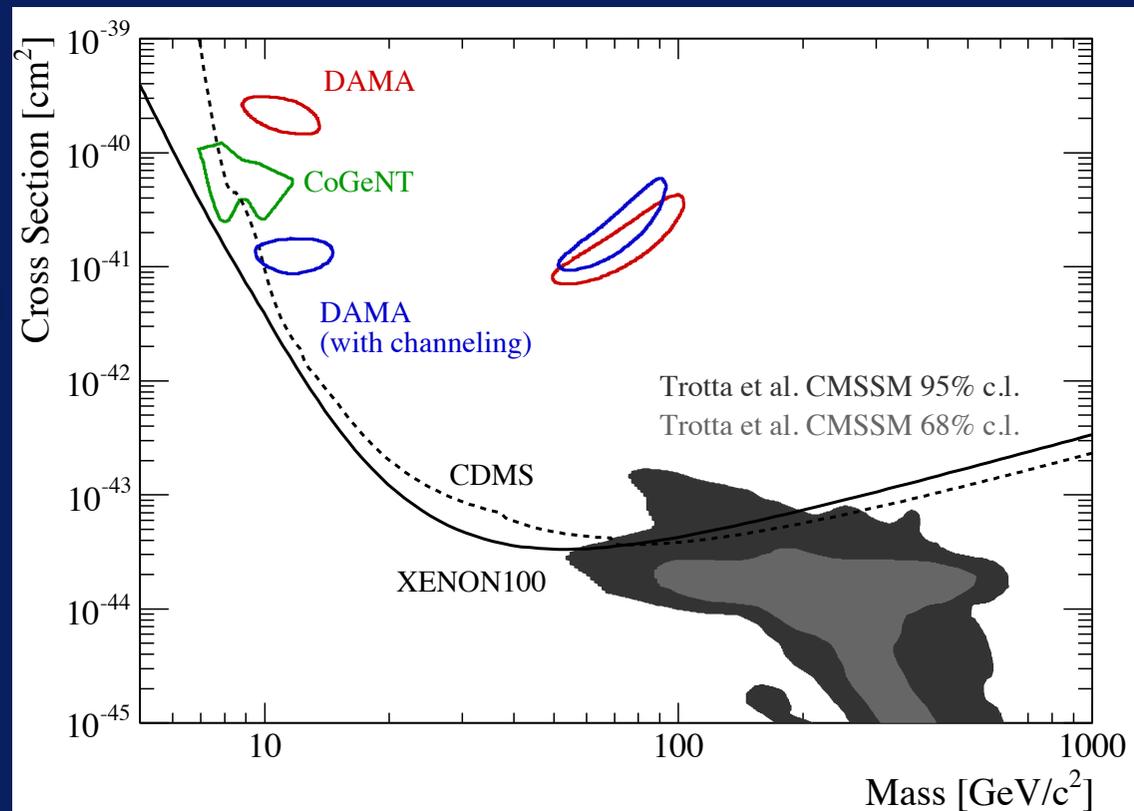
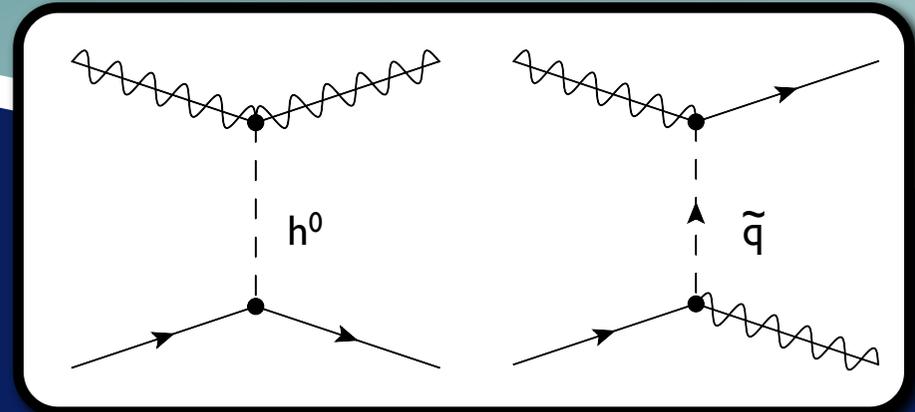
- The Majorana character also has important consequences for direct detection.

- No vector currents imply the Z exchange can only mediate spin-dependent interactions.

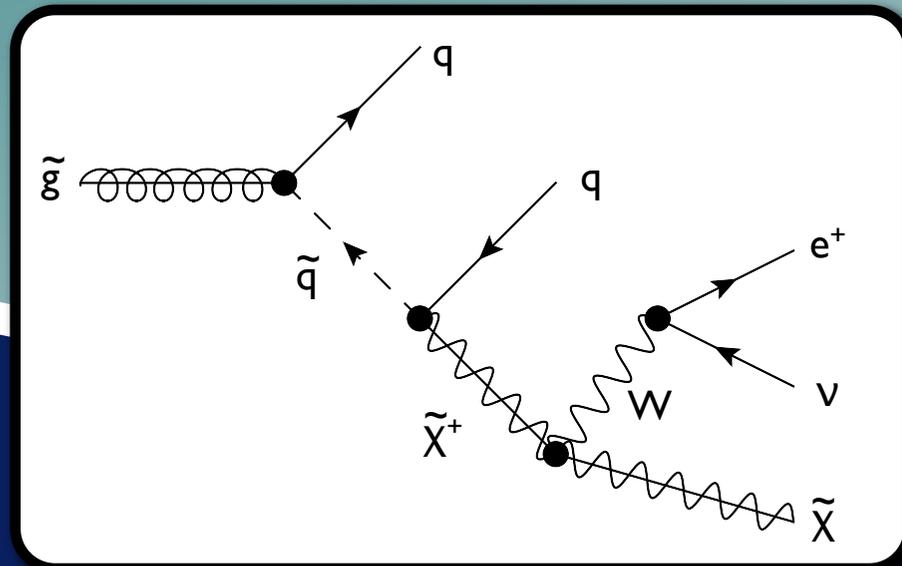
- The Higgs exchange requires both gaugino and higgsino admixture.

- (MSSM has a light Higgs!)

- Direct detection is sensitive to MSSM parameter space!



Collider Signals



At hadron colliders like the LHC, the largest signals tend to come from producing the colored superpartners.

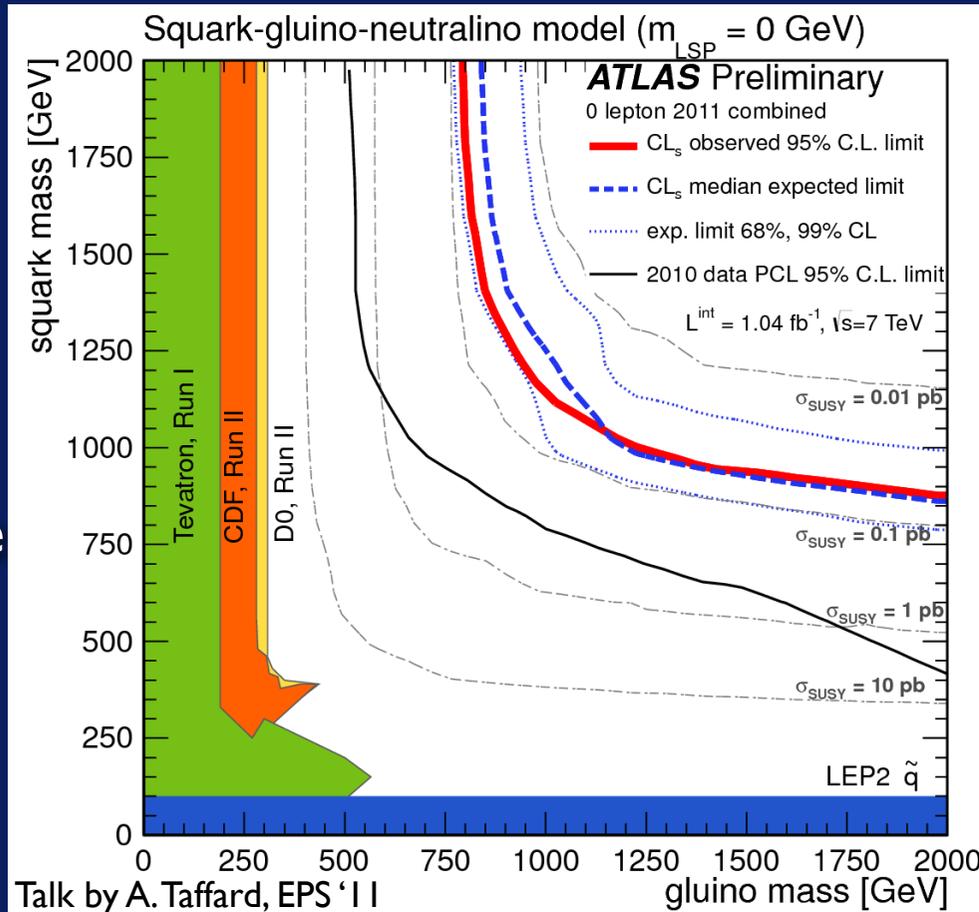
“Cascade” decays down to the LSP.

The LSP passes through the detector, leading to missing momentum.

Hard jets are also present.

Depending on the decay chain, there may be hard leptons as well.

Often pairs of leptons will have the same charge, a signal with small expected SM backgrounds.



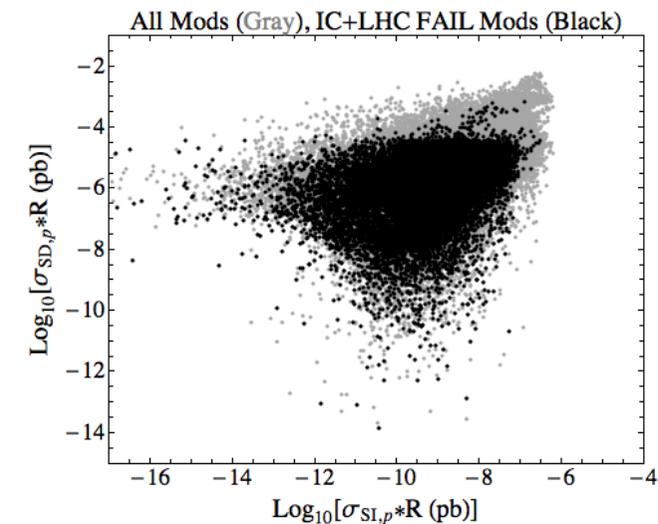
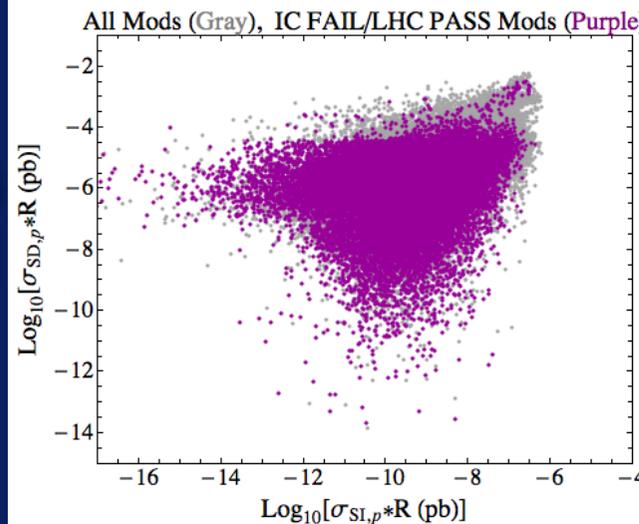
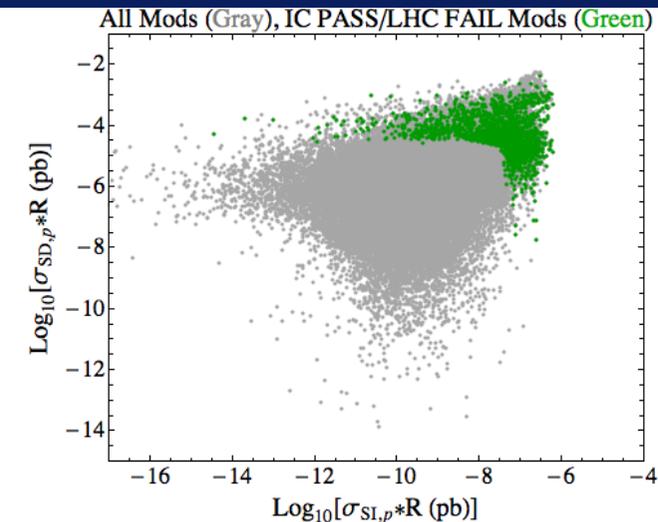
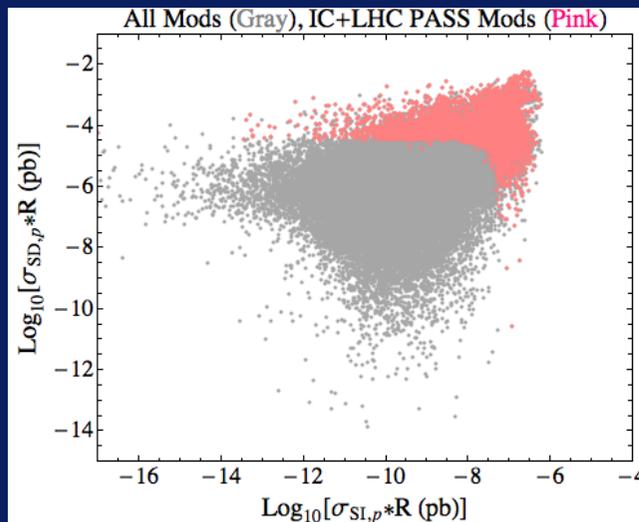
Reconstructing the MSSM

While we can hope to eventually have many, many signals to measure, the parameter space is also very large.

Even “simplified” versions have ~ 20 parameters!

Mapping from signal to parameter space is very complicated and not generally one to one.

Neutrinos from Ice Cube
Cotta, Howe, Hewett, Rizzo arXiv:1105.1199



Reconstructing SUSY DM

With enough measurement, we can reconstruct the SUSY model sufficiently to compute the relic density, and see if it checks out.

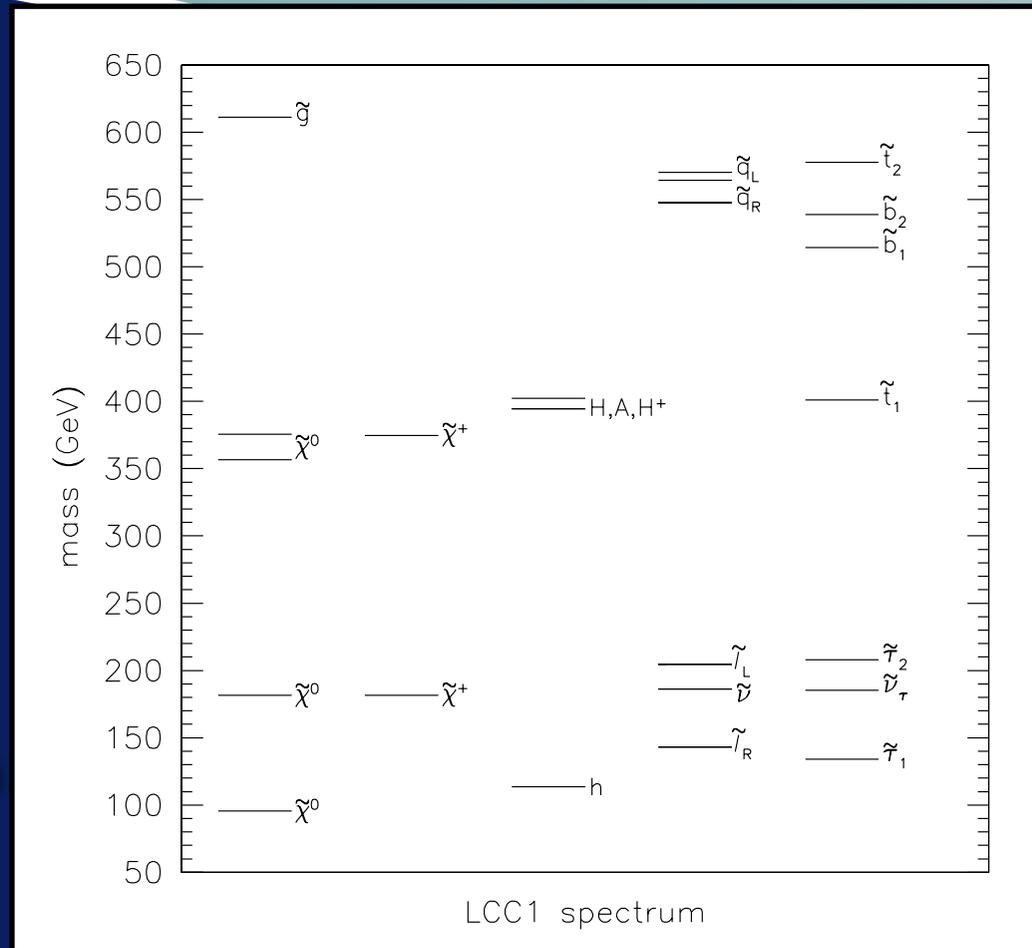
This is a precision measurement!

Strong dependence on the composition of the LSP.

(Perhaps) strong dependence on slepton masses.

The funnel region depends on the Higgs couplings.

A precision collider such as an ILC can help a lot...



Baltz, Battaglia, Peskin, Wizansky, hep-ph/0602187

LCC1: Model with light sleptons.

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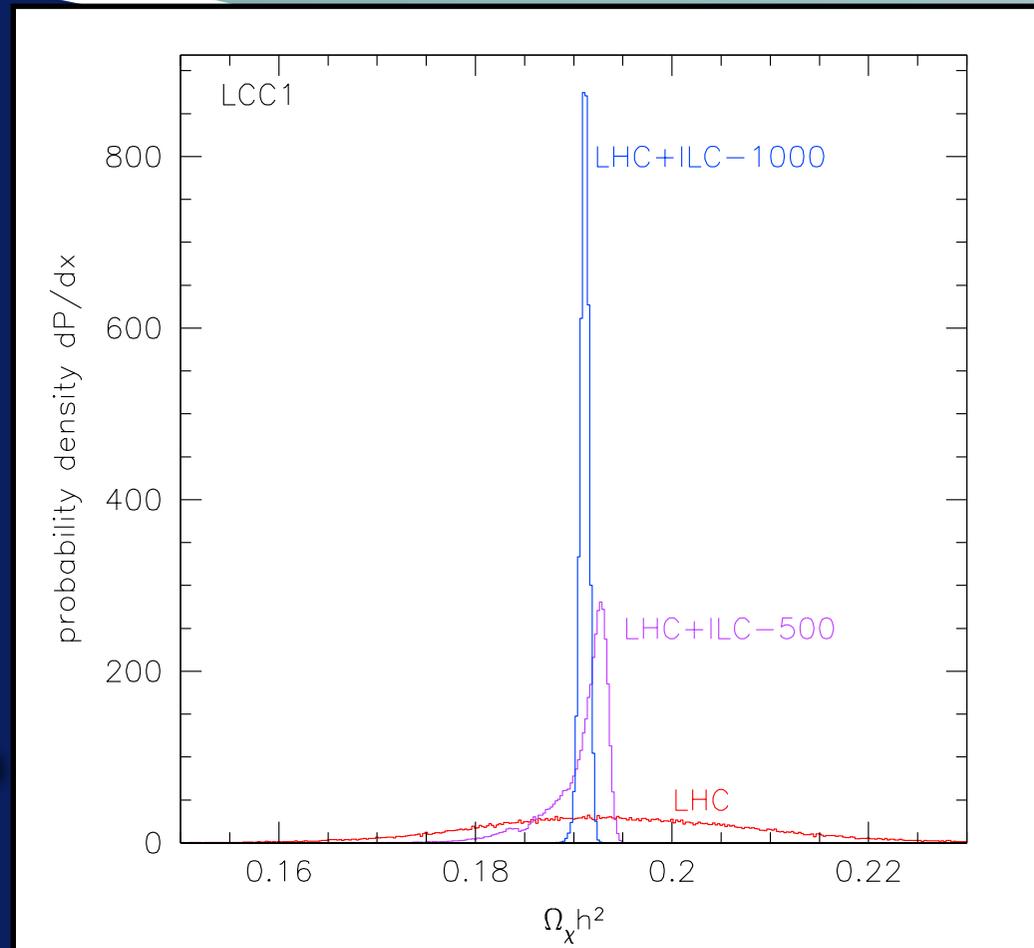
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LCCI: Model with light sleptons.

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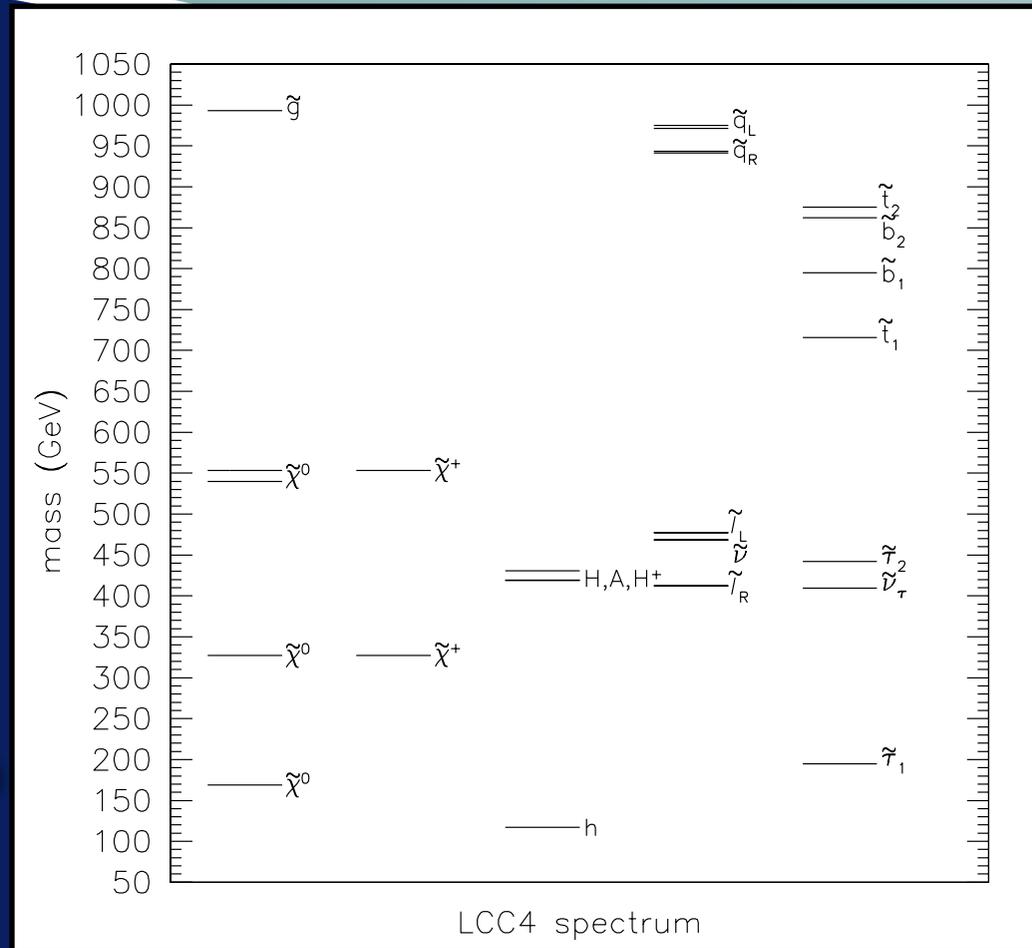
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LCC4: Annihilation through A^0 .

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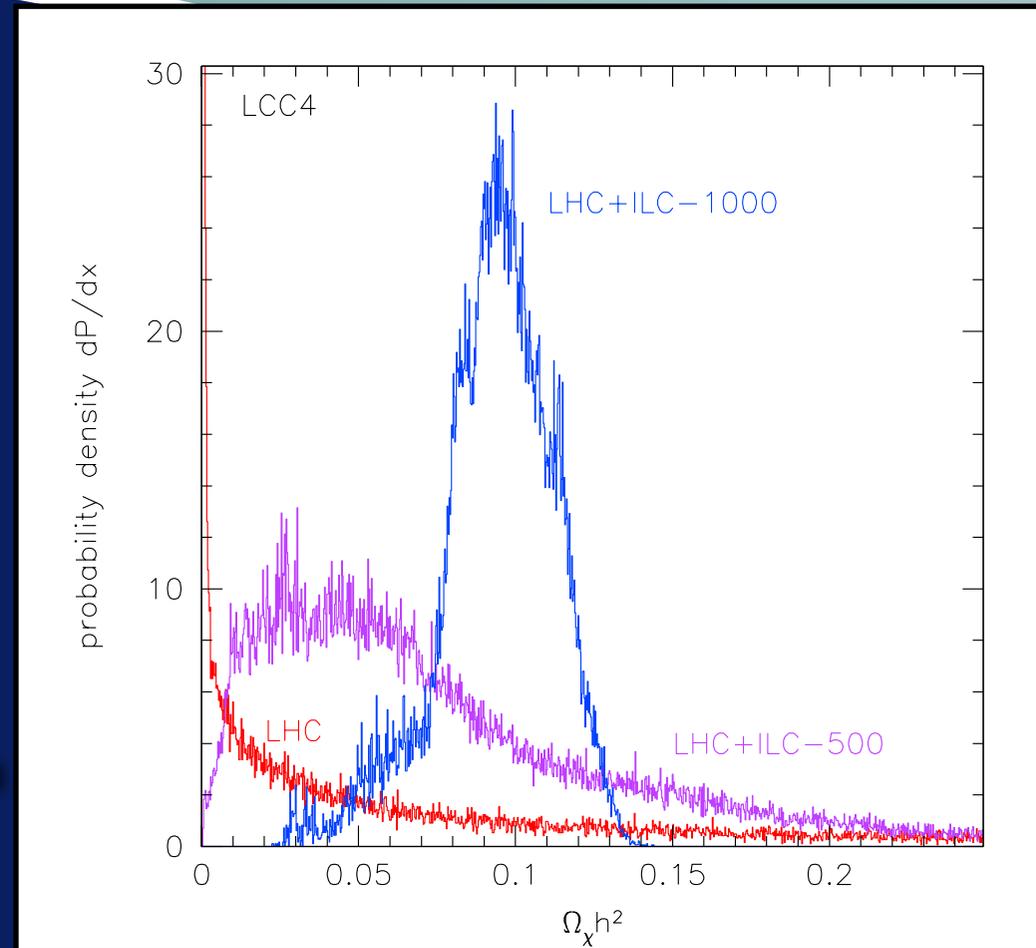
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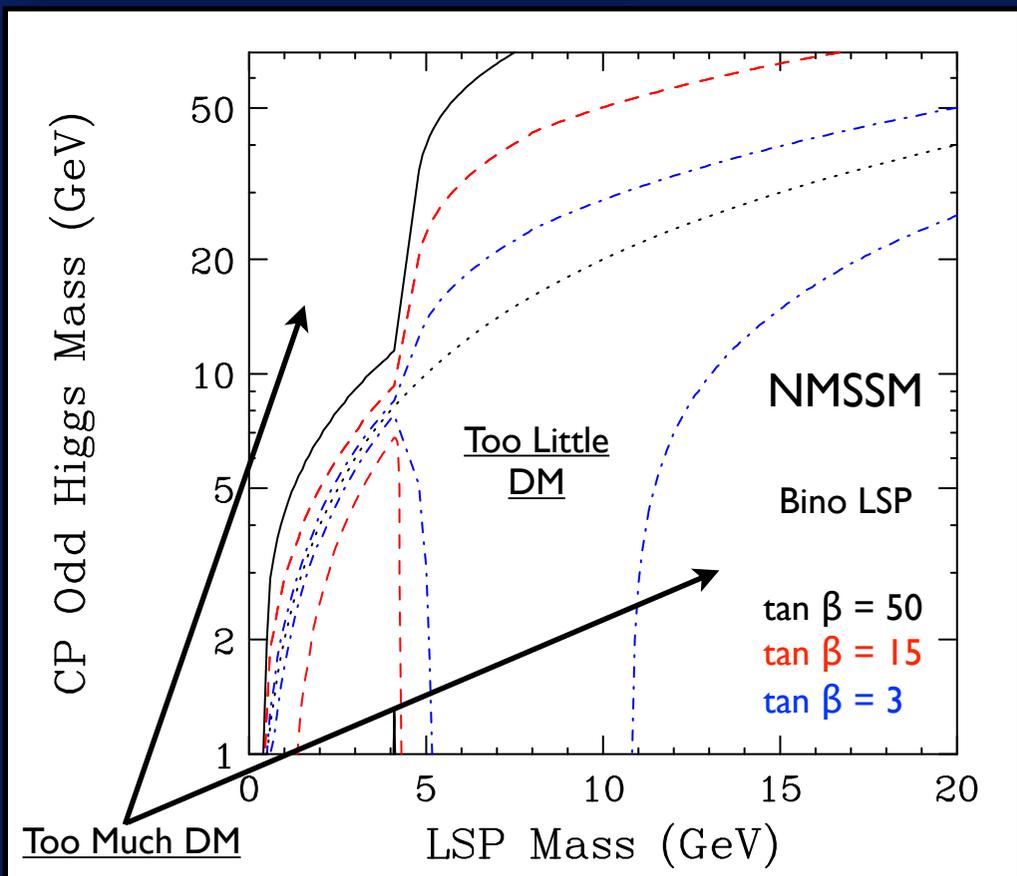
Baltz, Battaglia, Peskin, Wizansky, hep-ph/0602187

LCC4: Annihilation through A^0 .

Beyond the MSSM

- As we have seen, the minimal model already contains a lot of interesting physics.
- But nothing tells us Nature chose something minimal!
- Simple extensions such as adding a gauge singlet (i.e. the NMSSM) can have a big impact on the picture of dark matter.
- New neutralinos
- New Higgs bosons
- New couplings

Gunion, Hooper, McElrath, hep-ph/0509024



Recap

- There are many ideas for what dark matter could be.
 - Dark, Neutral, and Stable [Symmetry!].
 - WIMPs are a particularly attractive class of dark matter.
 - Their relic density explains the ballpark dark matter abundance.
 - Large interactions give us handles to search for them.
- Supersymmetry is an attractive, representative theory of dark matter.
 - We can explore the features of a Majorana fermion WIMP.
 - Interesting regions with the correct relic density.
 - Distinctive signals of direct, indirect, and collider searches.
 - We'll see contrasting features tomorrow when we discuss UED DM.

Not Likely to be Machos...

